



UVA

SCHOOL *of* ENGINEERING
& APPLIED SCIENCE

Controlling heat transfer through interfacial coupling of photons, electrons and phonons



Patrick E. Hopkins

Professor

Dept. Mech. & Aero. Eng.

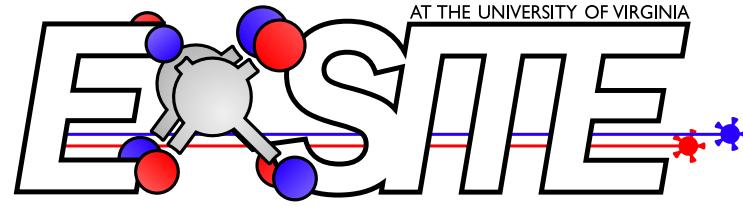
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Dept. Physics

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Group members and collaborators

Group members

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- J. Braun (Laser Thermal)
- H. Olson (Laser Thermal)
- J. Gaskins (Laser Thermal)
- K. Aryana (NASA)
- R. Cheaito (Google)
- Md. S. B. Hoque (UVA)
- W. Hutchins (UVA)
- T. Pfeifer (ONRL/UVA)



Collaborators

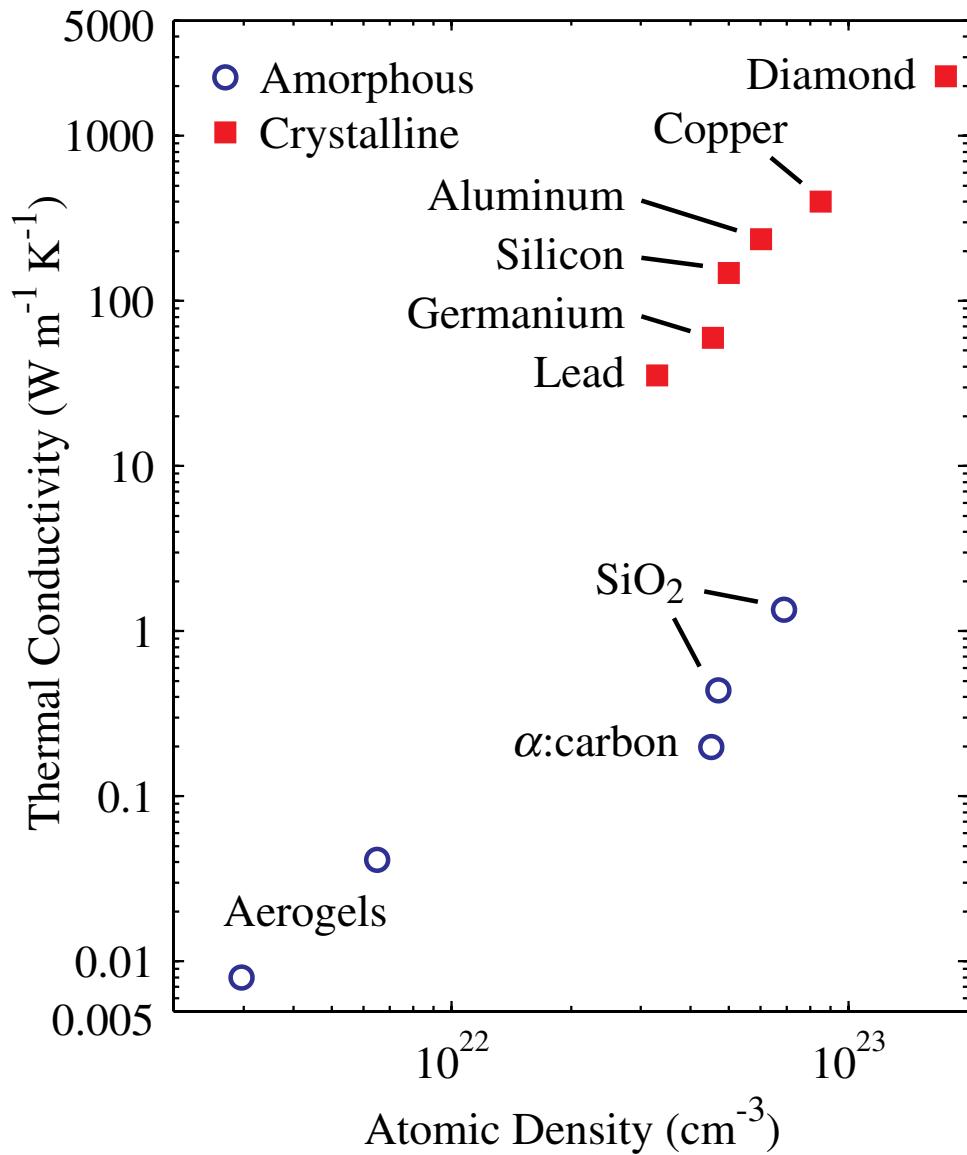
- X. Roy (Columbia)
- A. Majumdar (Stanford)
- R. Ramesh (Rice)
- J. Ravichandran (USC)
- C. Nuckolls (Columbia)
- P. Kim (Harvard)
- A. Khan (U. South Carolina)
- J.P. Maria (Penn State)
- S. King (Intel)
- S. Walton (NRL)
- O. Prezhdo (USC)
- E. Runnerstrom (ARO)
- J. Caldwell (Vanderbilt)
- S. Pantelides (Vanderbilt)

E. Hoglund (ORNL)



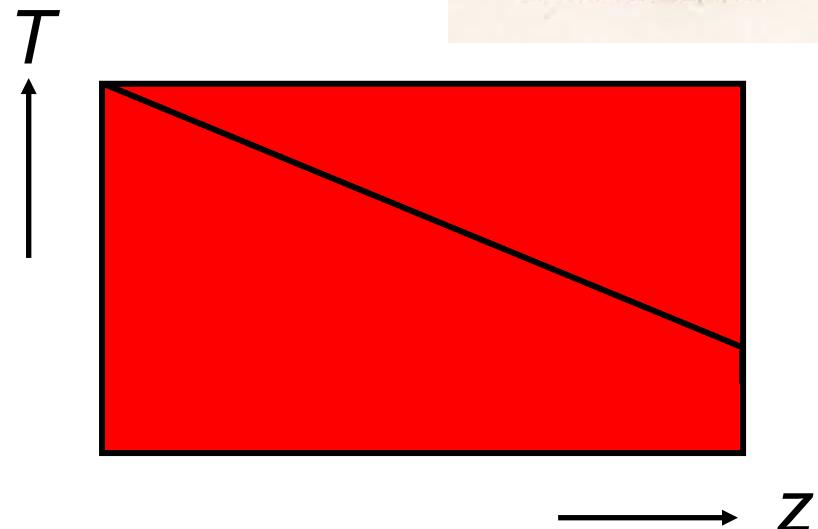
J. Hachtel (ORNL)

Thermal conductivity of materials – Macroscopic



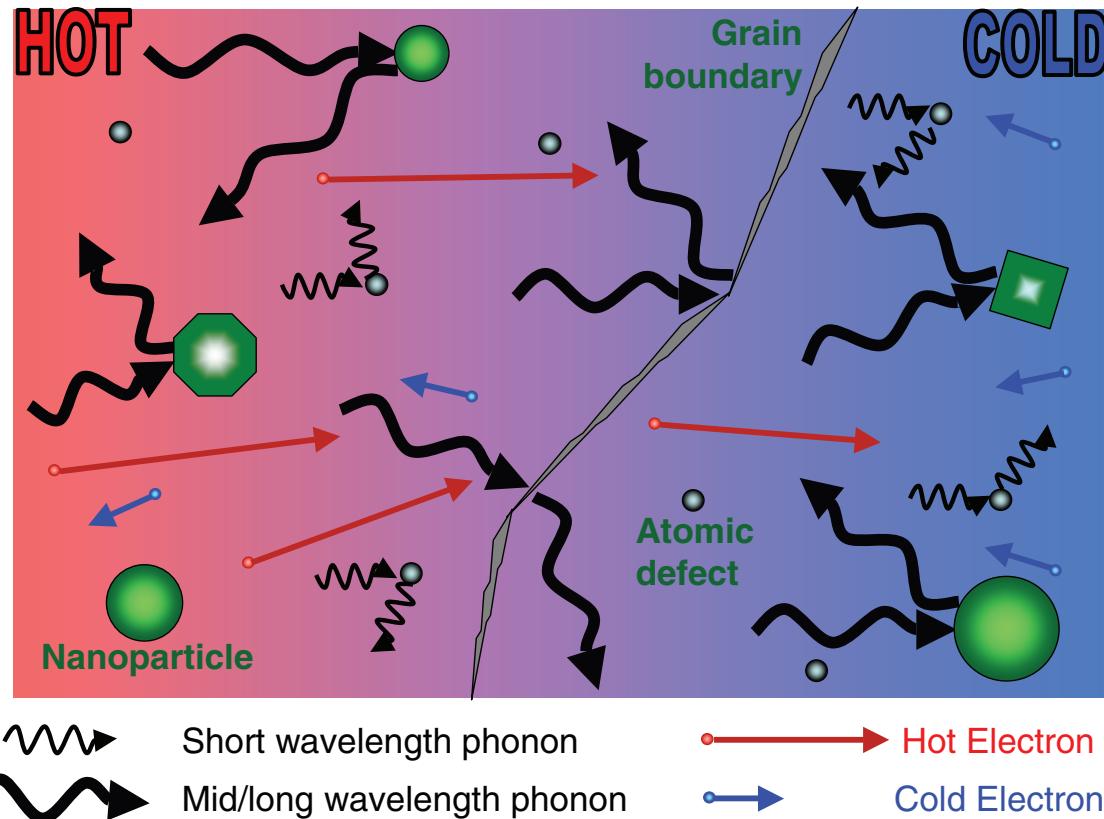
The Fourier Law

$$q = -\kappa \frac{\partial T}{\partial z}$$



A nanoscopic view with Kinetic Theory

$$\kappa = \frac{1}{3} Cv \lambda = \frac{1}{3} Cv_g^2 \tau$$

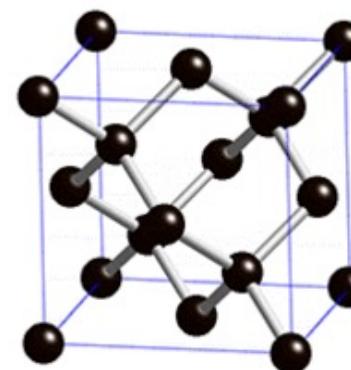
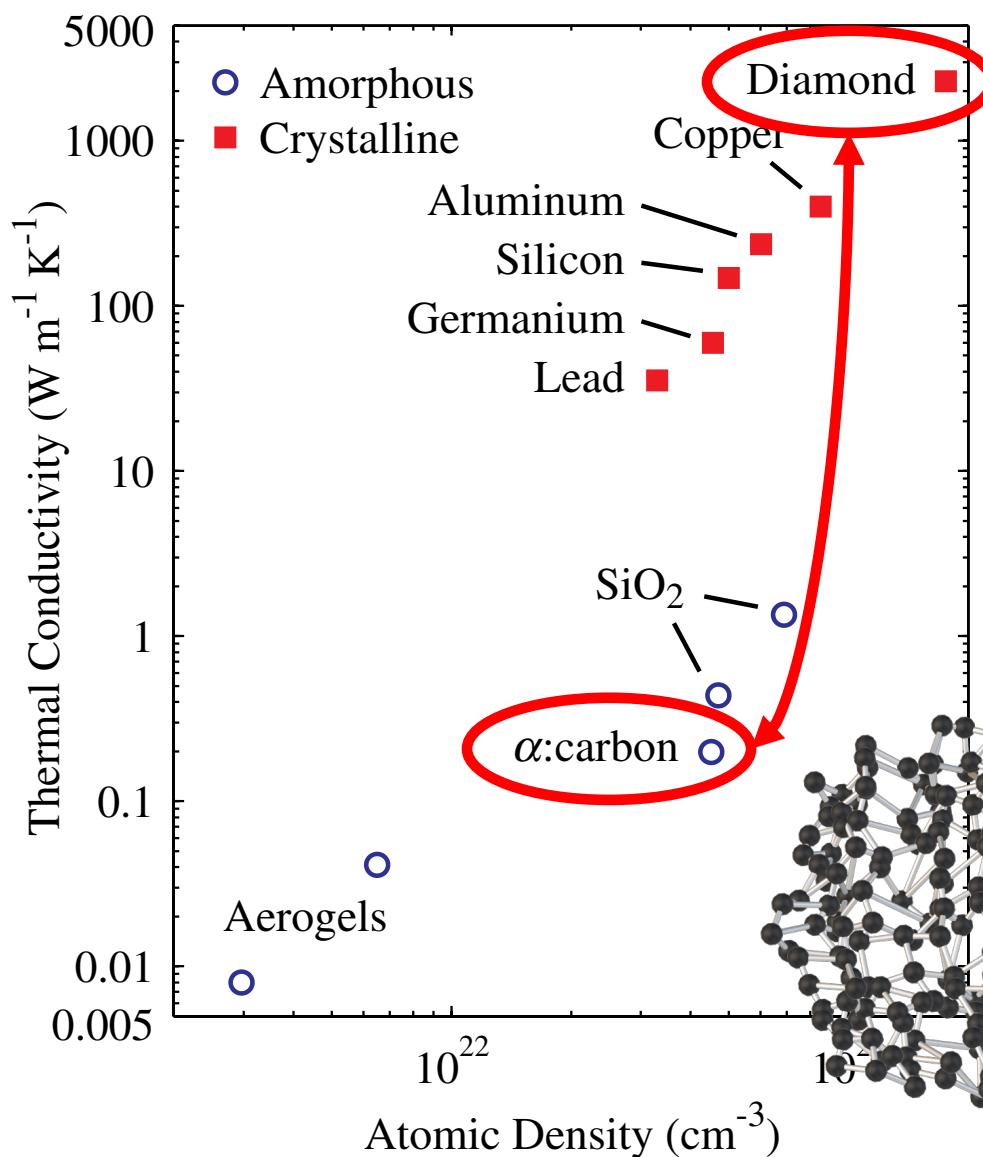


C: Heat capacity
“How much energy electrons/phonons store”

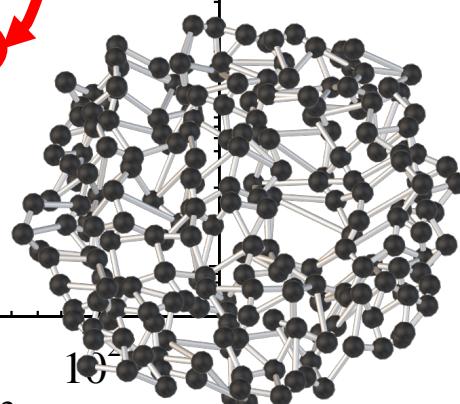
v: Velocity
“How fast the electrons/phonons move”

λ = Mean free path
“How far they move before losing energy/momentum”

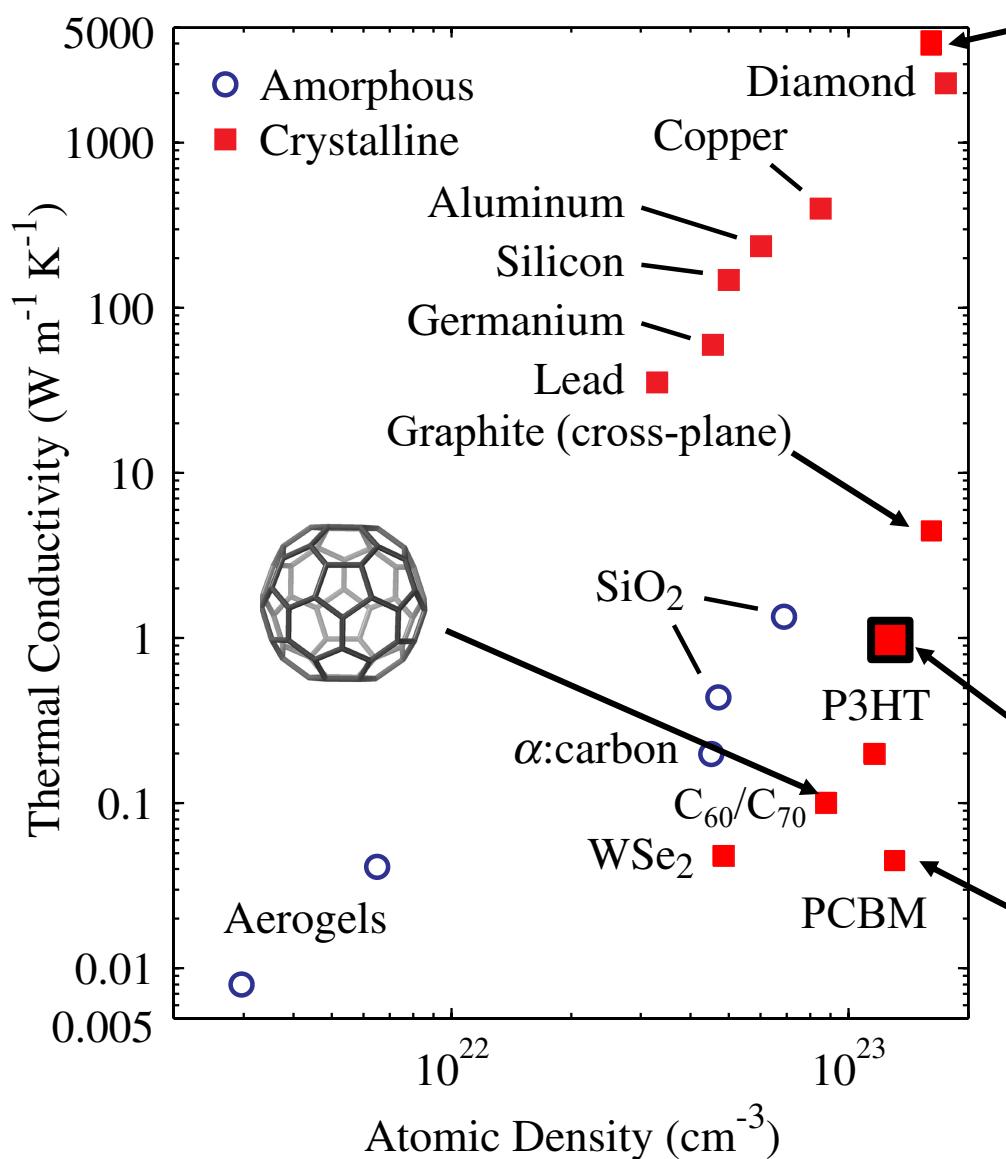
Thermal conductivity of materials – Nanoscopic



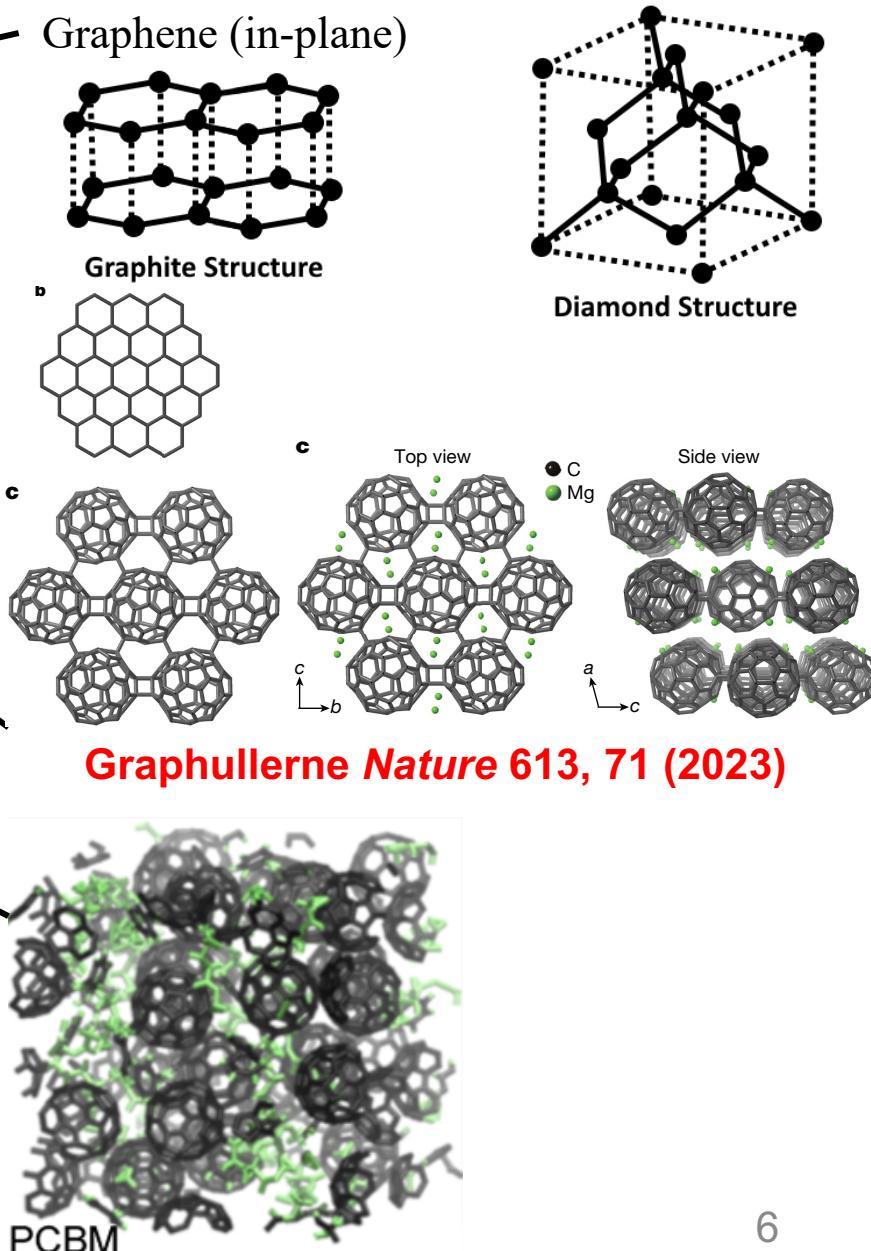
$$\kappa = \frac{1}{3} Cv\lambda = \frac{1}{3} Cv_g^2 \tau$$



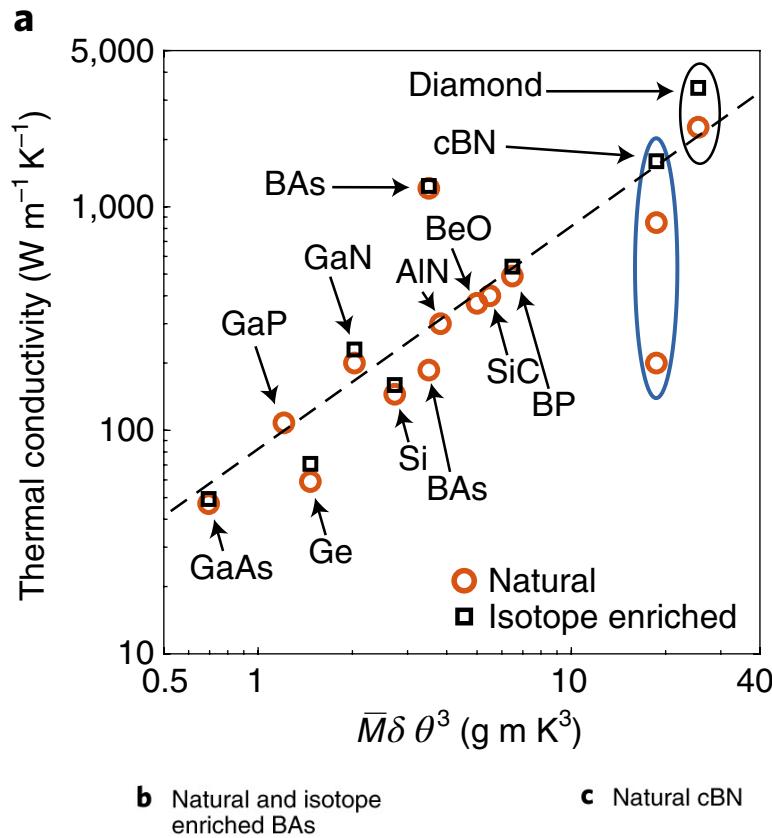
Engineering the thermal conductivity of materials: Nano HX



PRL 110, 015902 (2013)

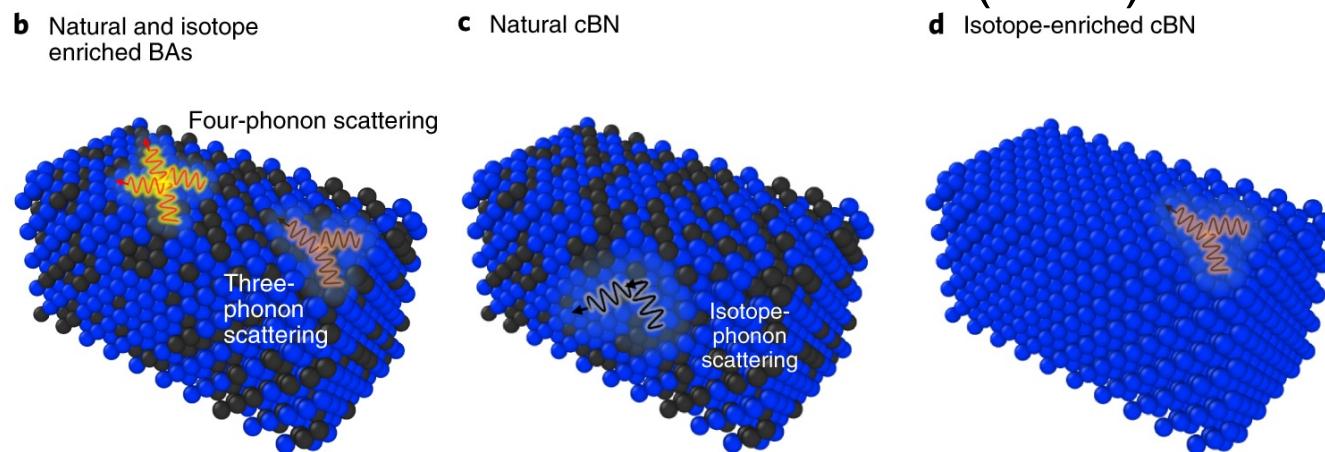


Thermal conductivity of materials: “Engineering defects”



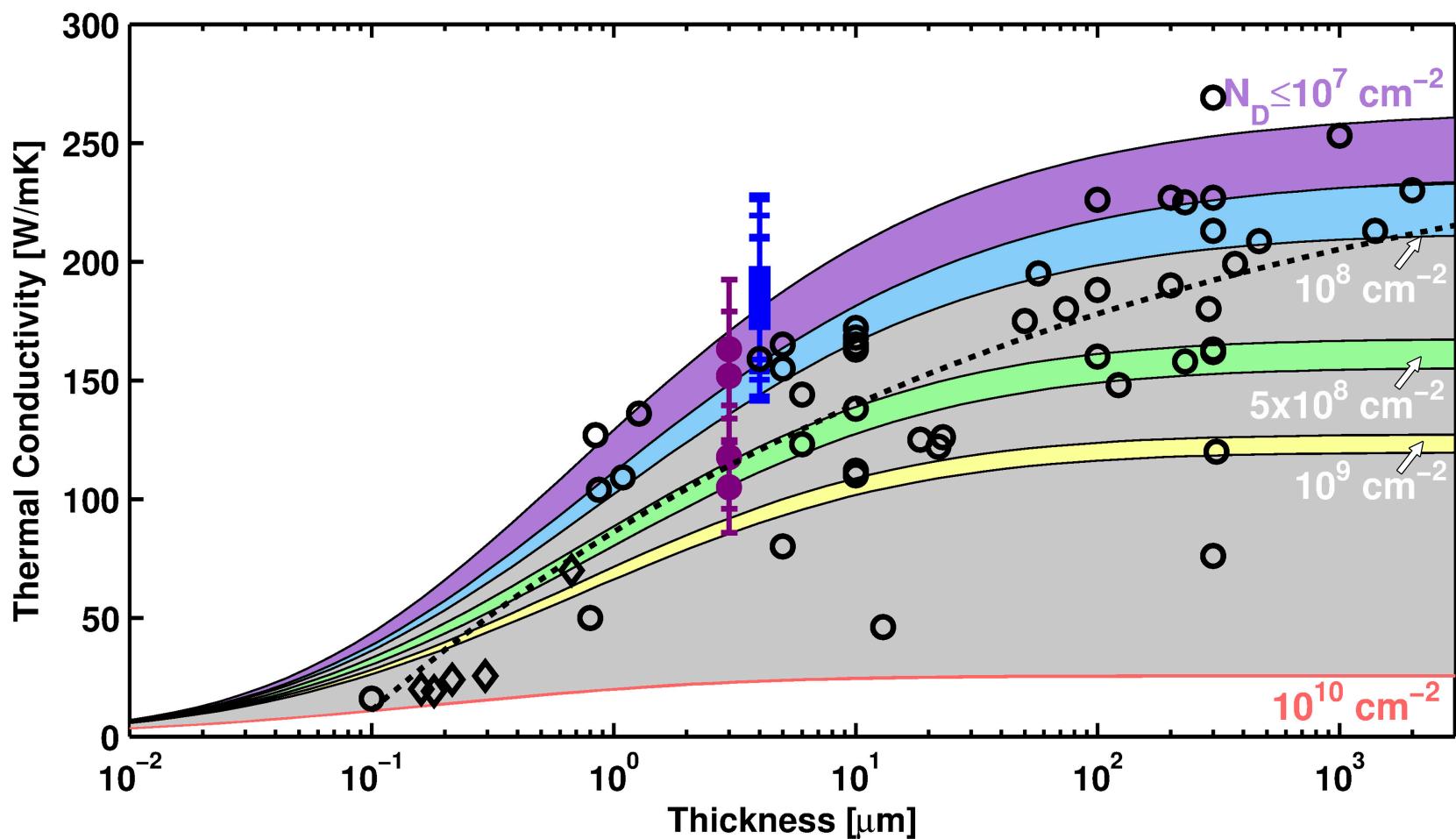
Giri and Hopkins
Nature Materials **19**, 482 (2020)

- Isotope enrichment reduces phonon-defect scattering
- Acoustic phonon bunching and large A-O bandgap reduces phonon-phonon scattering (originally discovered by Lindsay *et al.* *Phys. Rev. Lett.* **111**, 025901 (2013))

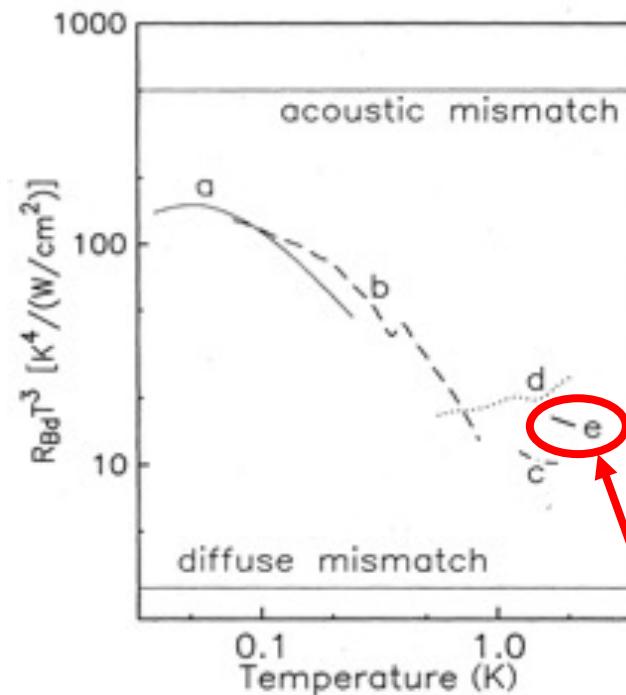
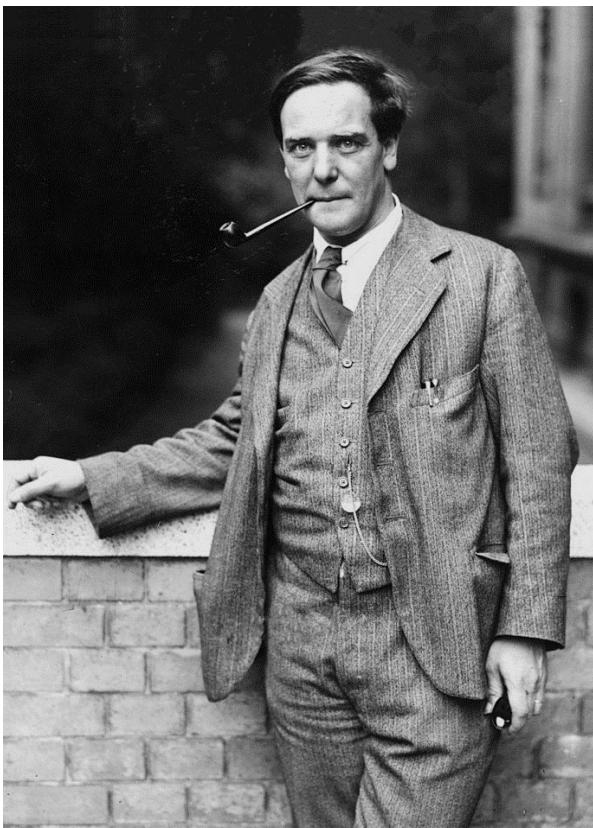


Thermal conductivity of materials: Defects and size effects

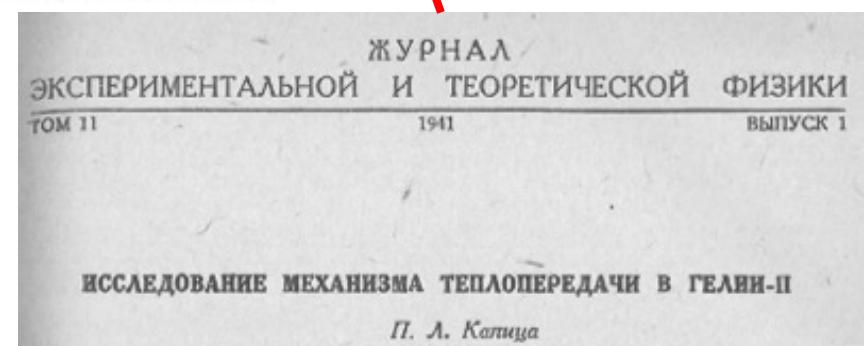
Survey of literature thermal conductivity of GaN *J. Appl. Phys.* **120**, 095104



The founding father of the “Kapitza Resistance”



Plot from Fig. 1 in
Swartz and Pohl,
“Thermal boundary
resistance,”
Rev. Mod. Phys.
61, 605 (1989)

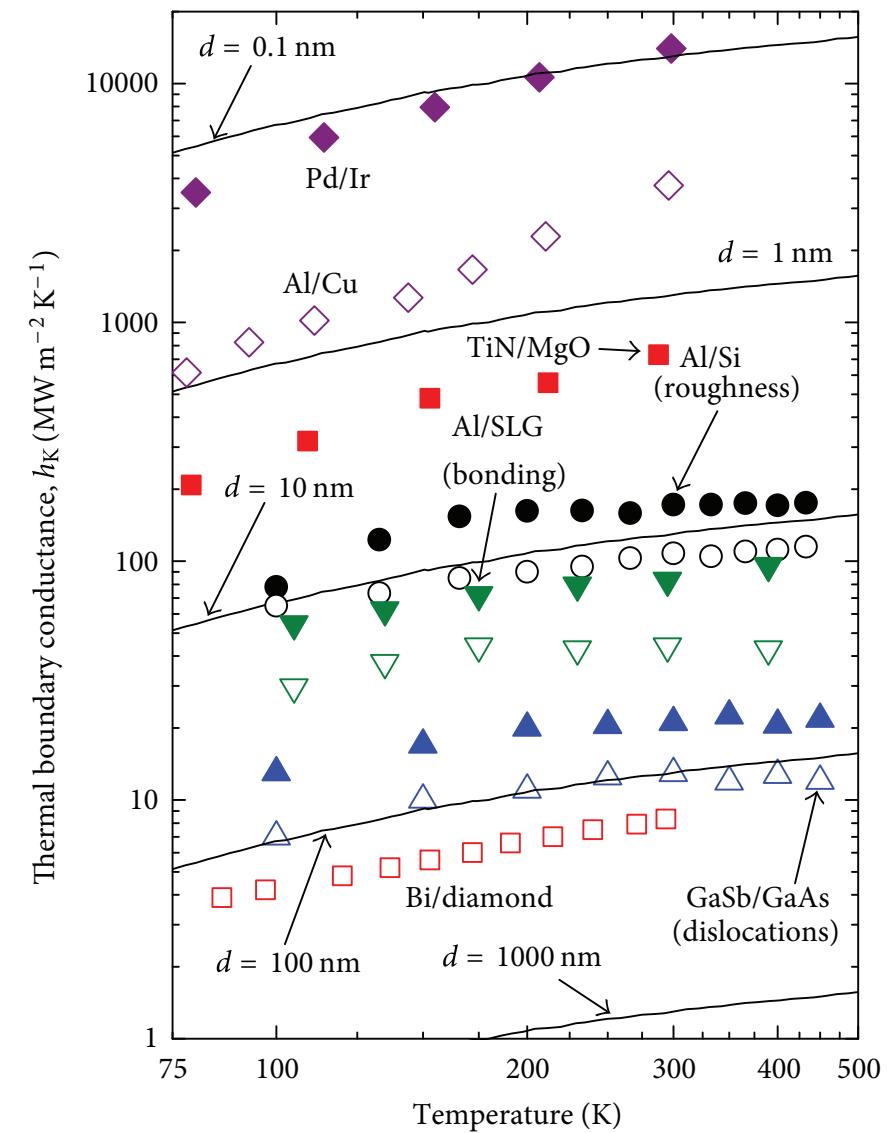
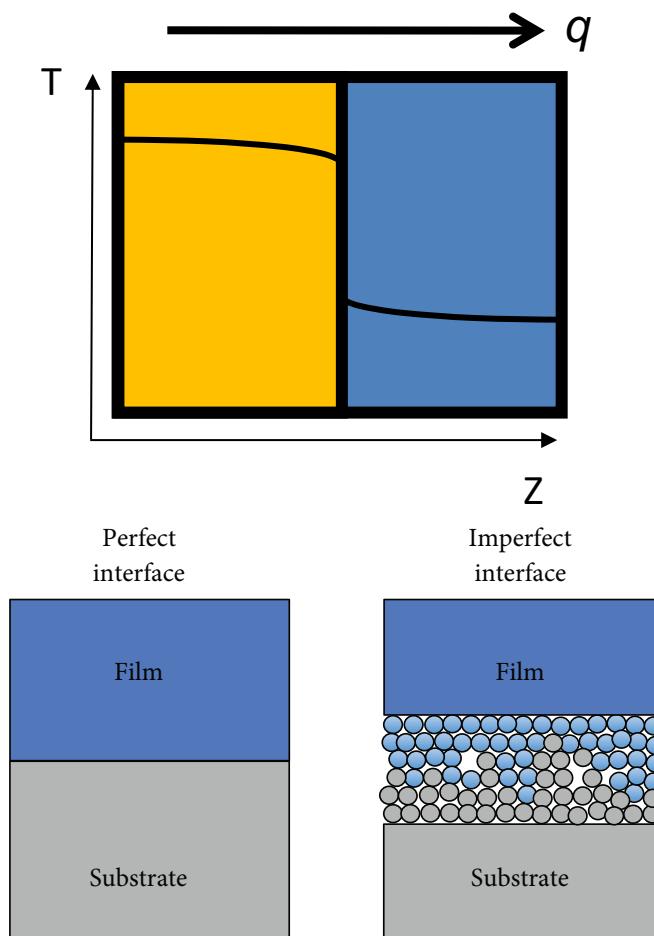


Pyotr Kapitsa

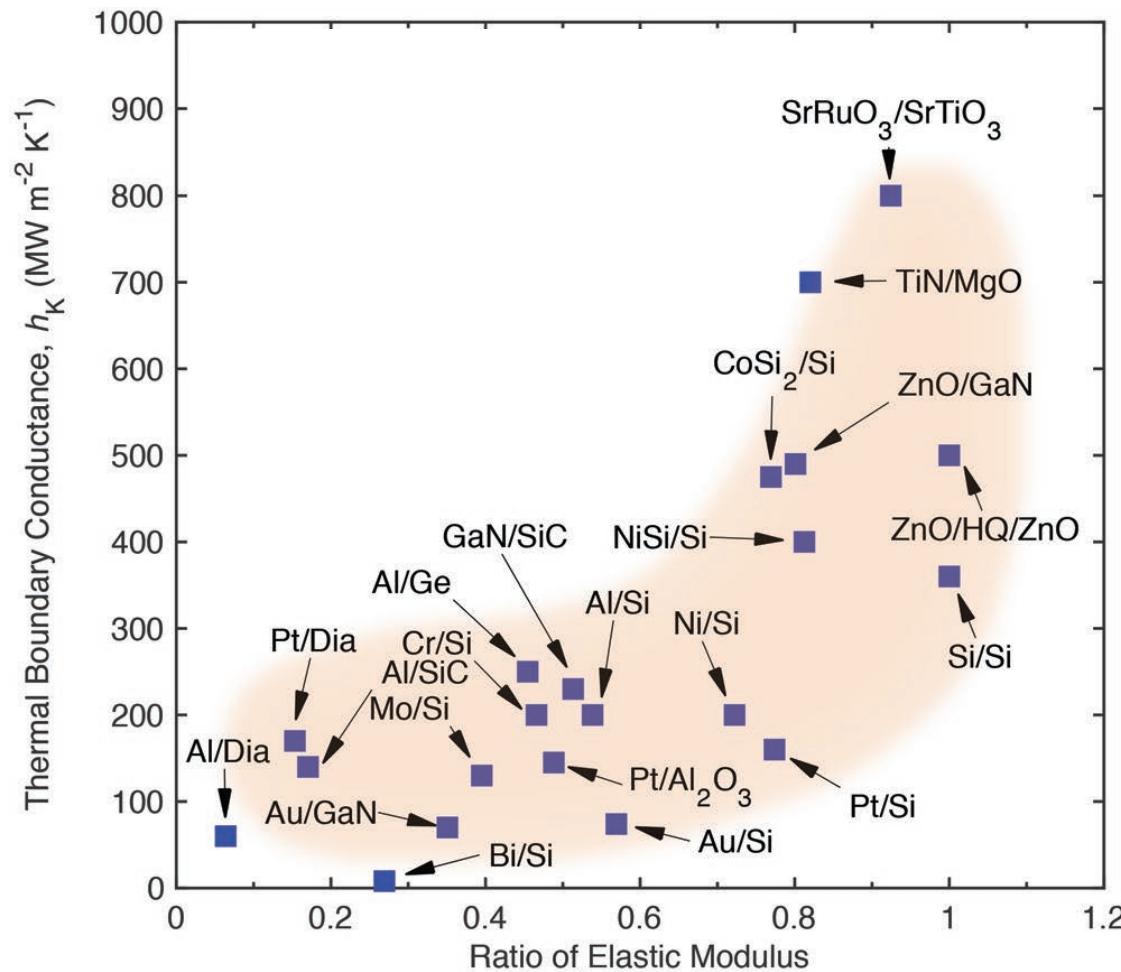
P. L. Kapitsa. The study of heat transfer in helium II. *Zhurnal eksperimentalnoi i teoreticheskoi fiziki*, 11:1–31, 1941.

Thermal boundary conductance – nanoscale resistances

$$q = h_K \Delta T = \frac{1}{R_K} \Delta T$$

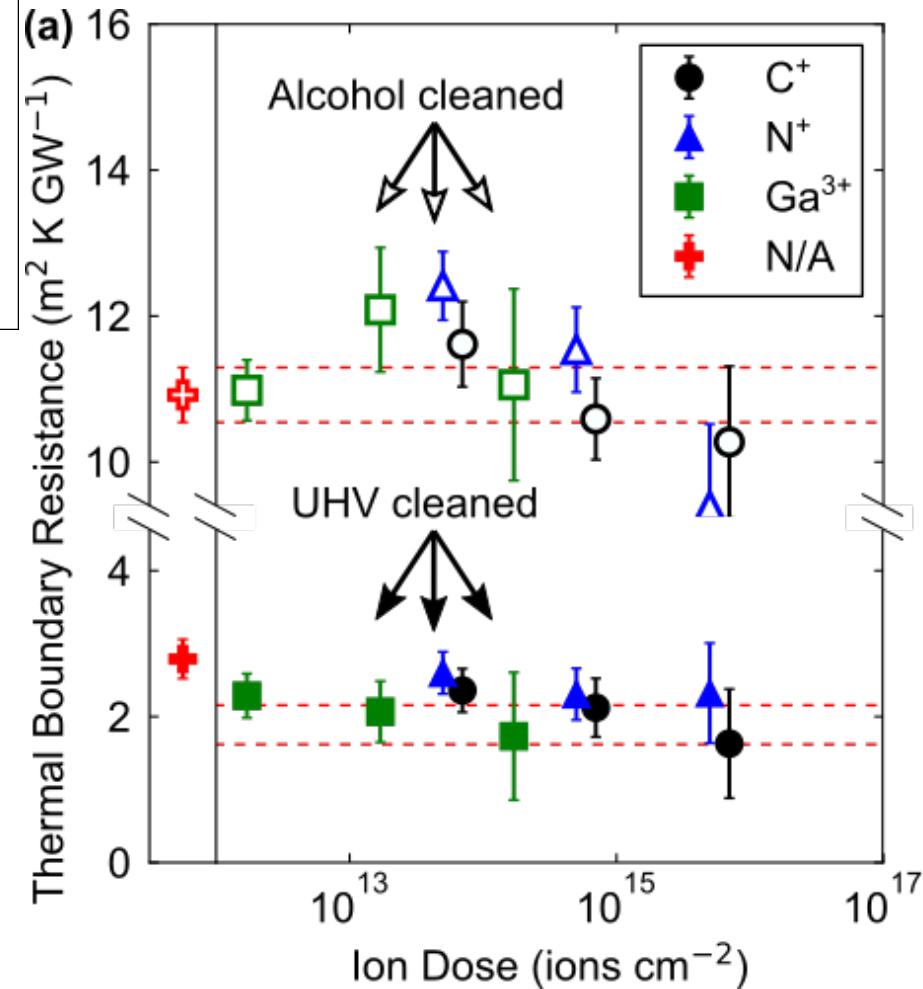
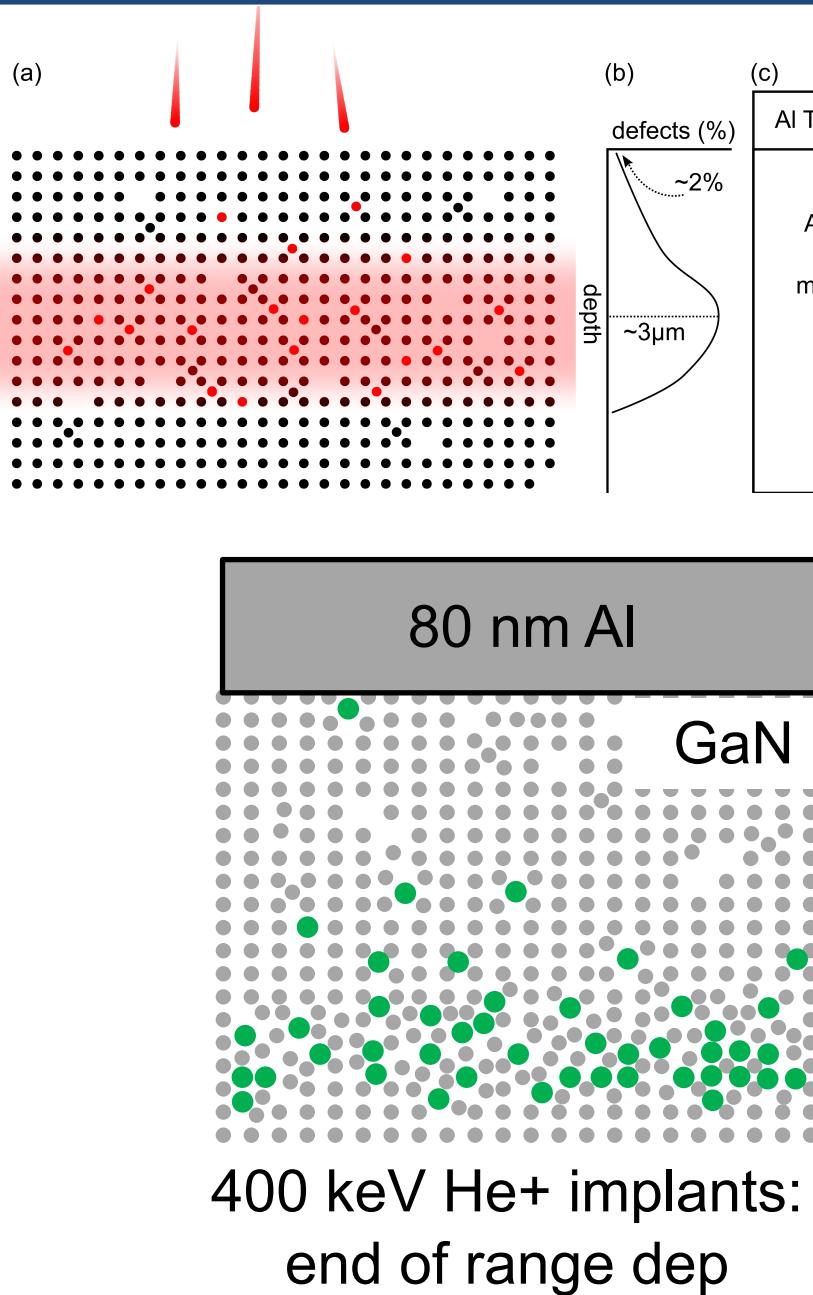


Phonon dominated TBC: metal/nonmetal interfaces

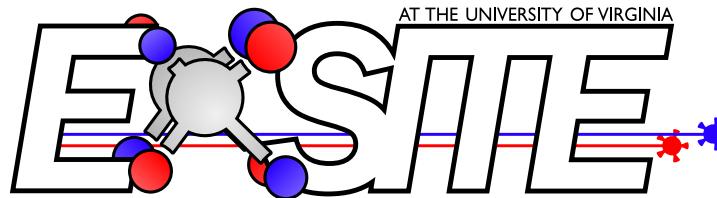


Recent Review
*Advanced Functional
Materials* **30**, 1903857 (2020)

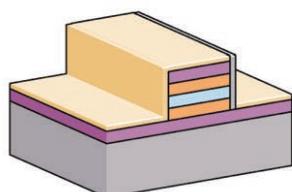
Phonon dominated TBC: “Engineering defects”



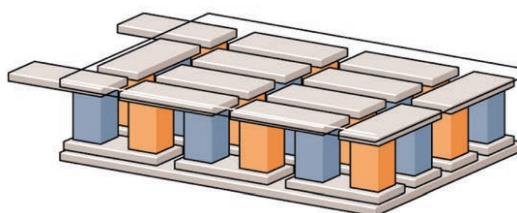
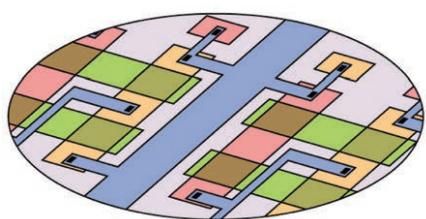
Thermal properties of nanomaterials and interfaces



Quantum cascade lasers

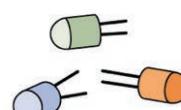


Computing and logic circuits

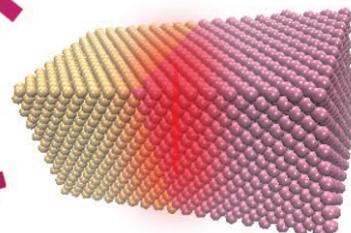


Thermoelectrics

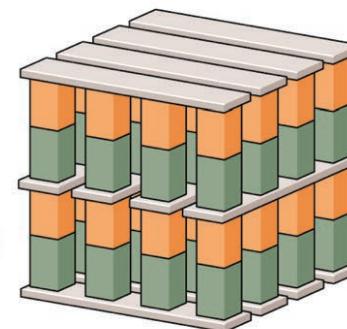
Light emitting diodes



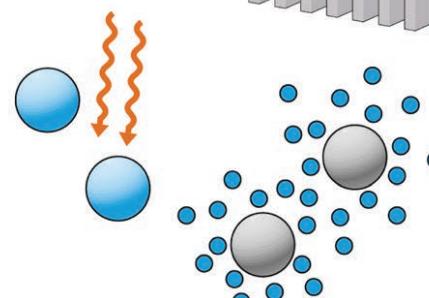
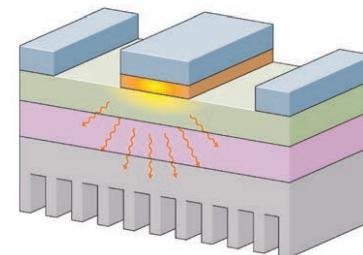
Nanoscale Interfaces



Phase change memory



High power electronic devices



Nanoparticles for Photothermal therapy

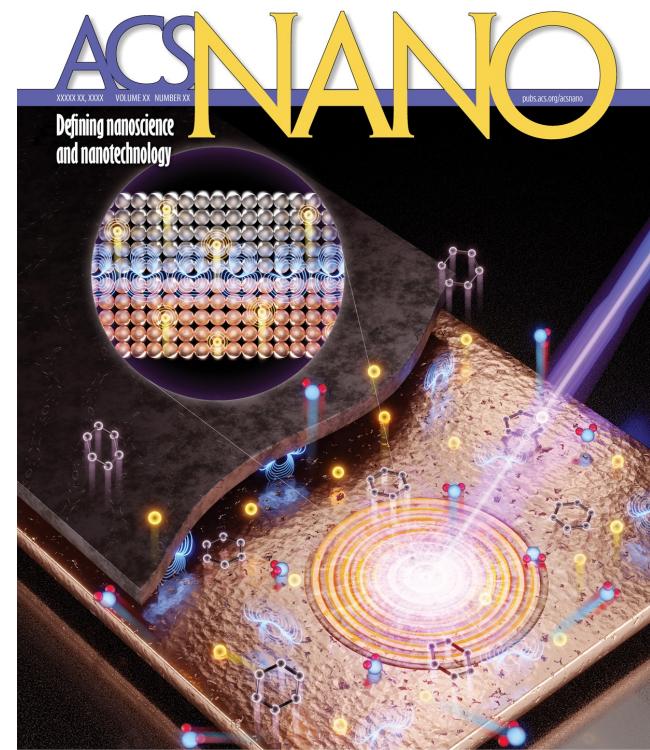
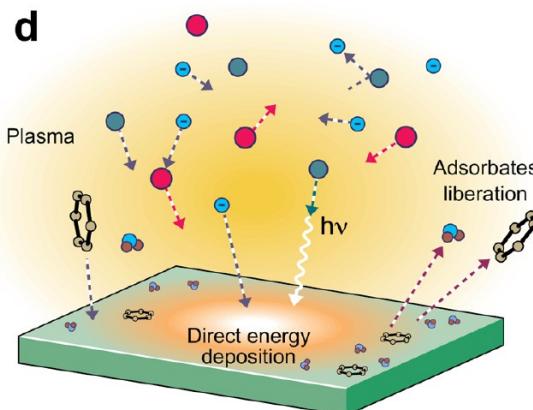
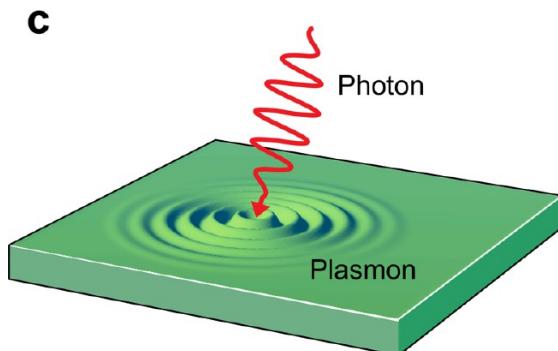
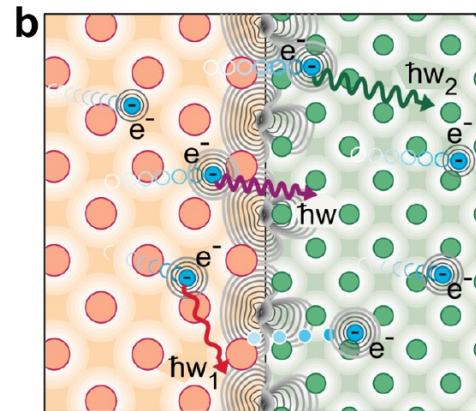
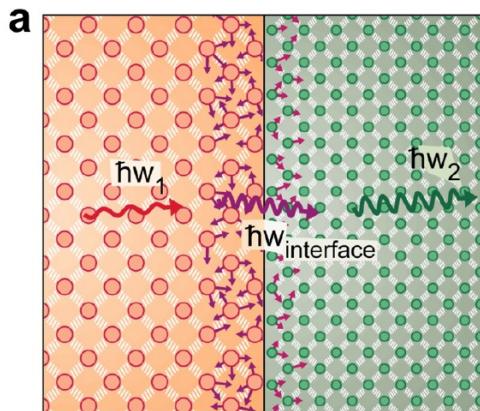
Outline

Theme: Coupling of carriers across interfaces lead to unique energy states and modes of heat transfer

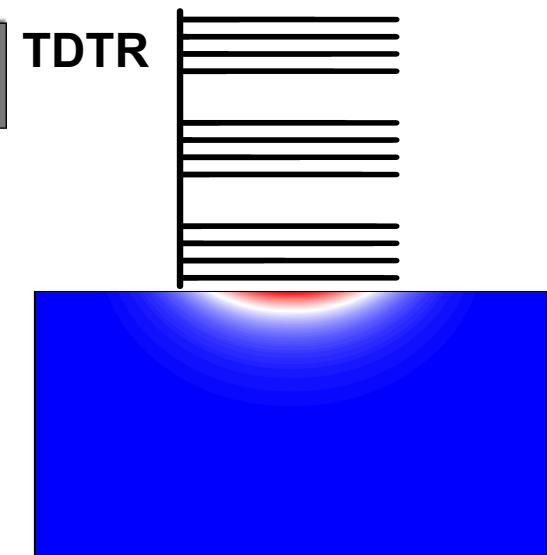
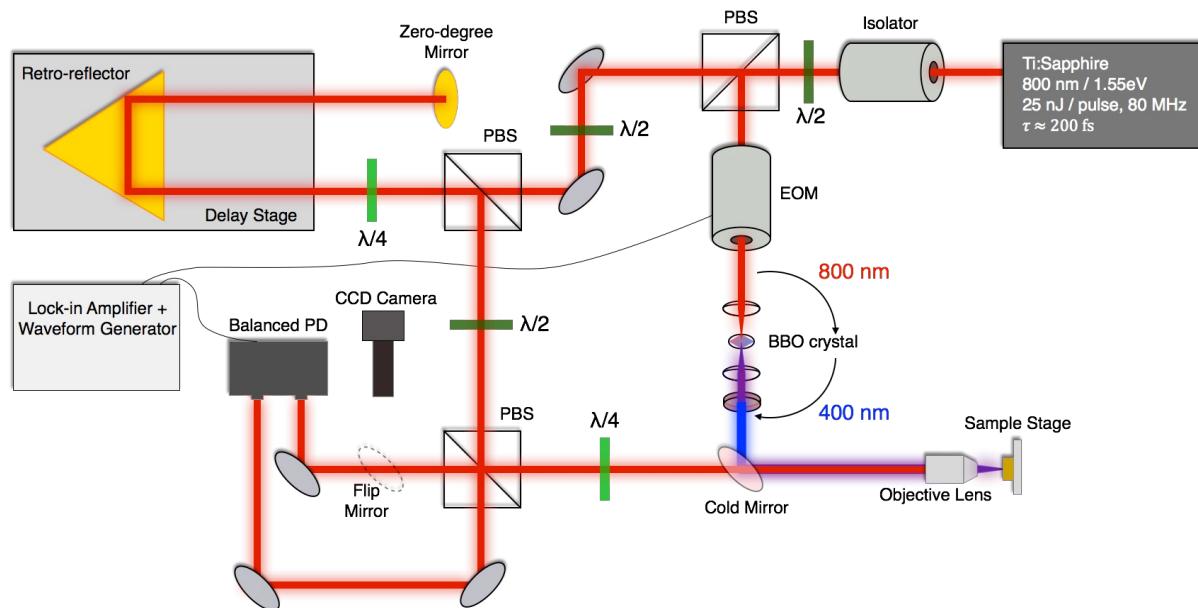
- How do we measure the thermal conductivity of materials and the thermal boundary resistance at their interfaces?
- Thermal conductivity of superlattices
 - Minimum thermal conductivity and creating “crystals of interfaces”
- Interfacial heat transfer control of the IR properties of solids
 - Near field radiative interfacial heat transfer with plasmon-polaritons and phonon-polaritons
- Transient temperature changes during plasma-surface interactions
 - “Plasma cooling”

Ultrafast and Nanoscale Energy Transduction Mechanisms and Coupled Thermal Transport across Interfaces

Ashutosh Giri,* Scott G. Walton, John Tomko, Niraj Bhatt, Michael J. Johnson, David R. Boris, Guanyu Lu, Joshua D. Caldwell, Oleg V. Prezhdo, and Patrick E. Hopkins*

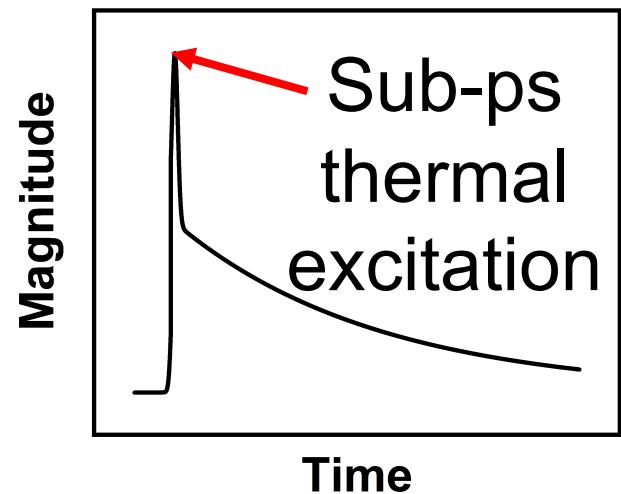


How do we measure the thermal properties? TDTR

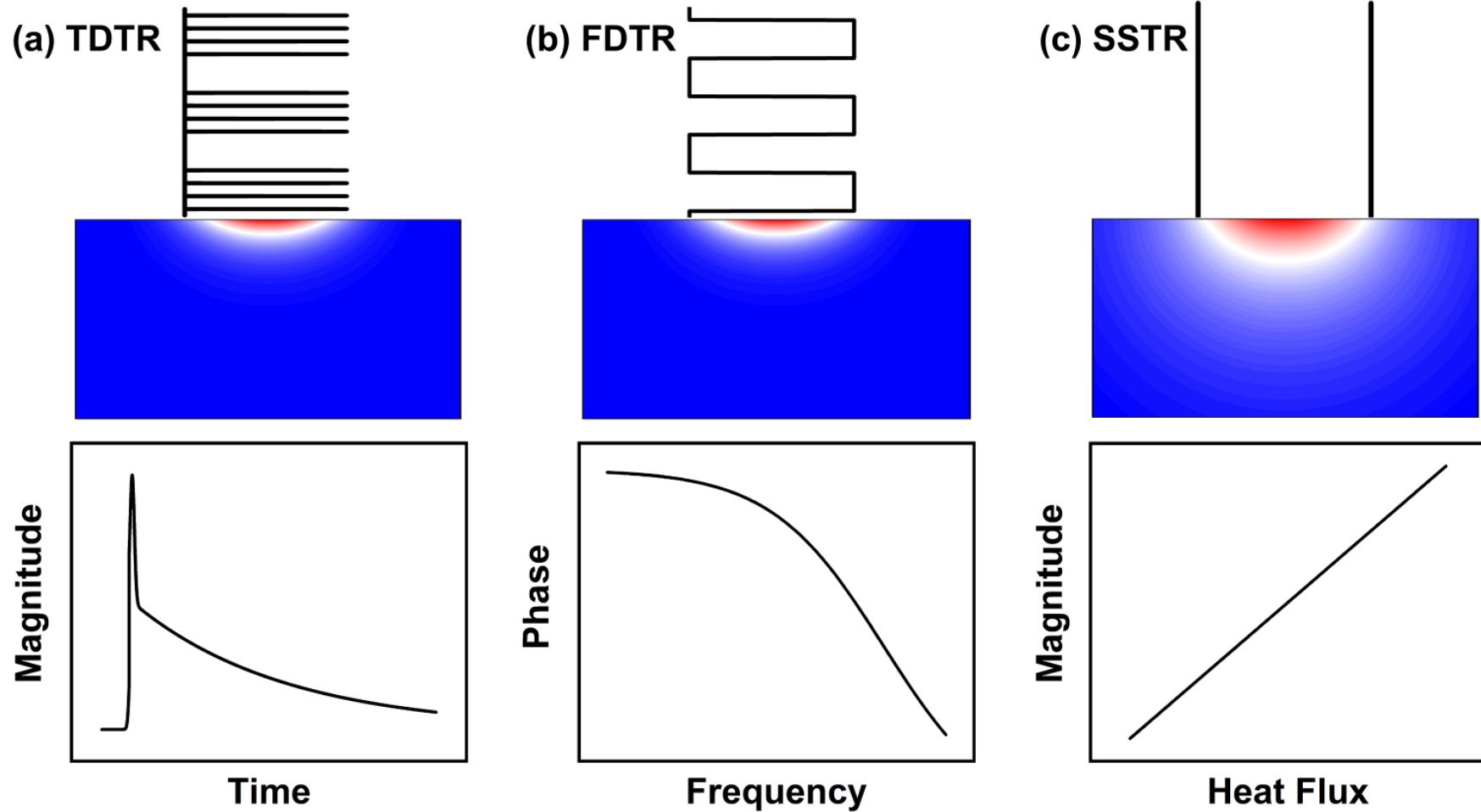


TDTR Reviews and Analyses

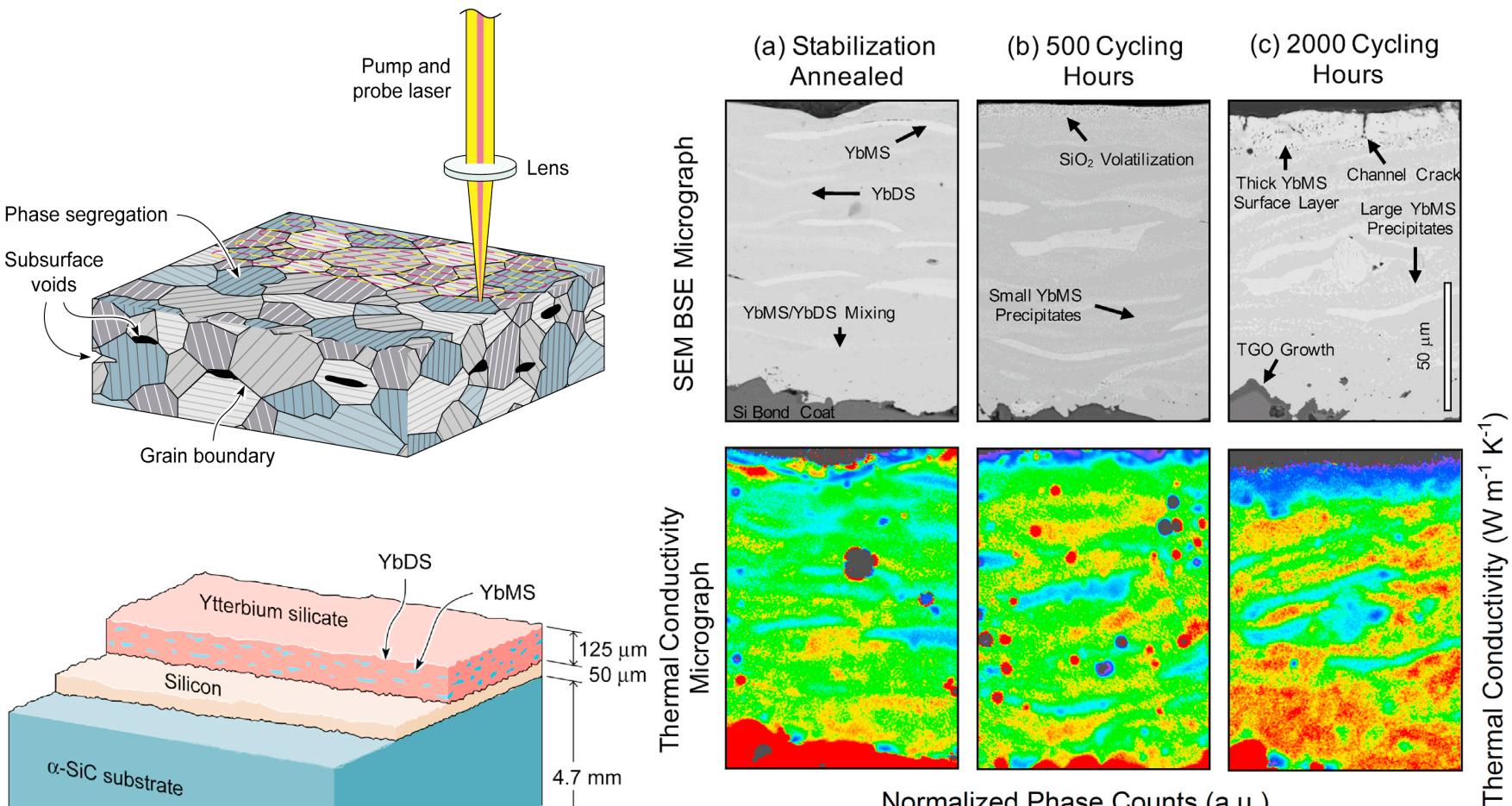
Rev. Sci. Instr. **75**, 5119;
Rev. Sci. Instr. **79**, 114902
J. Heat Trans. **132**, 081302;
Ann. Rev. Heat Trans. **16**, 159



The “flavors” of thermoreflectance

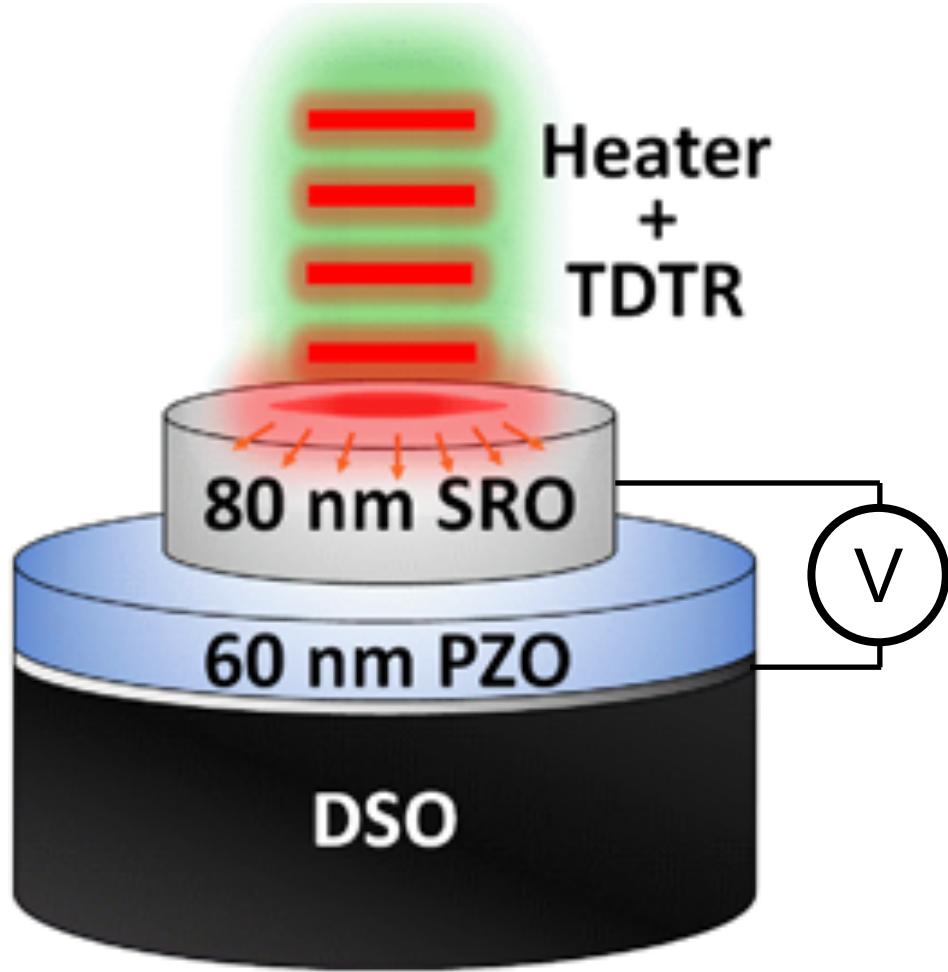


Local thermal conductivity measurements and spatial maps

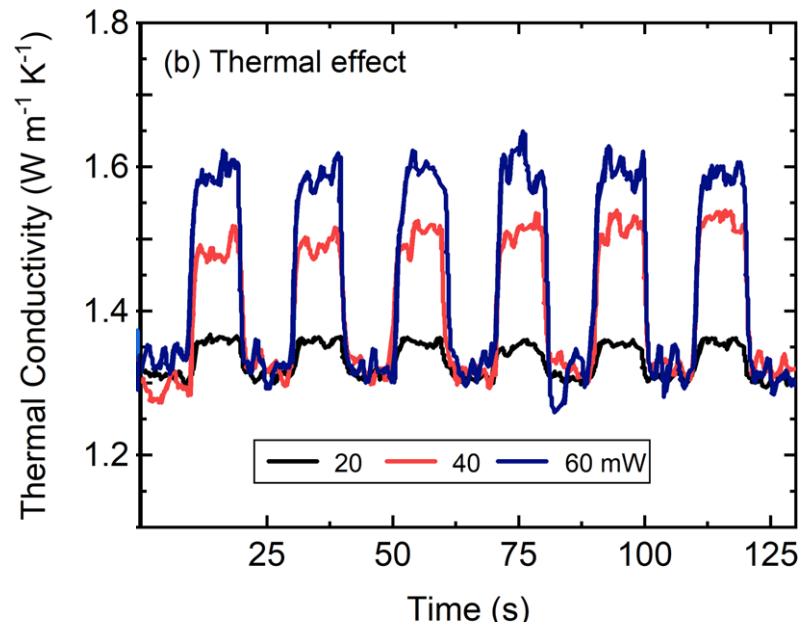
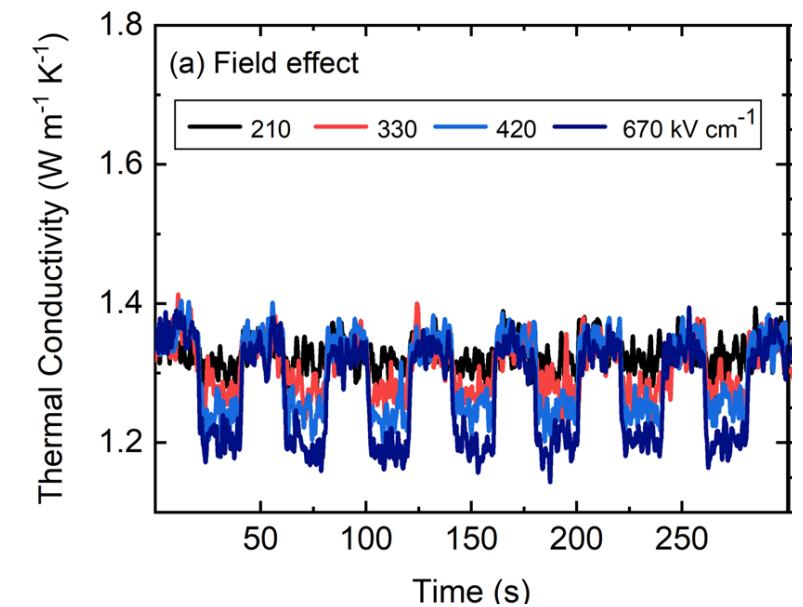


Olson *et al.*, *Materials Today Phys.*
17, 100304 (2021)

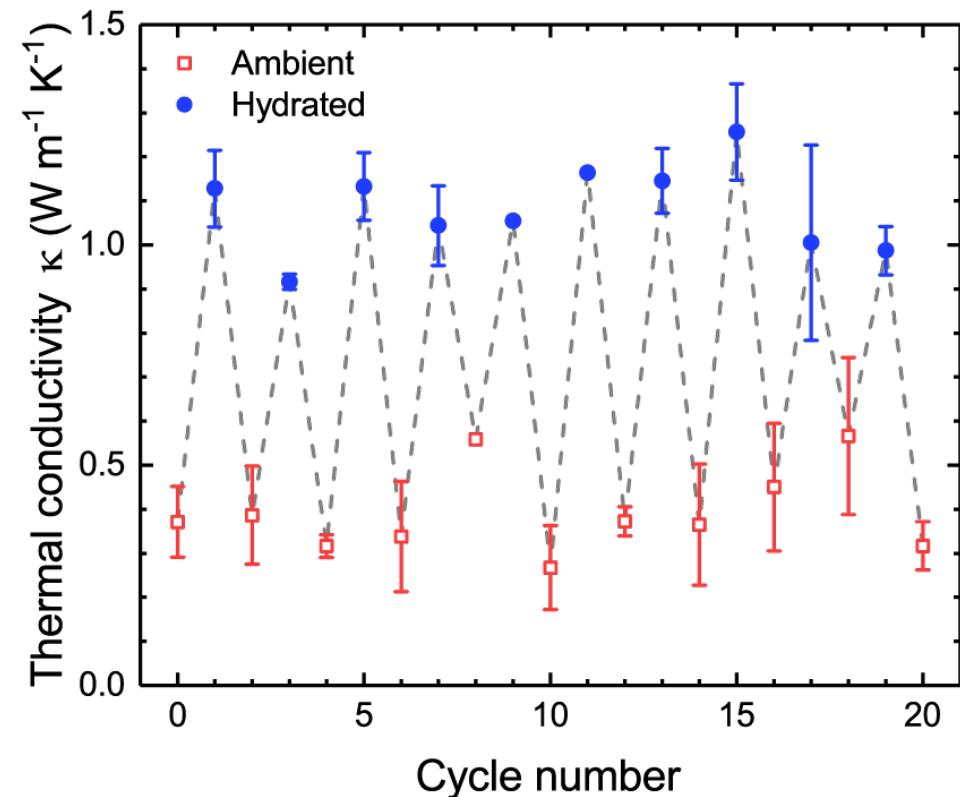
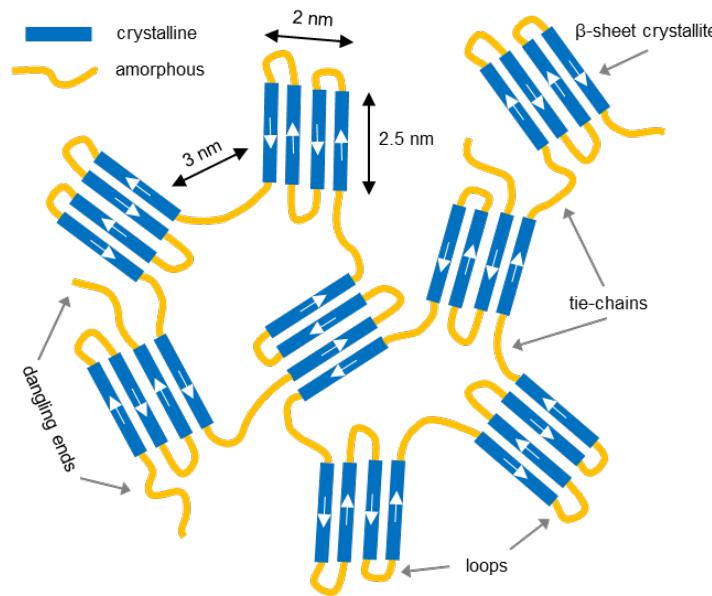
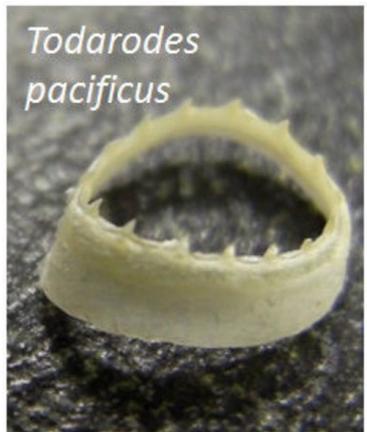
Actively switching thermal conductivity w/ FE/AFE domains



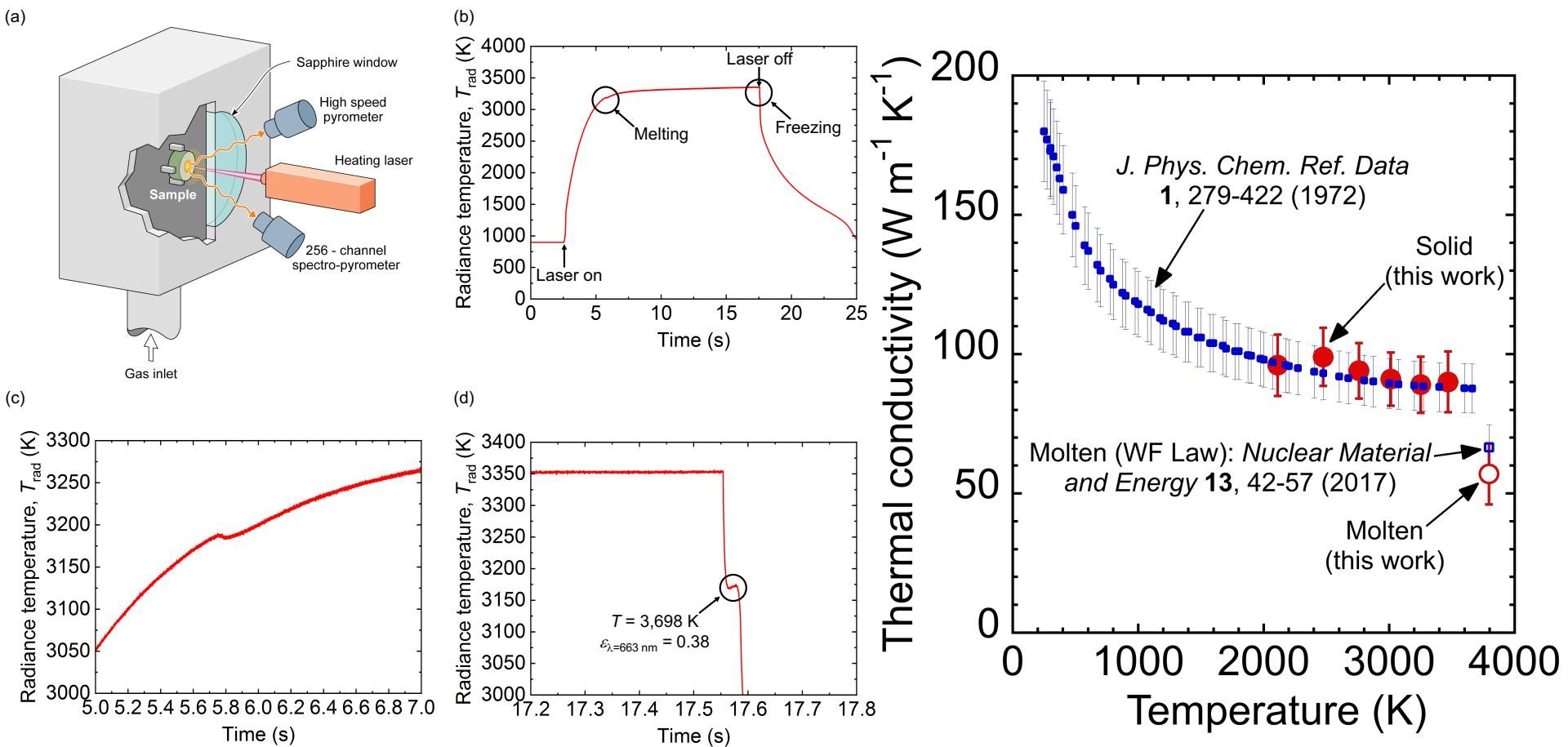
Aryana *et al.*, *Nat. Comm.*
13, 1573 (2022)



Actively switching thermal conductivity in bio materials

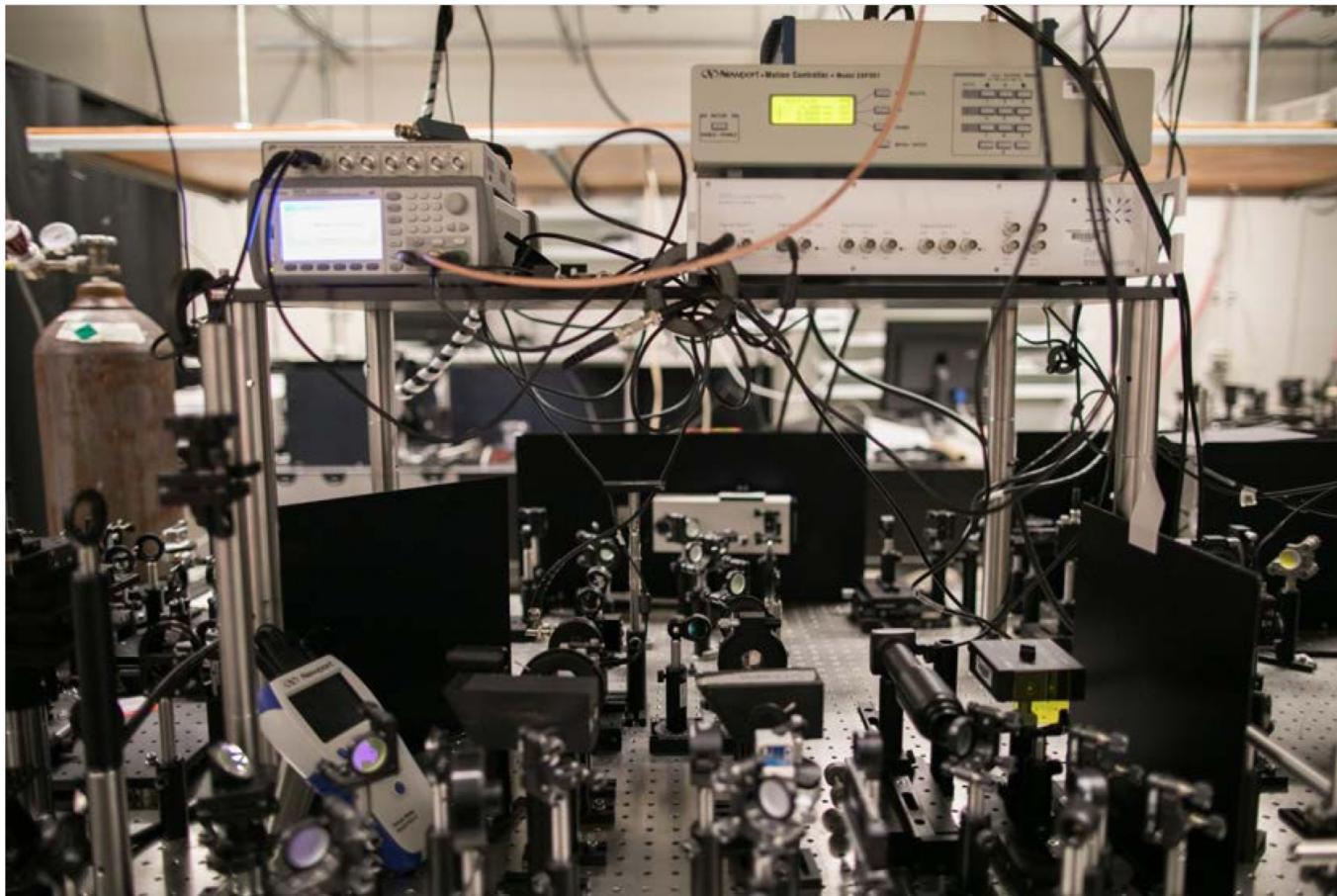


Thermometry at ultrahigh temperatures



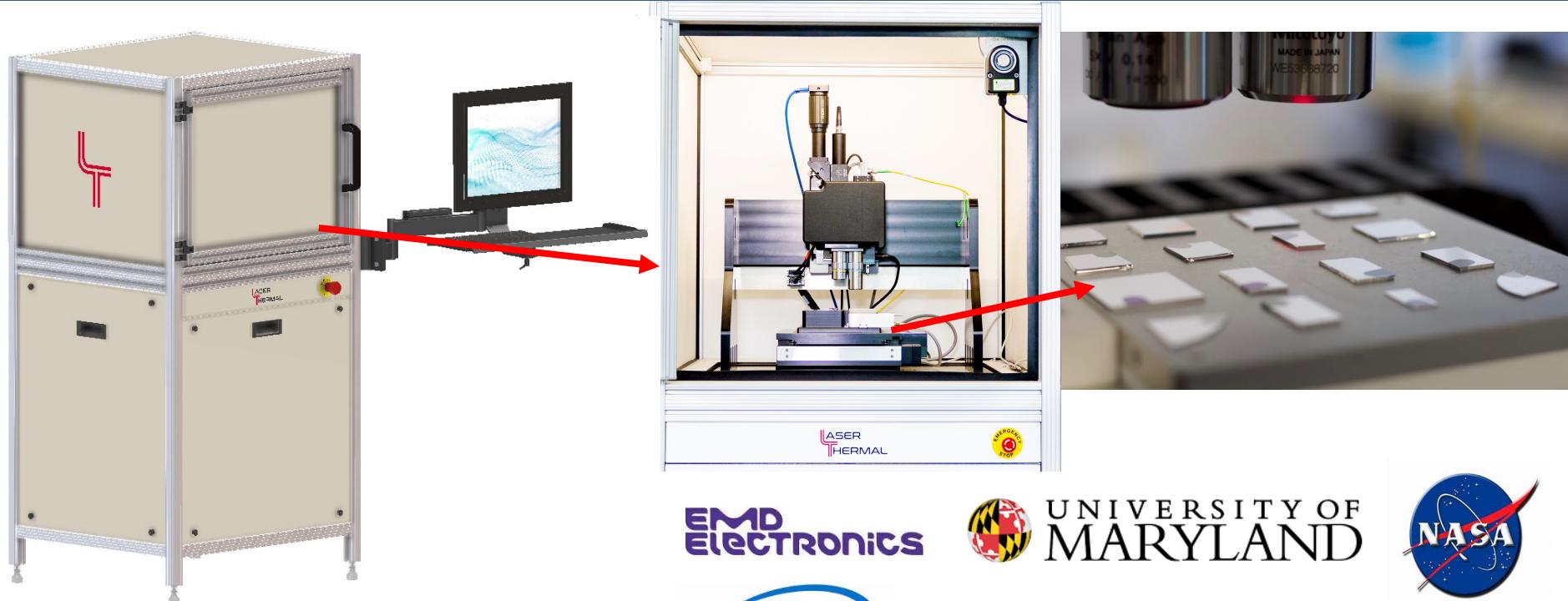
So why don't we all have thermoreflectance systems????

The typical thermoreflectance set up



A LOT of optics, upkeep and expertise for analysis

SSTR-F: Commercialized for turn-key thermal conductivity microscope for bulk materials, thin films and interfaces



LASER
Thermal

EMD
ELECTRONICS

UNIVERSITY OF
MARYLAND



intel

NIST NATIONAL INSTITUTE OF
STANDARDS AND TECHNOLOGY
U.S. DEPARTMENT OF COMMERCE



U.S. NAVAL
RESEARCH
LABORATORY

OAK RIDGE
National Laboratory

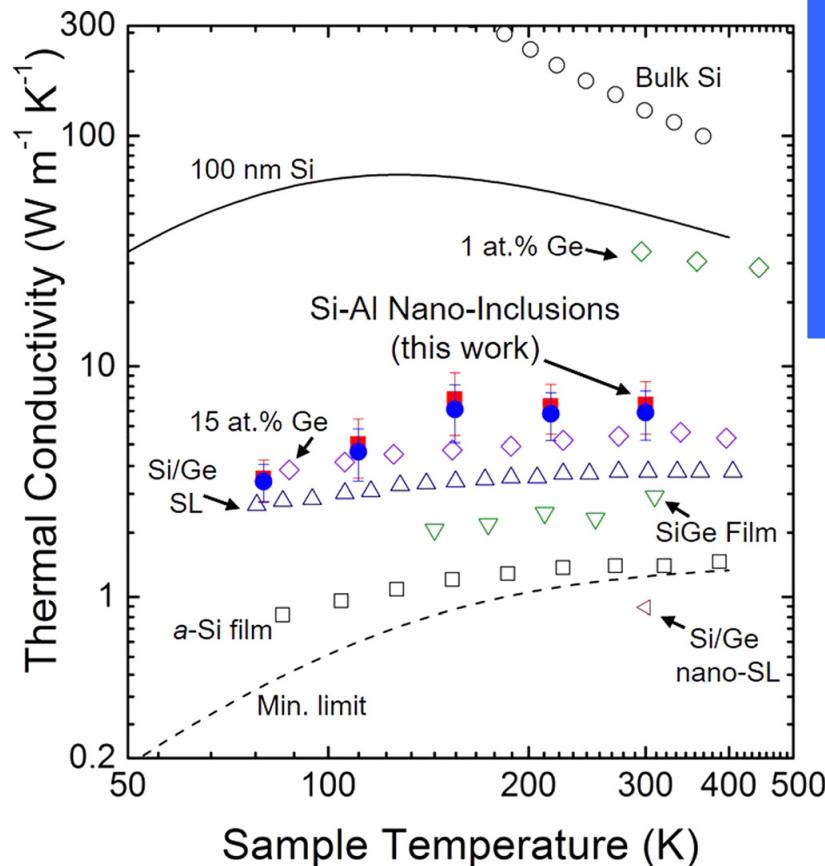
Disclosure: Hopkins co-Founder

Outline

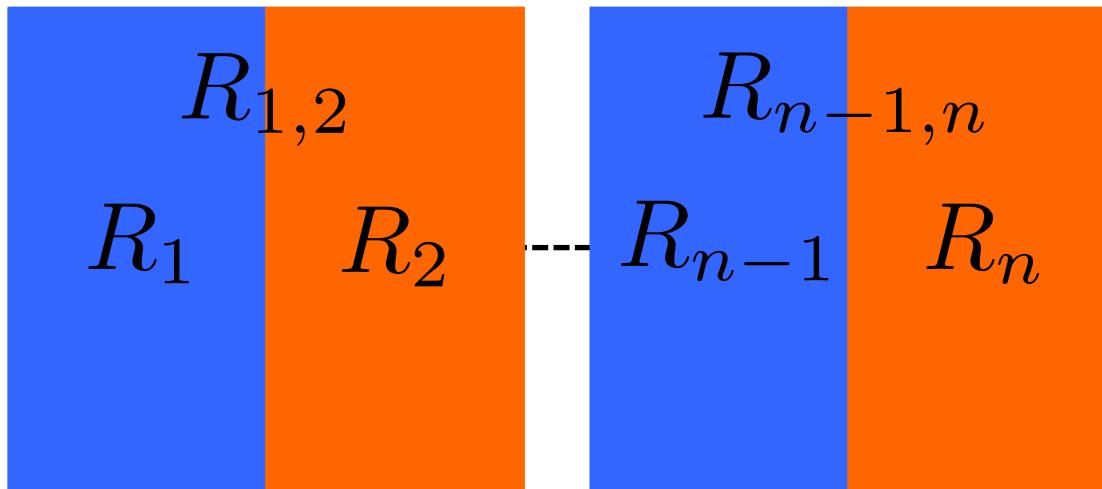
Theme: Coupling of carriers across interfaces lead to unique energy states and modes of heat transfer

- How do we measure the thermal conductivity of materials and the thermal boundary resistance at their interfaces?
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 - **Minimum thermal conductivity and creating “crystals of interfaces”**
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Phonon thermal conductivity in superlattices: incoherent



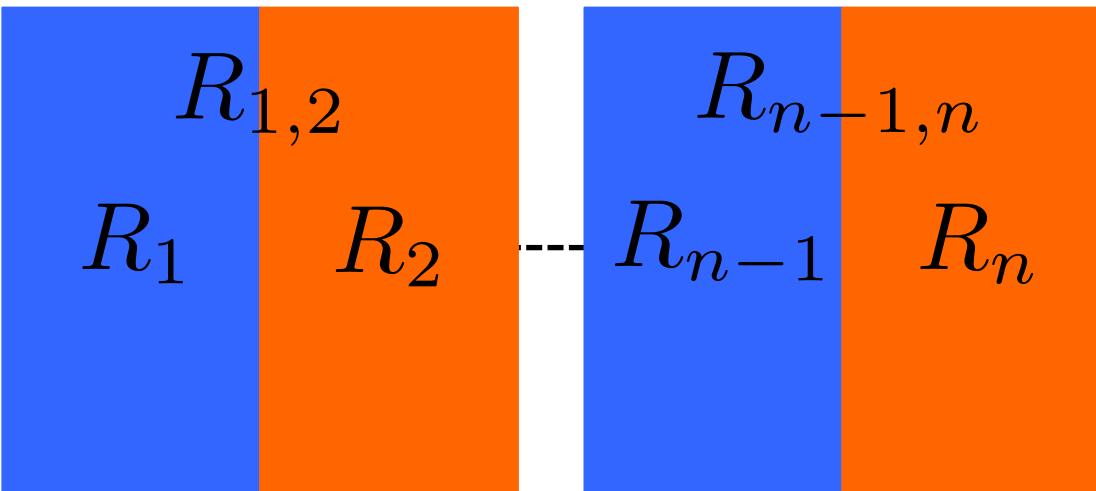
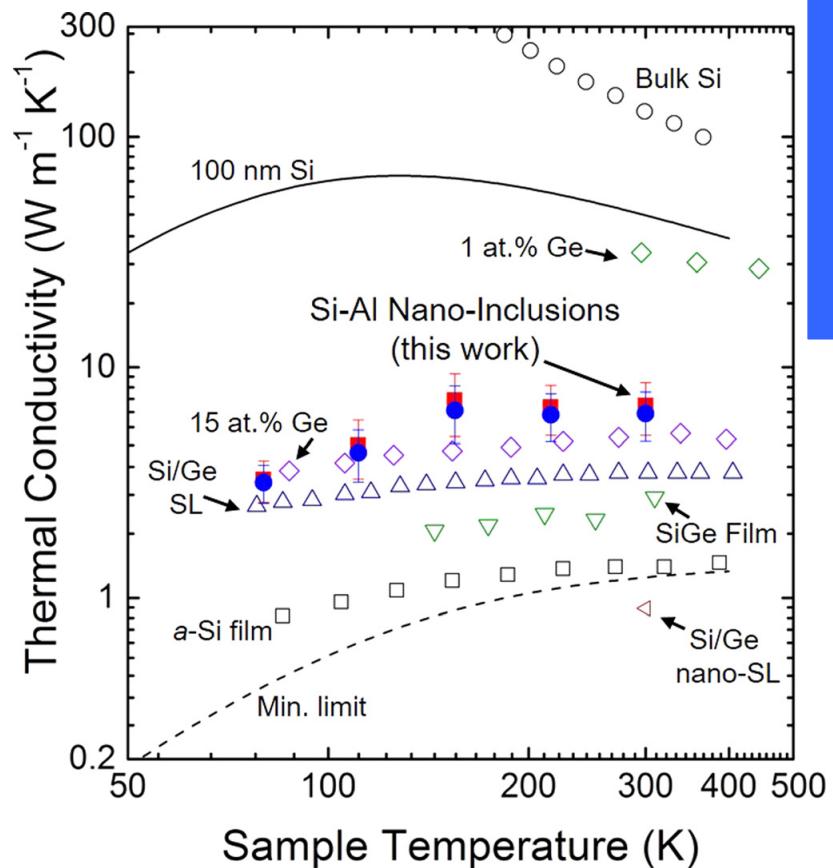
APL 112, 213103



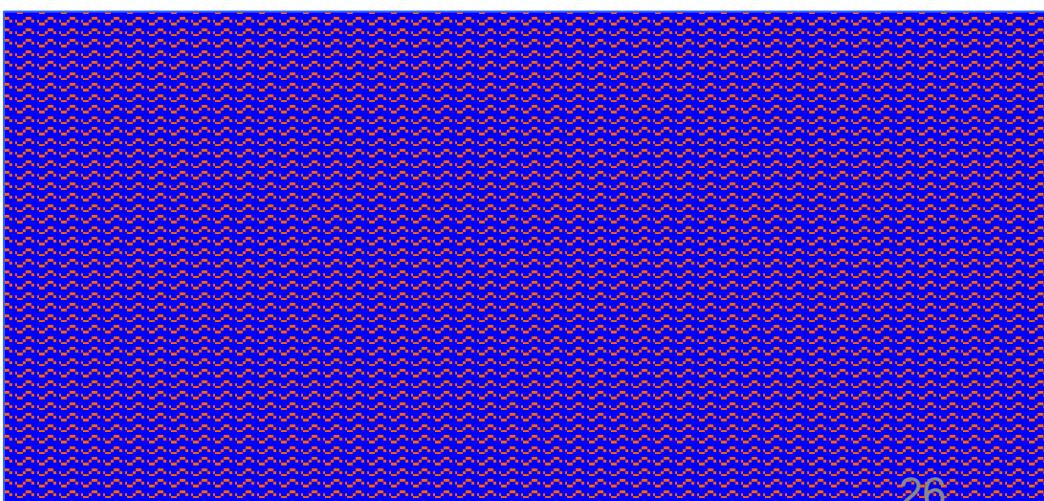
$$1/\kappa \propto 1/h = \sum_n R_j$$

Phonons scatter at every interface, and thus each interface offers a resistance to the overall thermal conductivity

Phonon thermal conductivity in superlattices: coherent



What if layers are “linked”?
– coherent transport



The minimum thermal conductivity of superlattices

PHYSICAL REVIEW B

VOLUME 25, NUMBER 6

15 MARCH 1982

VOLUME 84, NUMBER 5

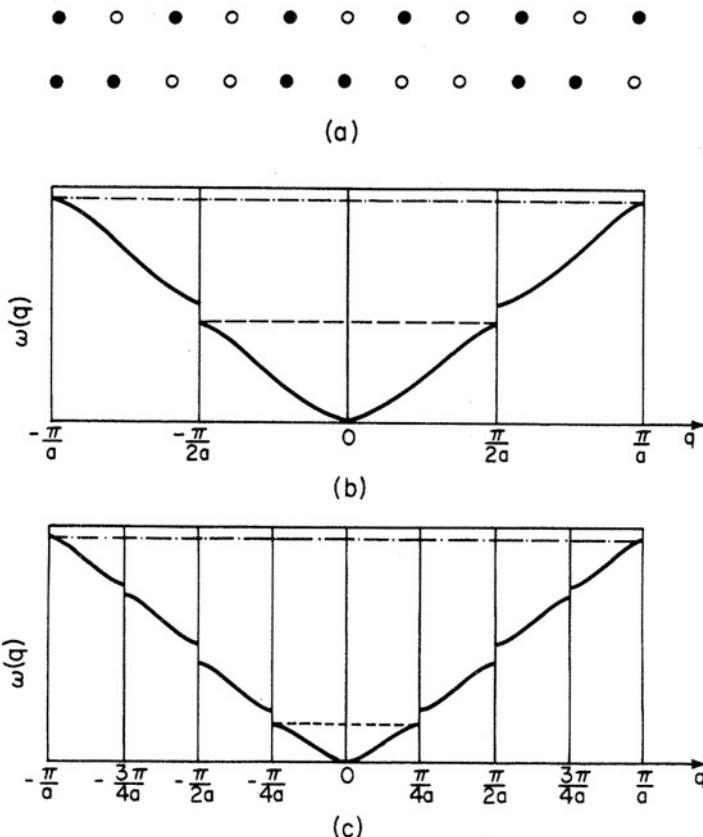
PHYSICAL REVIEW LETTERS

31 JANUARY 2000

Thermal conductivity of superlattices

Shang Yuan Ren* and John D. Dow

Department of Physics and Coordinated Science Laboratory, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801
(Received 21 September 1981)

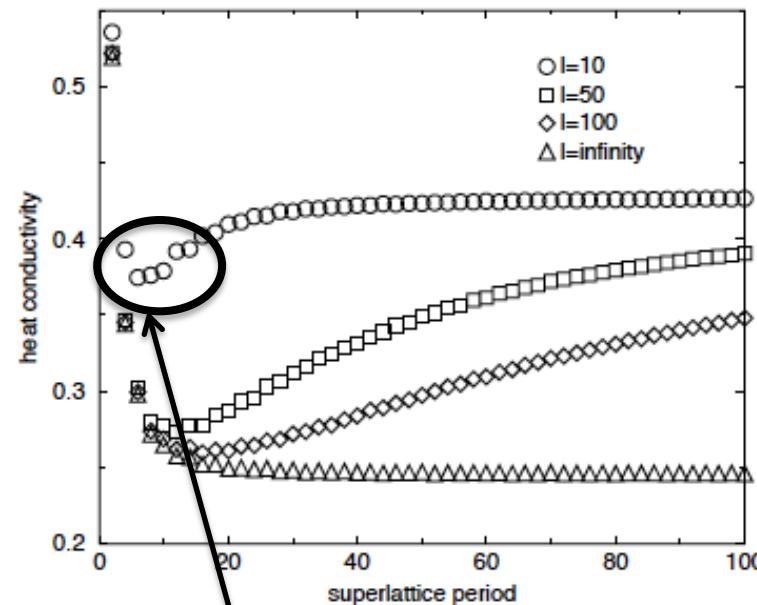


Interfacial periodicity can lead to “mini-band” formation

Minimum Thermal Conductivity of Superlattices

M. V. Simkin and G. D. Mahan

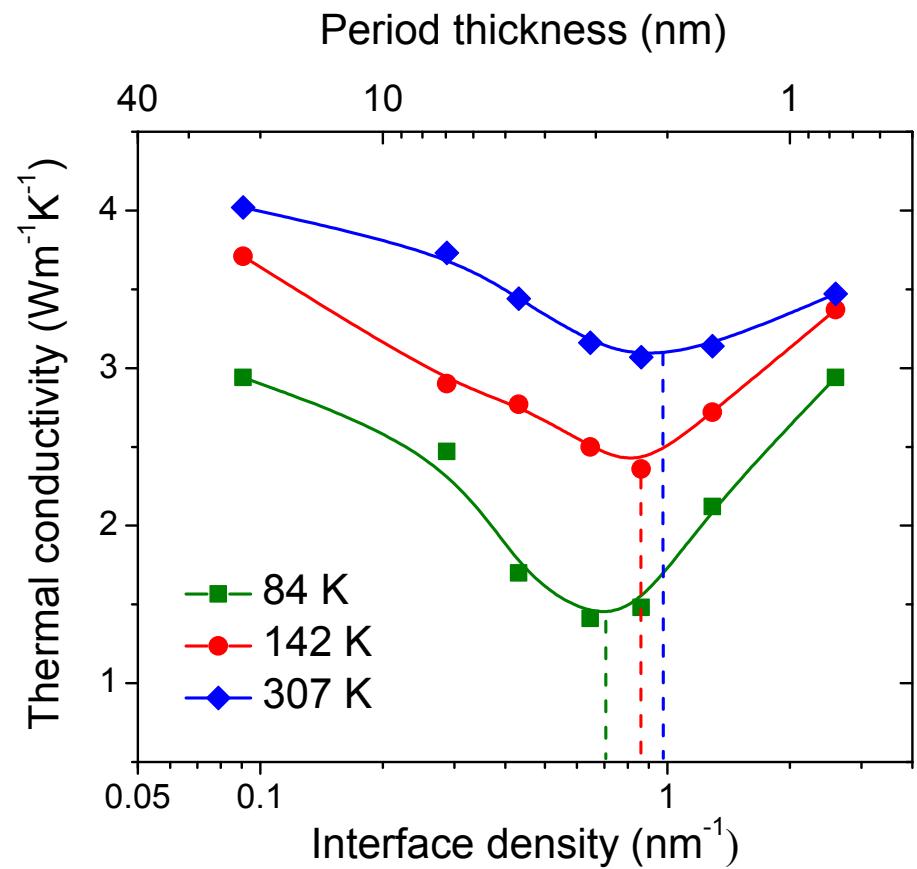
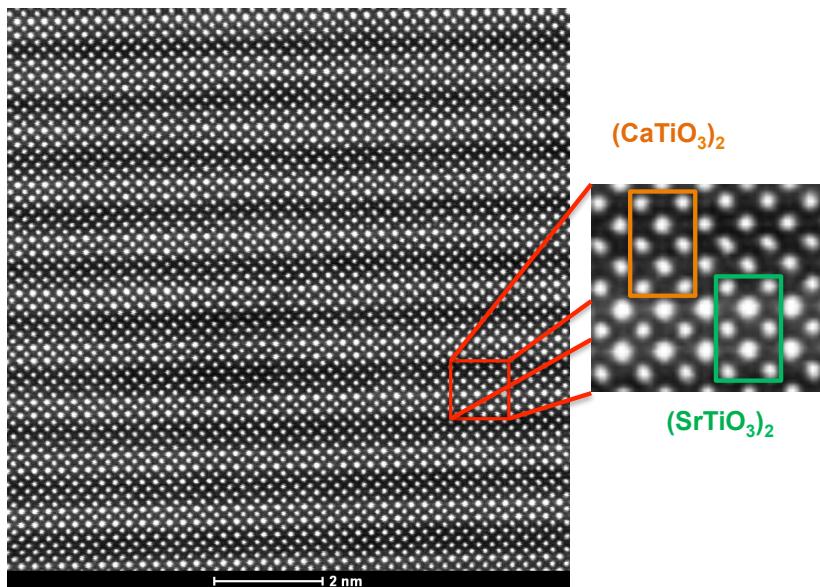
Department of Physics and Astronomy, University of Tennessee, Knoxville, Tennessee 37996-1200
and Solid State Division, Oak Ridge National Laboratory, P.O. Box 2008, Oak Ridge, Tennessee 37831
(Received 23 July 1999)



Mini-band formation leads to a minimum in the superlattice thermal conductivity

Experimental evidence of minimum thermal conductivity

- More pronounced minimum at low T, thermal conductivity measurements show trends of mini-band formation

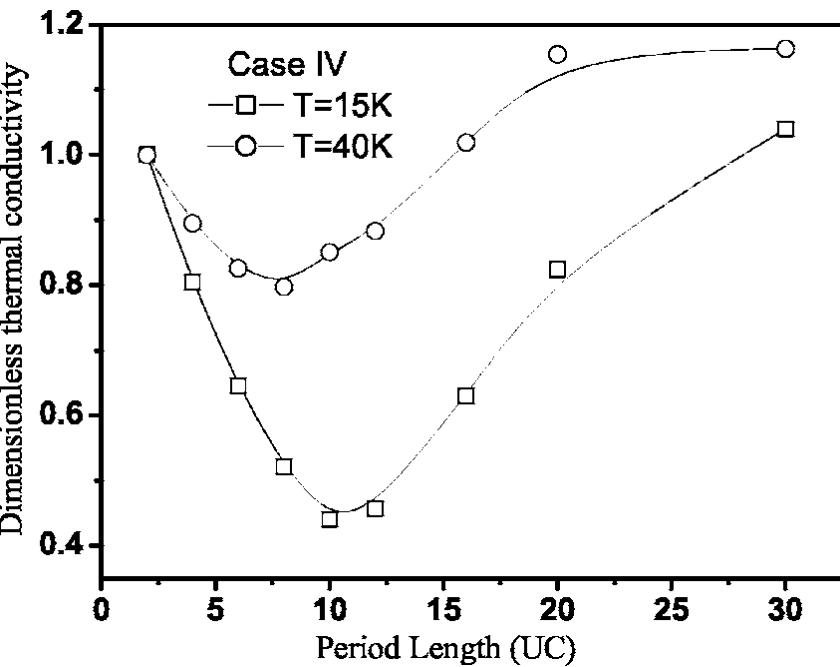


SL design to manipulate coherent phonon transport

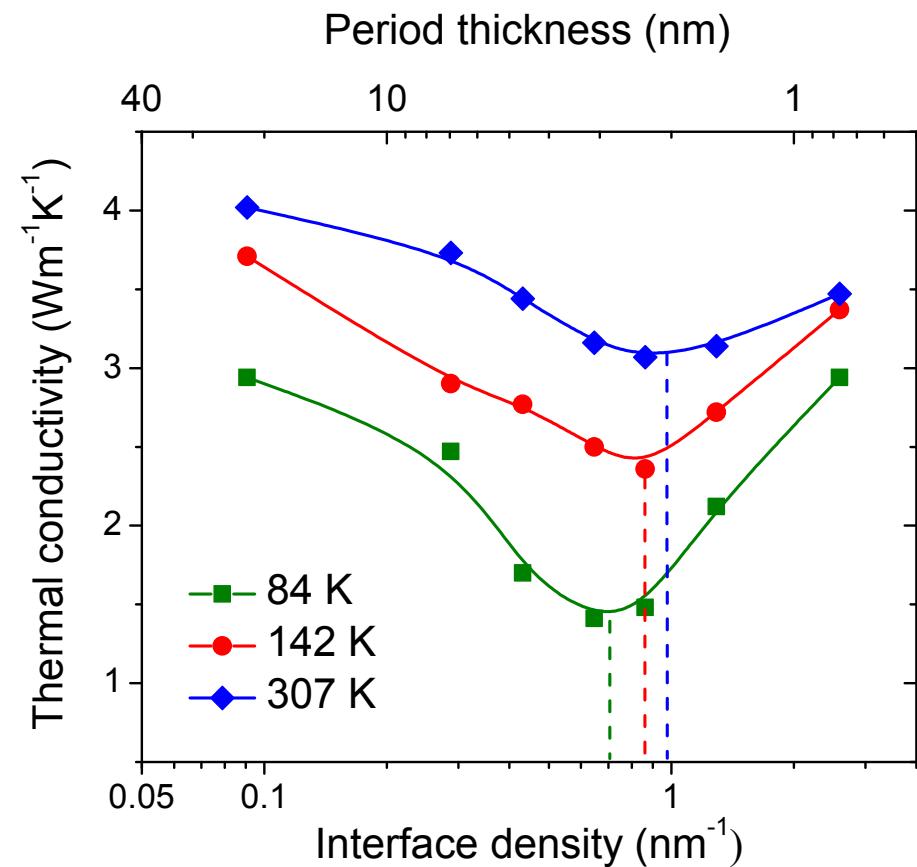
Ravichandran *et al.* *Nature Materials* 13, 168 (2013)

Experimental evidence of minimum thermal conductivity

- More pronounced minimum at low T, thermal conductivity measurements show trends of mini-band formation
- MD simulation (left), mini-band = phonon bandgap (PRB **72**, 174302)



Phys. Rev. B **72**, 174302



SL design to manipulate coherent phonon transport

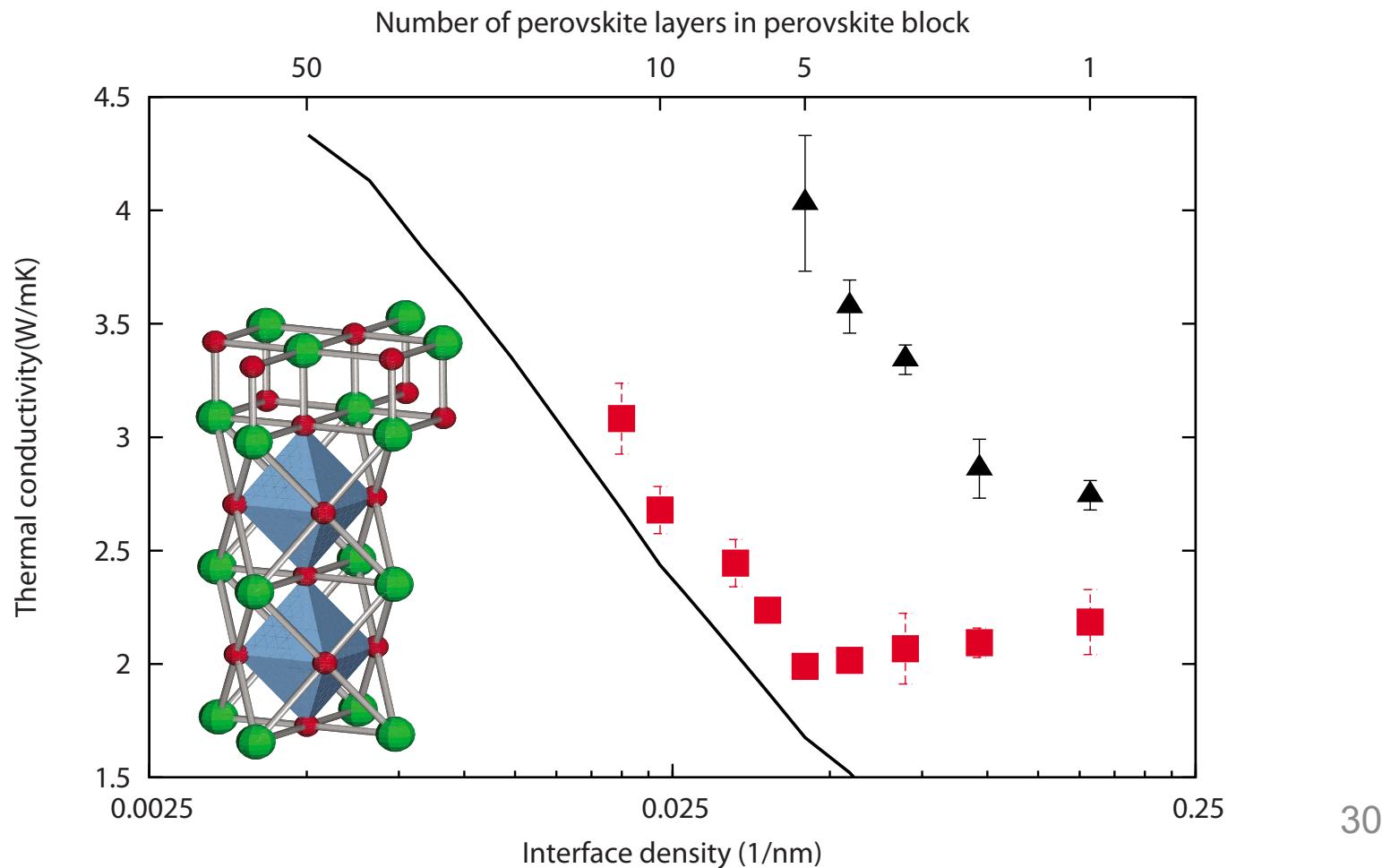
Ravichandran *et al.* Nature Materials **13**, 168 (2013)

"Naturally occurring superlattices": Layered phase crystals

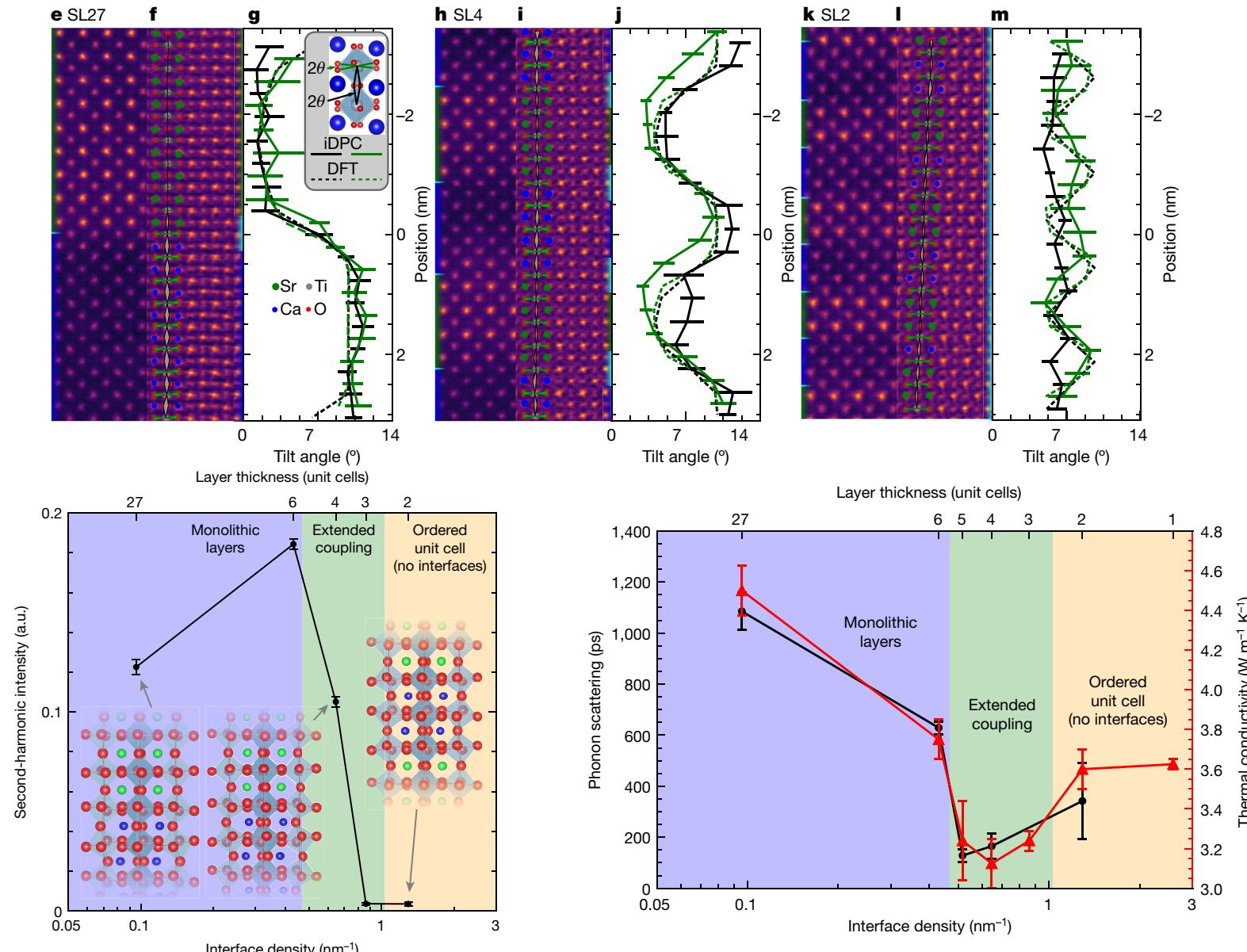
APPLIED PHYSICS LETTERS **95**, 161906 (2009)

Crossover in thermal transport properties of natural, perovskite-structured superlattices

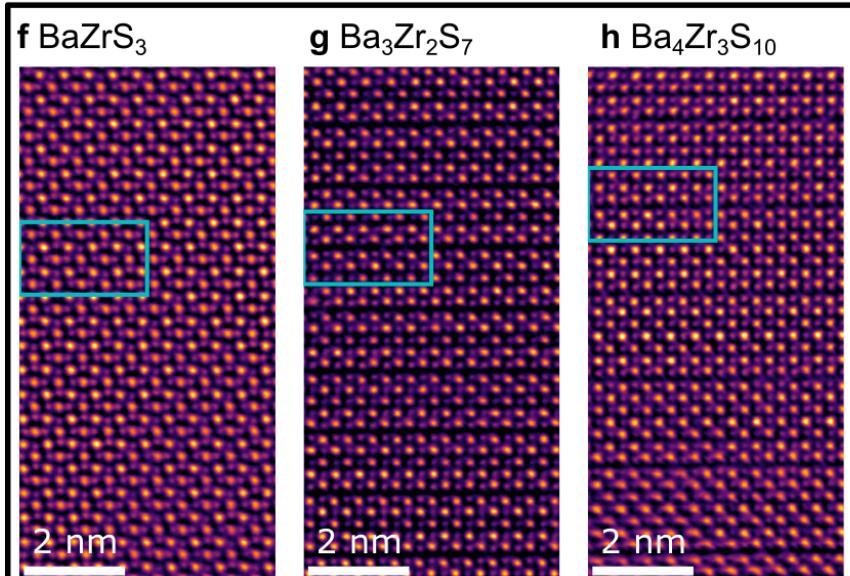
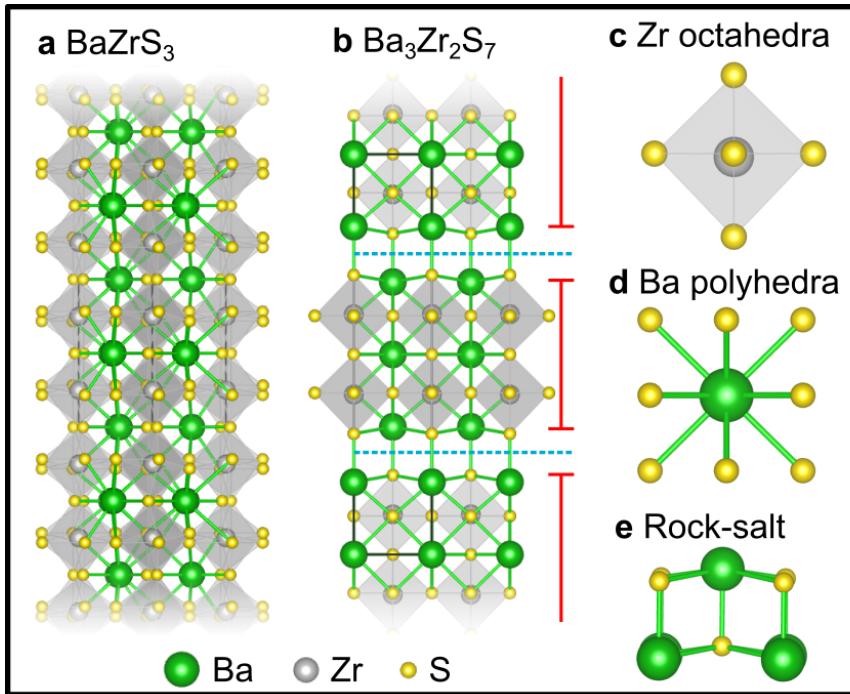
Aleksandr Chernatynskiy,¹ Robin W. Grimes,² Mark A. Zurbuchen,³ David R. Clarke,⁴ and Simon R. Phillpot^{1,a)}



Experimental evidence of phonon coherence

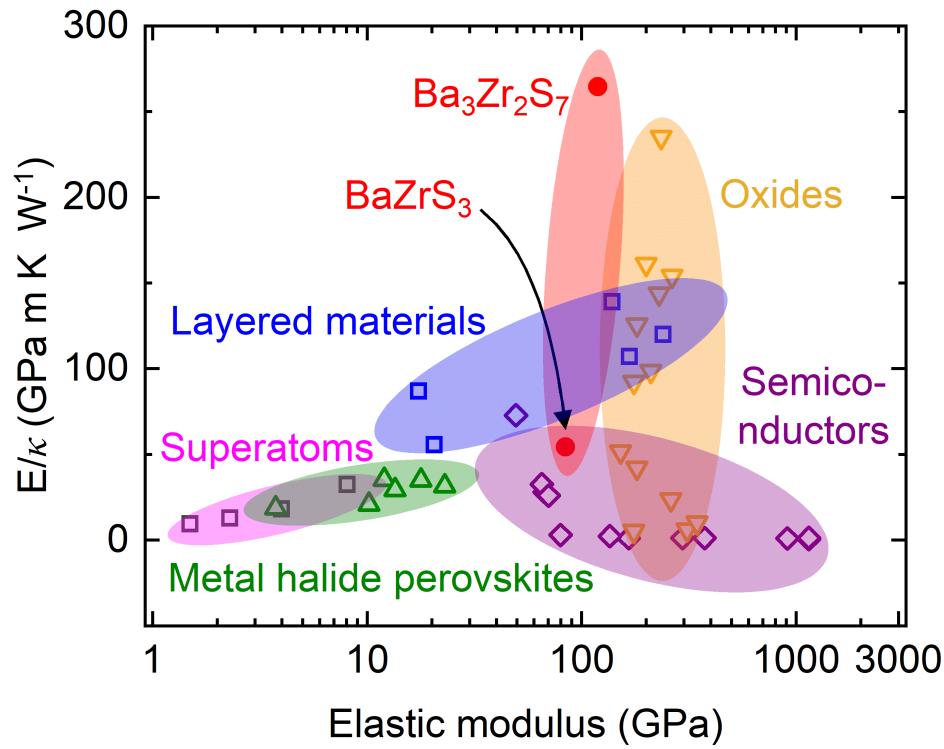


Barium Zirconium Sulfide: Ruddlesden-Popper chalcogenides



$\text{Ba}_3\text{Zr}_2\text{S}_7$: “stiff phonon glass”

- Large anharmonicity
- Layered structure
- Covalent bonding

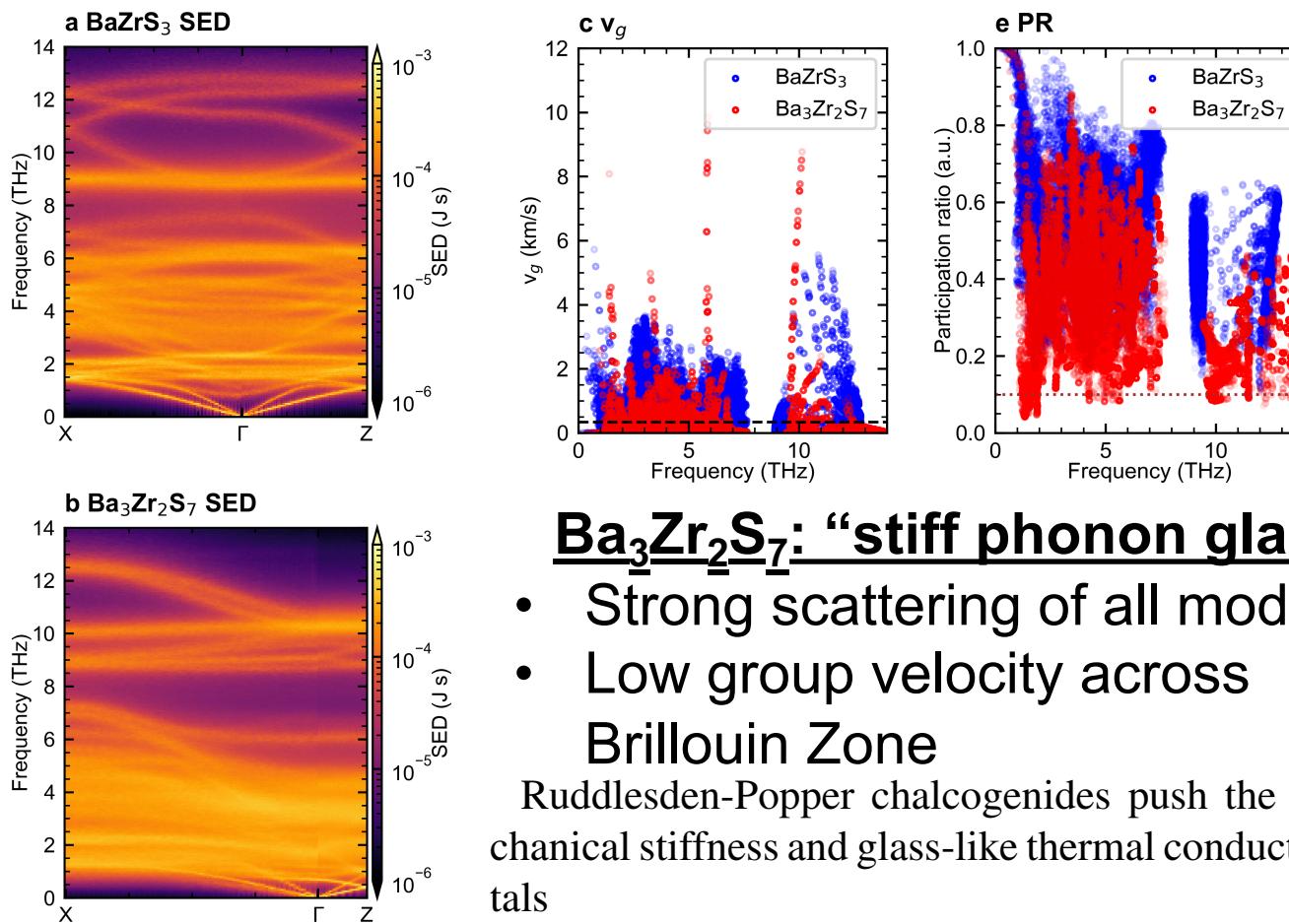


JMR 34, 3819 (2019)

arXiv: 2312.02534

RP BZS: Scattering, low group velocity, localization

$$\kappa = \frac{1}{3} Cv\lambda = \frac{1}{3} Cv_g^2 \tau$$



Ba₃Zr₂S₇: “stiff phonon glass”

- Strong scattering of all modes
- Low group velocity across Brillouin Zone

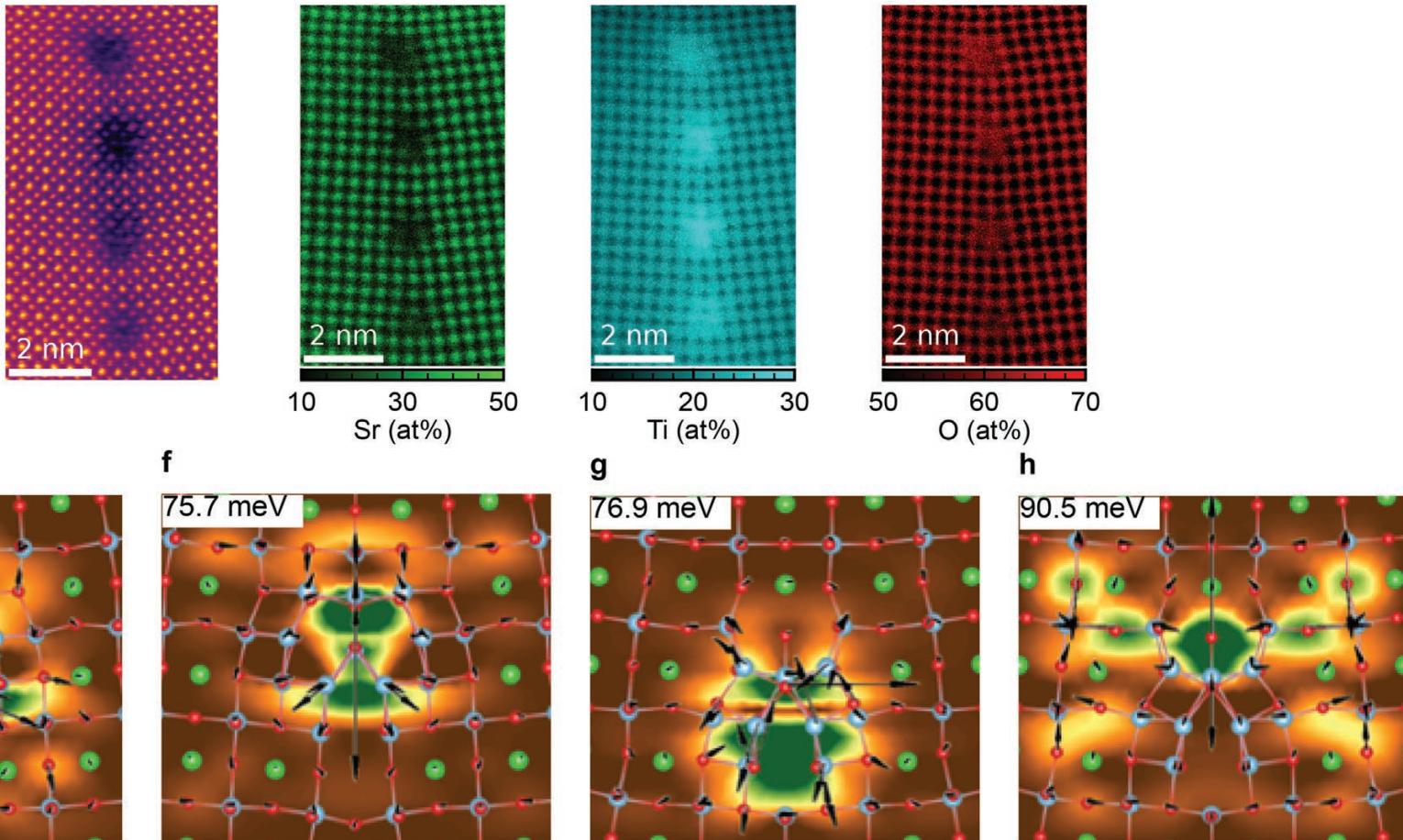
Ruddlesden-Popper chalcogenides push the limit of mechanical stiffness and glass-like thermal conductivity in crystals

Md Shafkat Bin Hoque, Eric R. Hoglund*, Boyang Zhao*, De-Liang Bao, Hao Zhou, Sandip Thakur, Eric Osei-Agyemang, Khalid Hattar, Ethan A. Scott, Mythili Surendran, John A. Tomko, John T. Gaskins, Kiumars Aryana, Sara Makarem, Ganesh Balasubramanian, Ashutosh Giri, Tianli Feng, Jordan A. Hachtel, Jayakanth Ravichandran, Sokrates T. Pantelides, and Patrick E. Hopkins*

Advanced microscopy (at ORNL) leads to discovery and understanding of emergent phononic properties

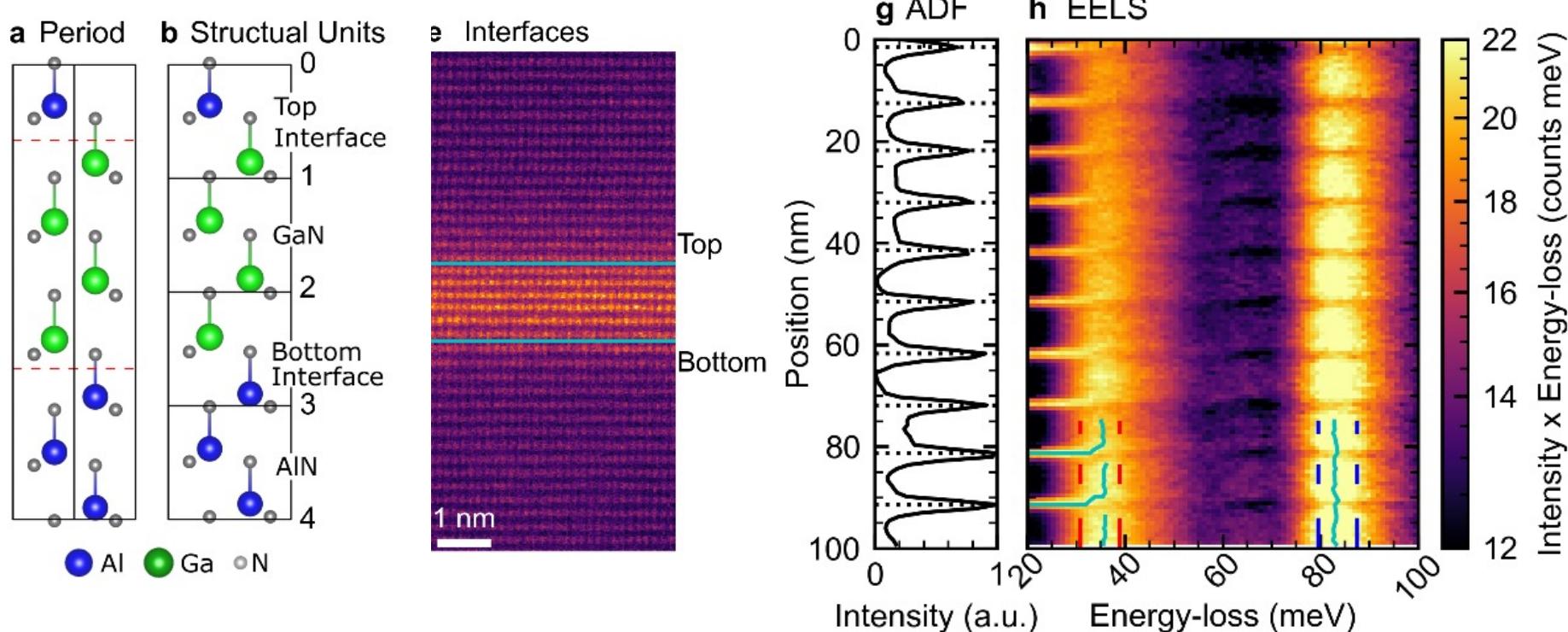
Localized vibrational modes at grain boundaries (STO)

Vibrational modes at grain boundaries are different from “bulk”



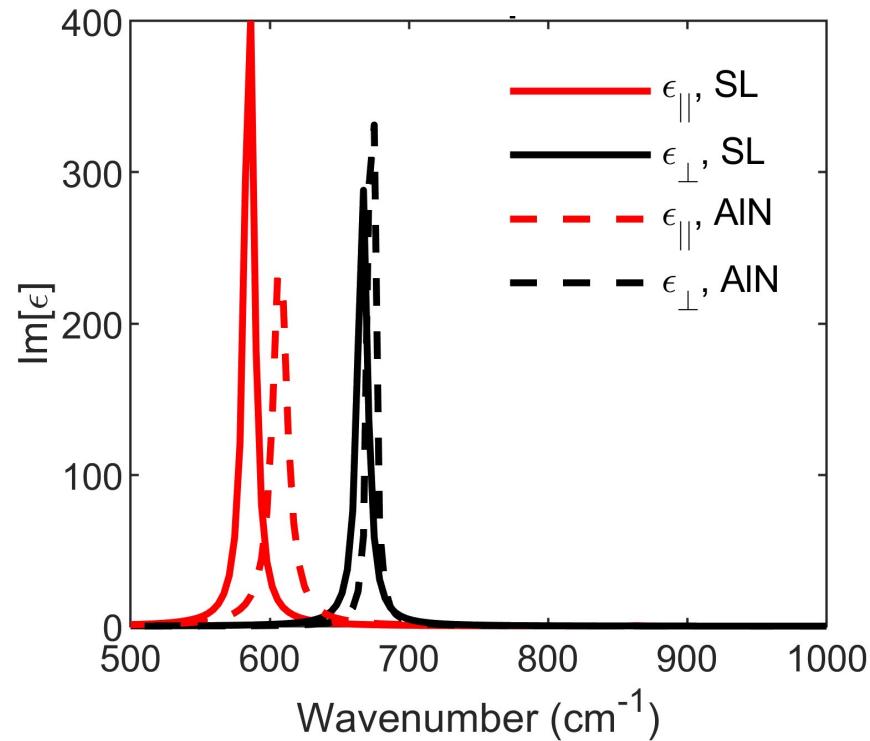
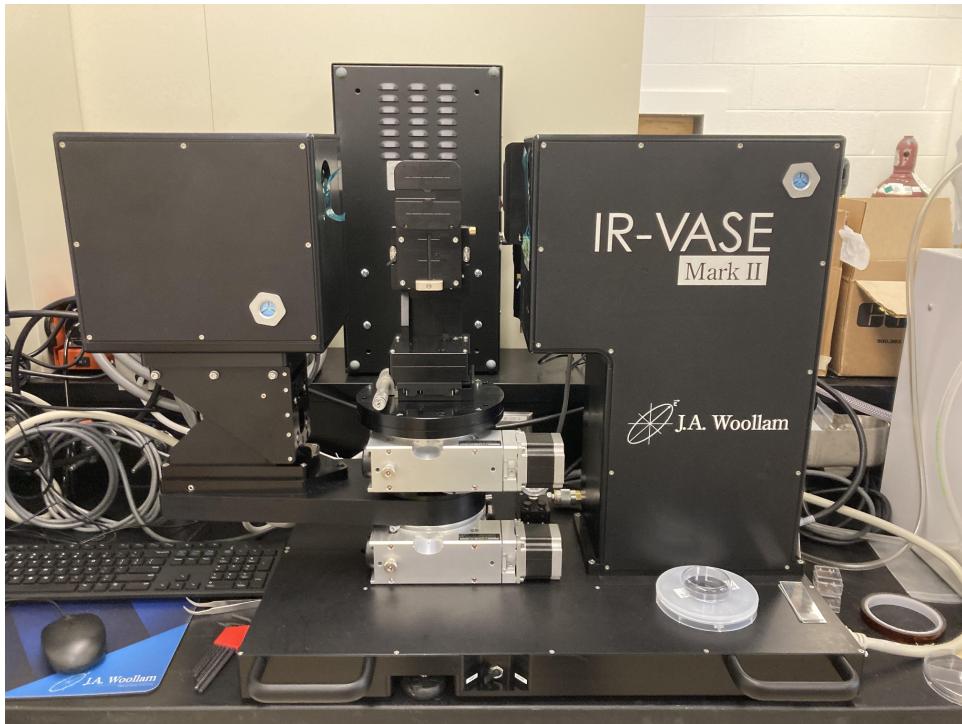
Advanced microscopy (at ORNL) leads to discovery and understanding of emergent phononic properties

Asymmetry in vibrational modes at polar interfaces
Phononically, AlN/AlGaN is different from AlGaN/AlN



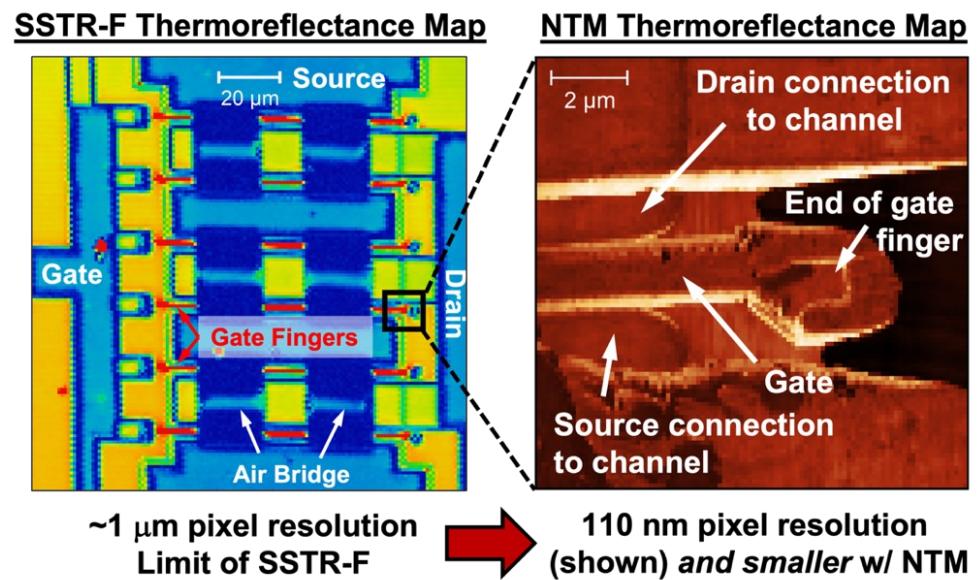
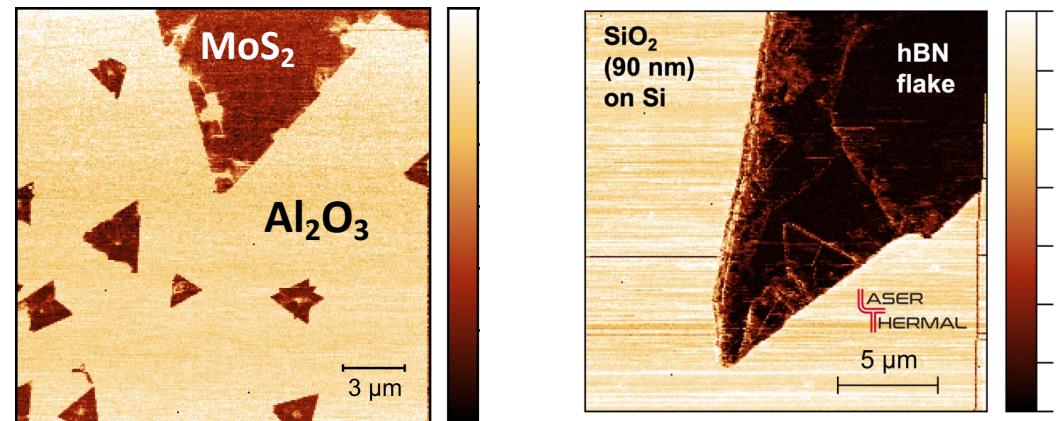
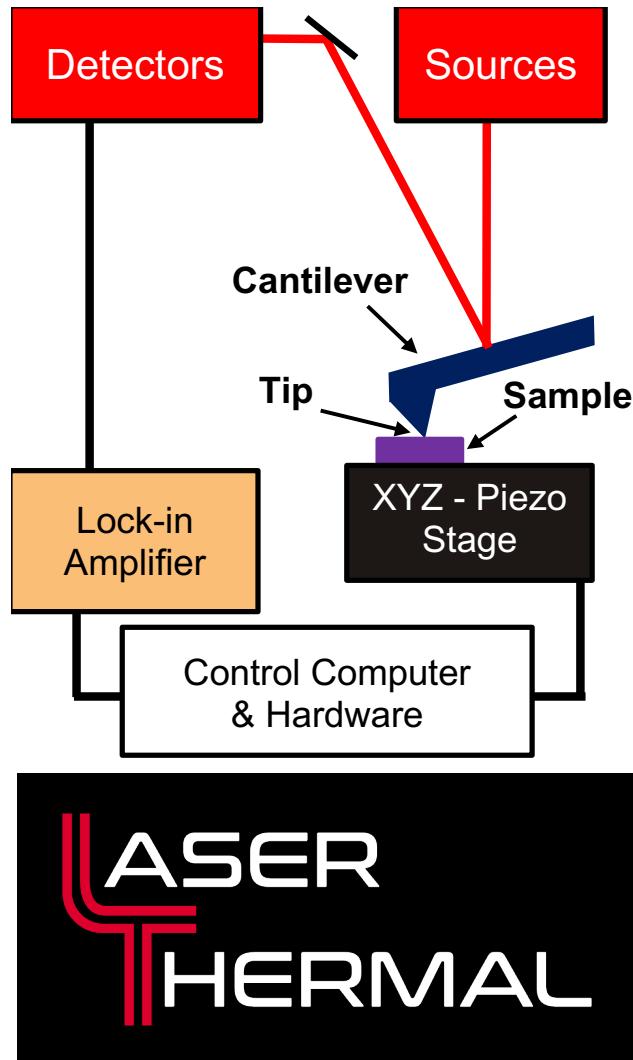
New techniques we're standing up at UVA to complement microscopy collaborations at ORNL

Infrared variable angle spectroscopic ellipsometry



New techniques we're standing up at UVA to complement microscopy collaborations at ORNL

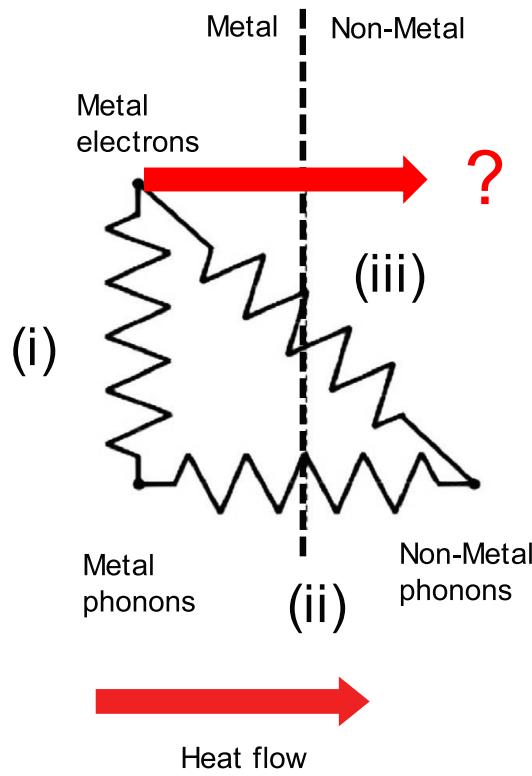
Nanoscale Thermal Microscopy (NTM) Pump-probe (SSTR) with ~ 10 nm areal resolution



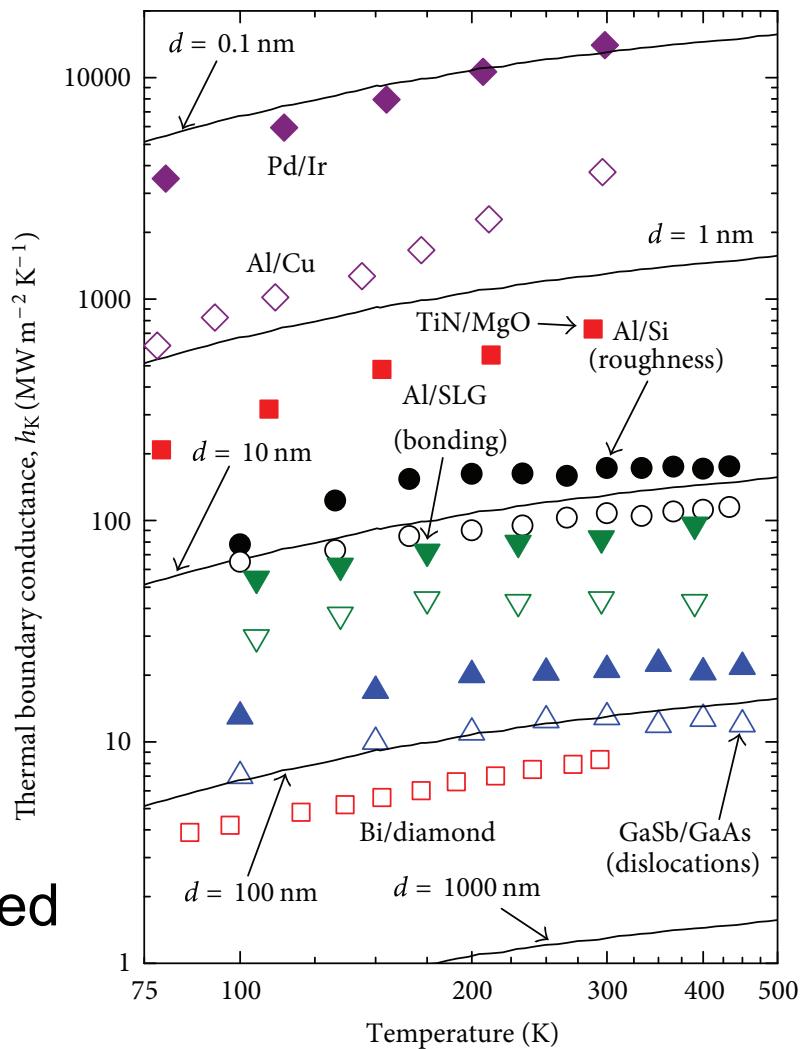
Theme: Coupling of carriers across interfaces lead to unique energy states and modes of heat transfer

- How do we measure the thermal conductivity of materials and the thermal boundary resistance at their interfaces?
- Thermal conductivity of superlattices
 - Minimum thermal conductivity and creating “crystals of interfaces”
- **Interfacial heat transfer control of the IR properties of solids**
 - **Near field radiative interfacial heat transfer with plasmon-polaritons and phonon-polaritons**
- Transient temperature changes during plasma-surface interactions
 - “Plasma cooling”

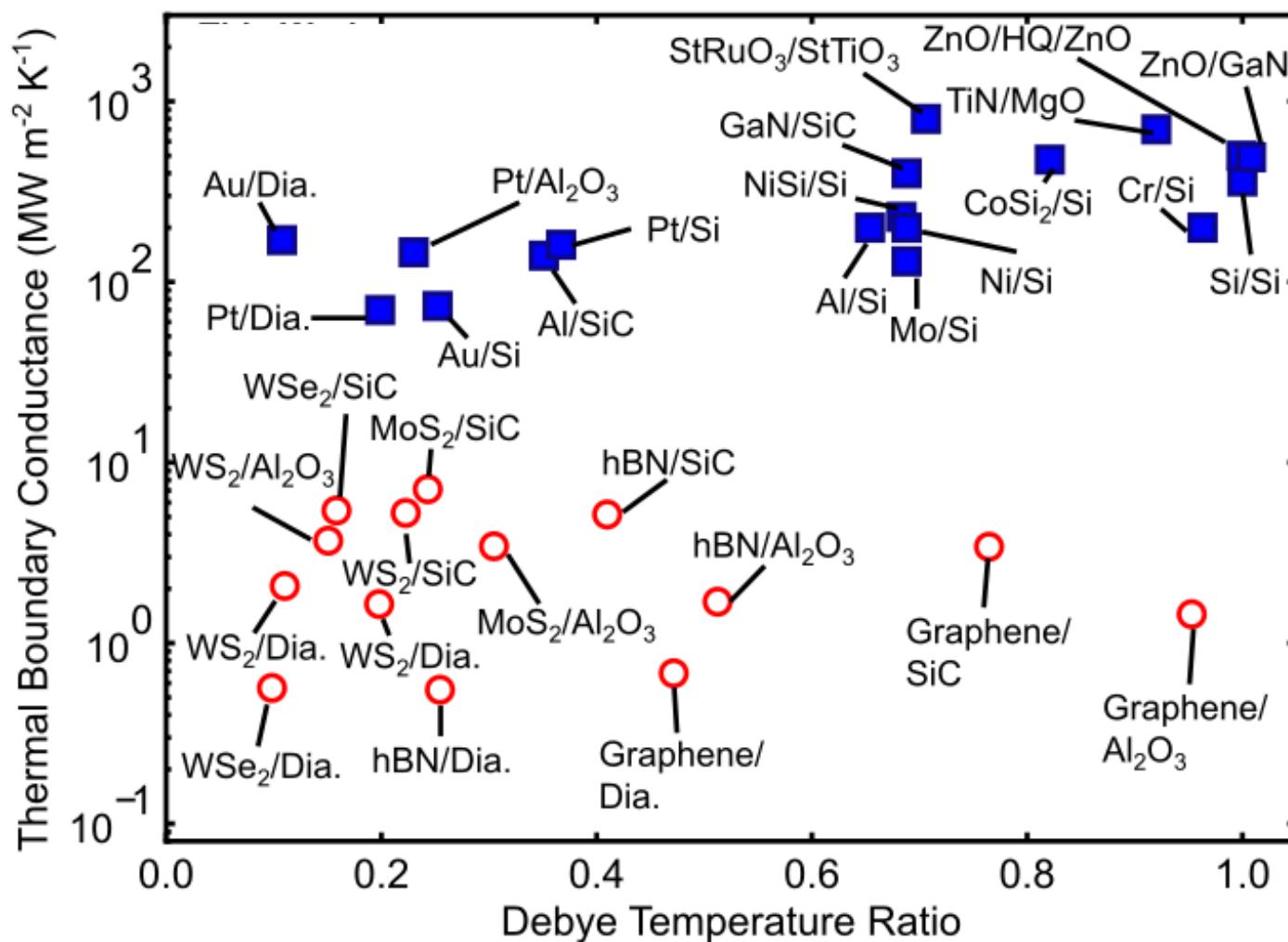
Phonon dominated TBC < electron dominated TBC



Metal/non-metal = Phonon dominated
Does it have to be?



TBCs are even lower at interfaces of 2D materials

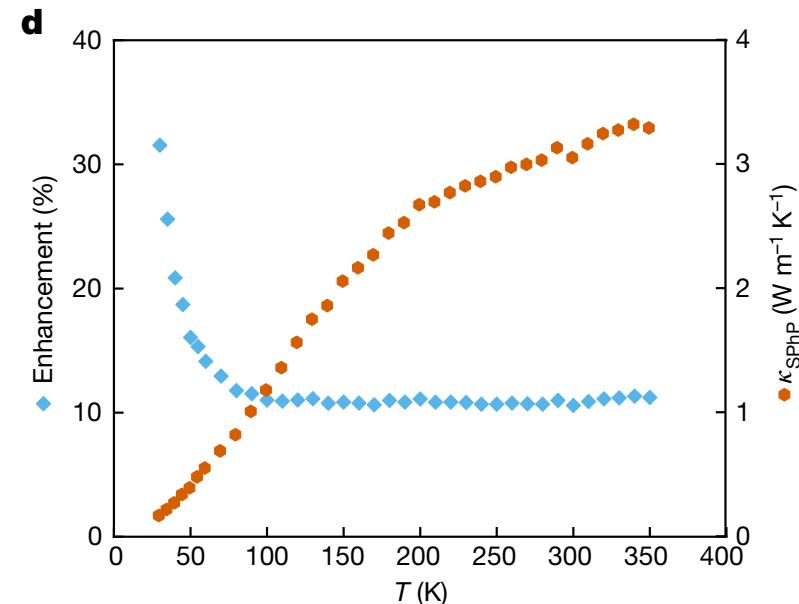
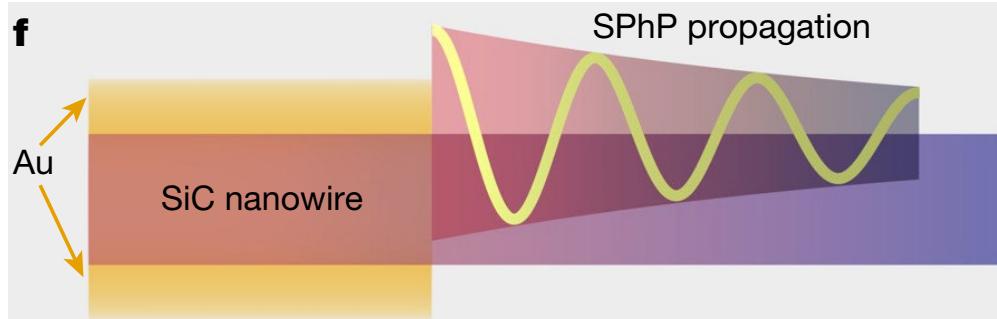


How can we increase TBC across metal/nonmetal interfaces?

Remarkable heat conduction mediated by non-equilibrium phonon polaritons

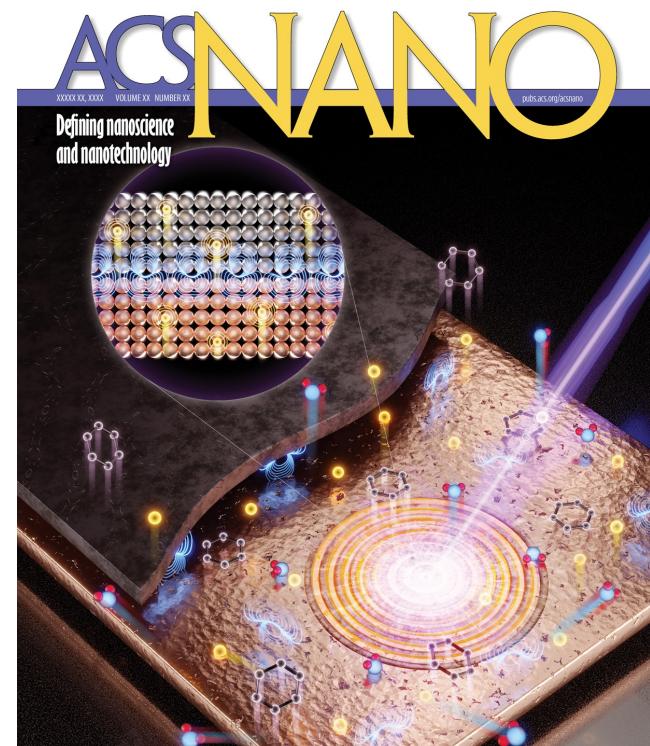
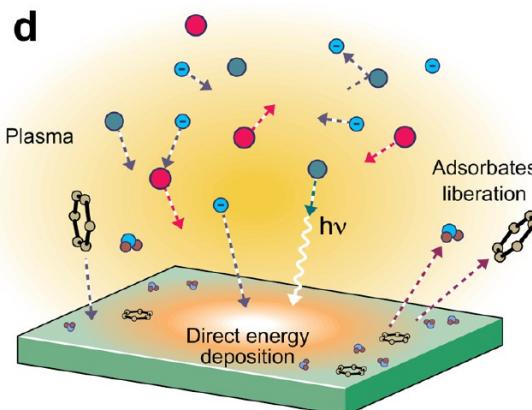
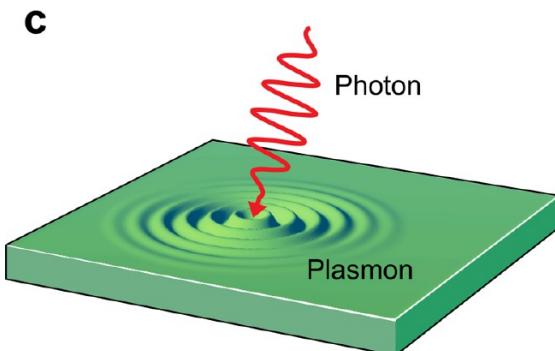
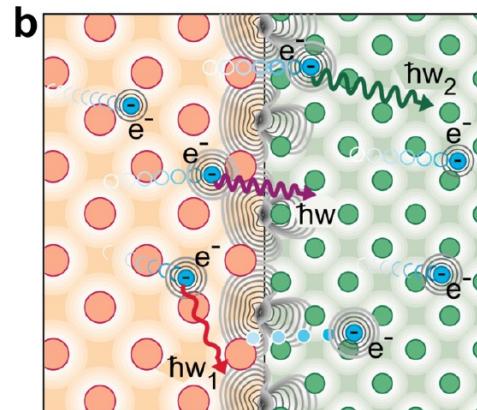
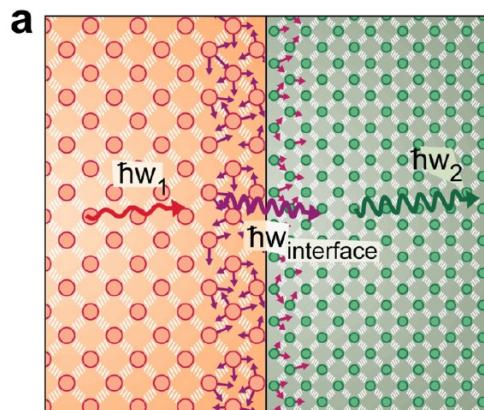
Zhiliang Pan¹, Guanyu Lu¹, Xun Li², James R. McBride³, Rinkle Juneja², Mackey Long⁴, Lucas Lindsay², Joshua D. Caldwell¹ & Deyu Li¹✉

Nature | Vol 623 | 9 November 2023 | **307**

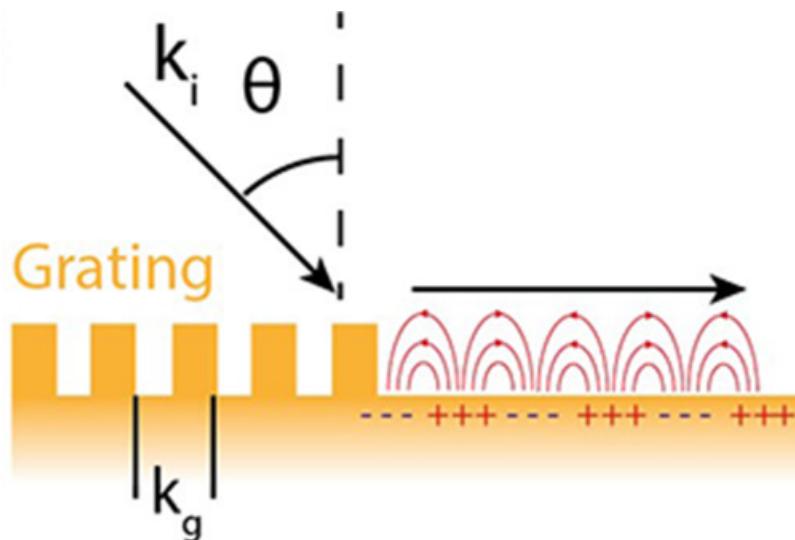


Ultrafast and Nanoscale Energy Transduction Mechanisms and Coupled Thermal Transport across Interfaces

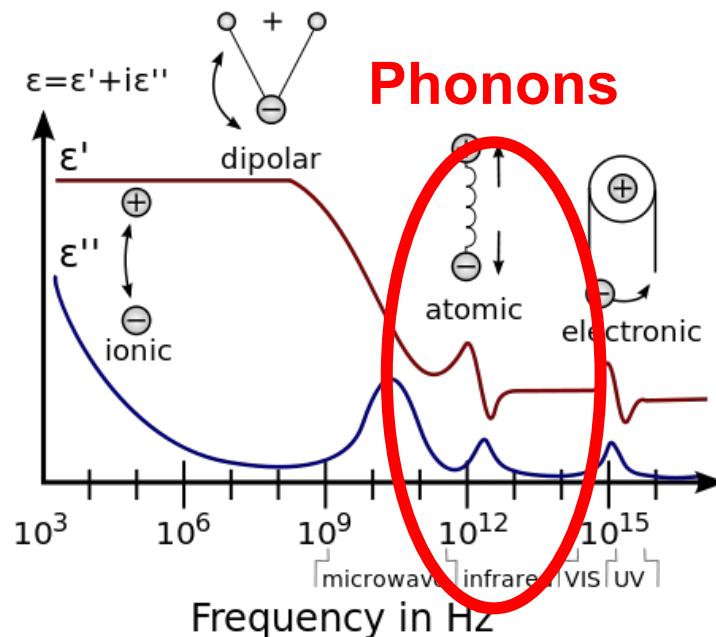
Ashutosh Giri,* Scott G. Walton, John Tomko, Niraj Bhatt, Michael J. Johnson, David R. Boris, Guanyu Lu, Joshua D. Caldwell, Oleg V. Prezhdo, and Patrick E. Hopkins*



Phonon-polaritons: enhancing heat transfer?



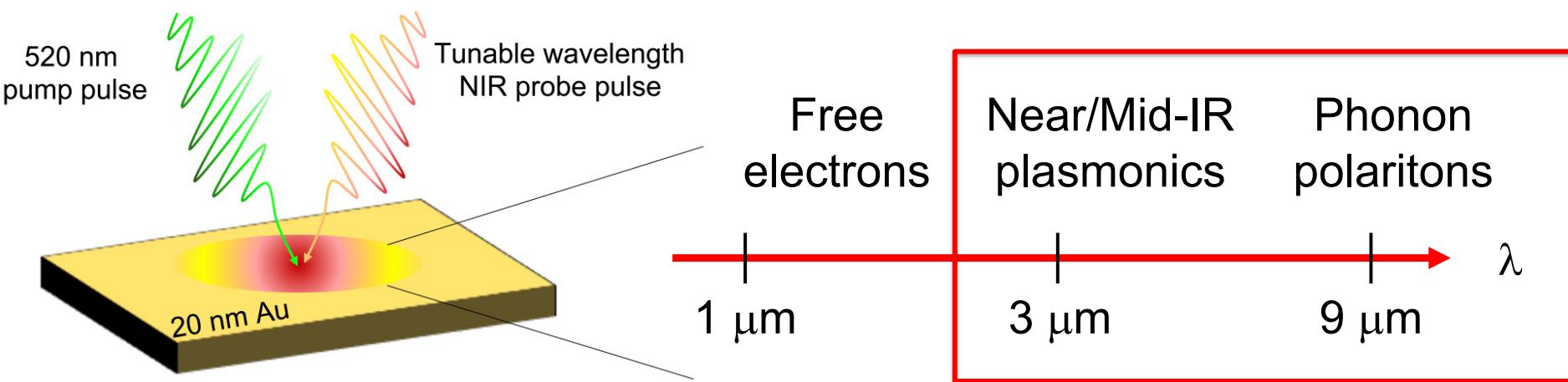
J. Appl. Phys. **125**, 191102 (2019)



- Light couples with electric dipole creating quasiparticle
- At IR wavelengths: Phonon polariton!
- PhP quasiparticles can propagate at $\sim 1\%$ of the speed of light!
- Prior evidence of thermal conductivity enhancement during excitation of PhPs

$$\kappa = \frac{1}{3} Cv\lambda = \frac{1}{3} Cv_g^2 \tau$$

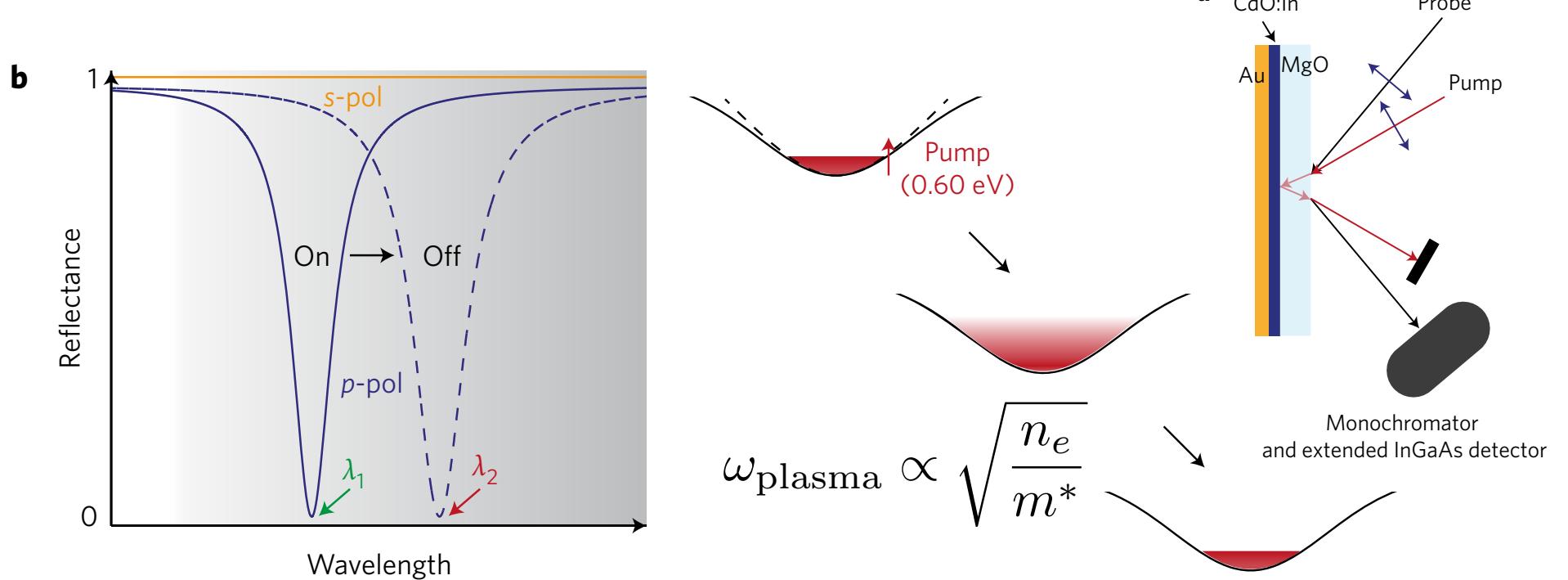
Pumping with heat, probing in IR away from e^- transitions



- **Thermally modulated plasmon-polariton in CdO**
 - Electron TBC controls plasmon response
- **Thermally modulated phonon-polariton in h-BN**
 - Broadband radiative flux from Au heat PhP

Can we use interfacial heat transfer to manipulate the optical properties of materials?

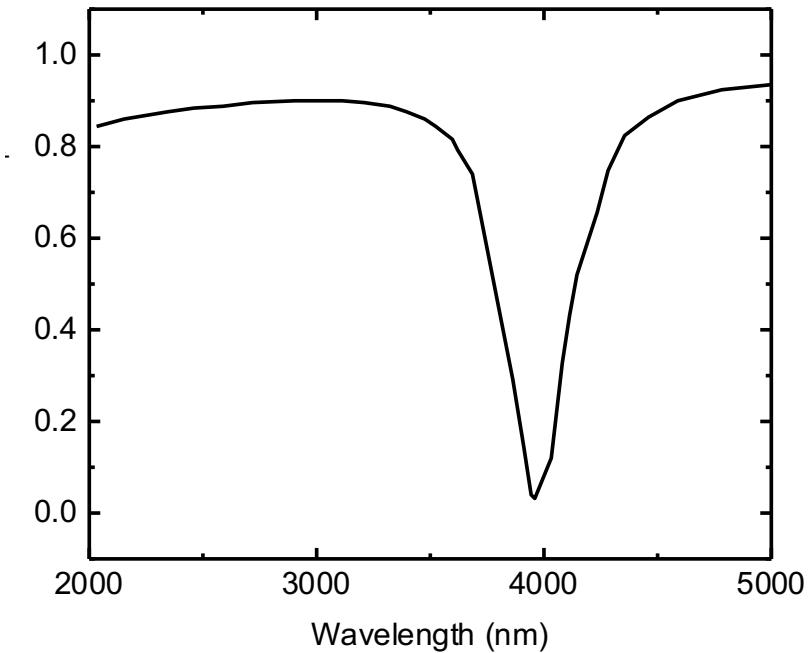
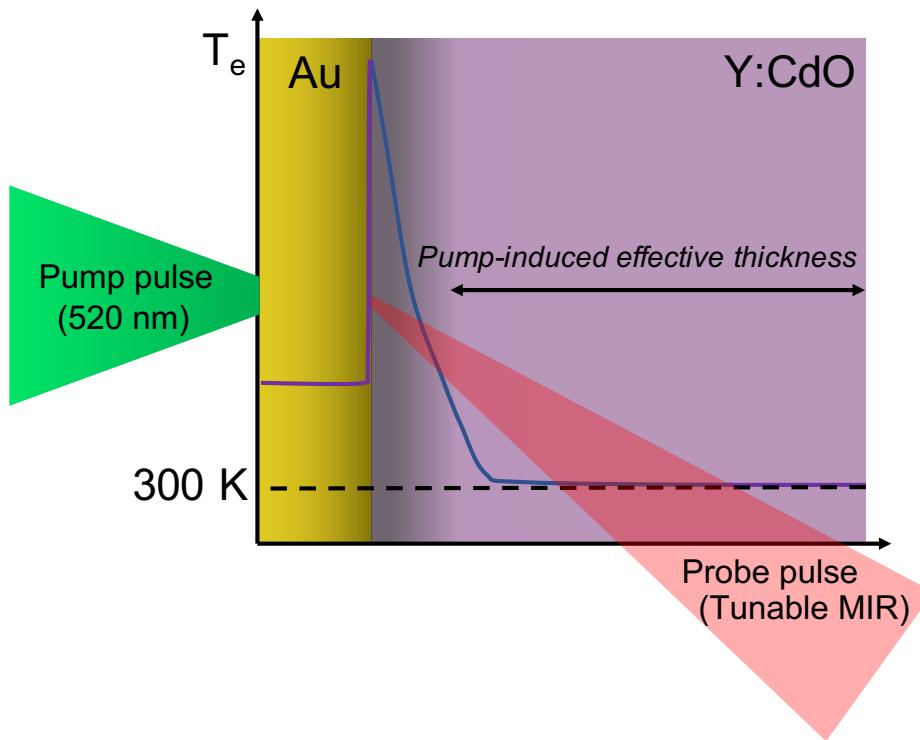
Carrier distributions modulate optical properties of ENZ plasmonic modes in CdO



Carrier scattering and relaxation drives optical properties

Recall earlier example

- Changing in carrier density via short pulse absorption can modulate plasmon resonance in CdO
- **Can we modulate and control plasmon lifetimes with heat?**

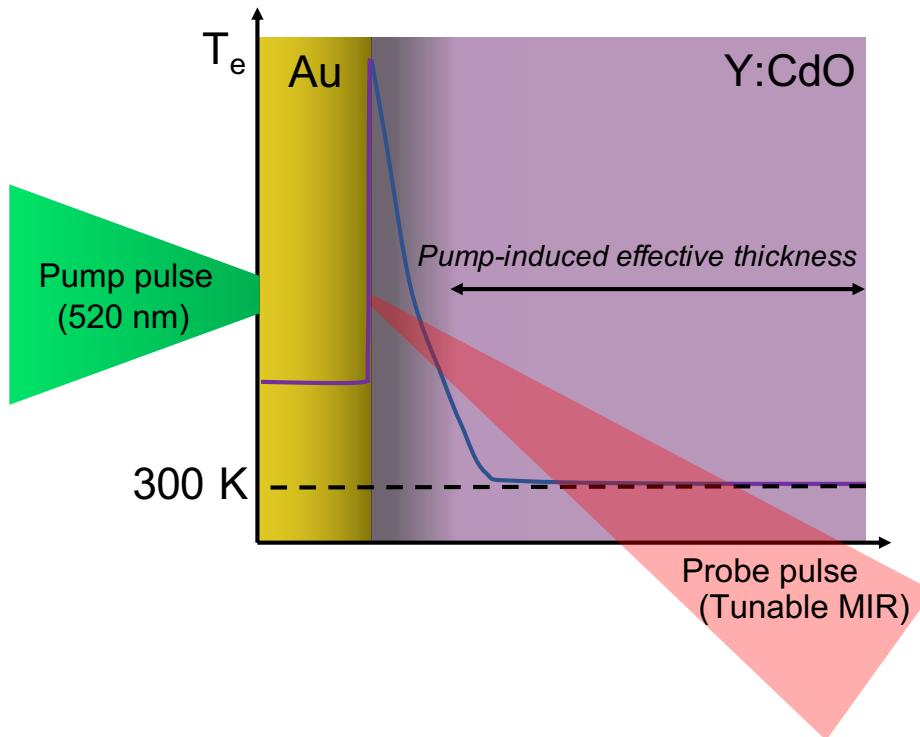


Tomko *et al.* *Nature Nano.* **16**, 47

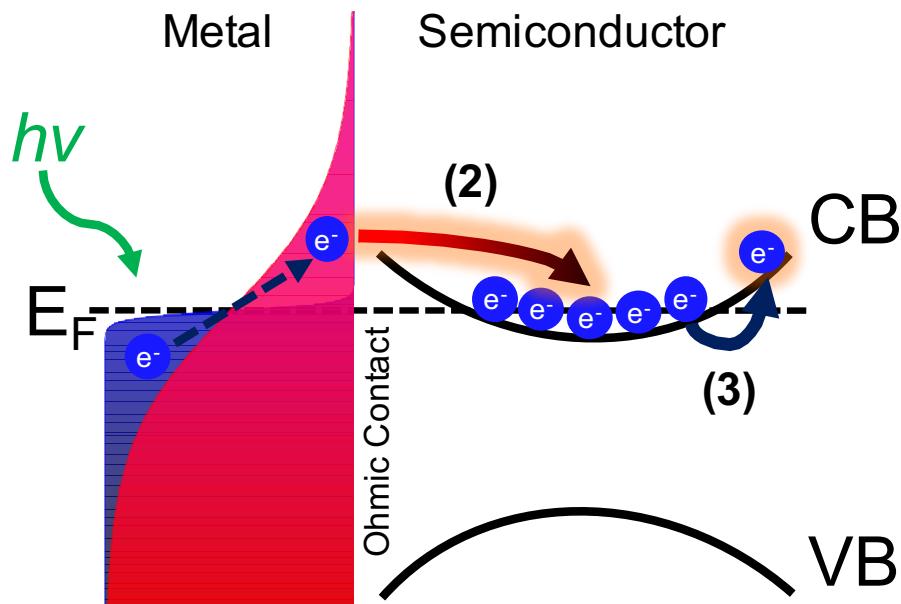
Collaboration: Caldwell (Vanderbilt), Maria (PSU), Prezhdo (USC) ⁴⁶

Electron mediated TBC at metal/doped non-metal is the key

- Energy easily transmitted across interface when out of equilibrium with phonons
- Slowly “goes back” across the interface when diffusive
- **This “ballistic thermal injection” is different than charge injection**



Ballistic thermal injection (Energy transfer)



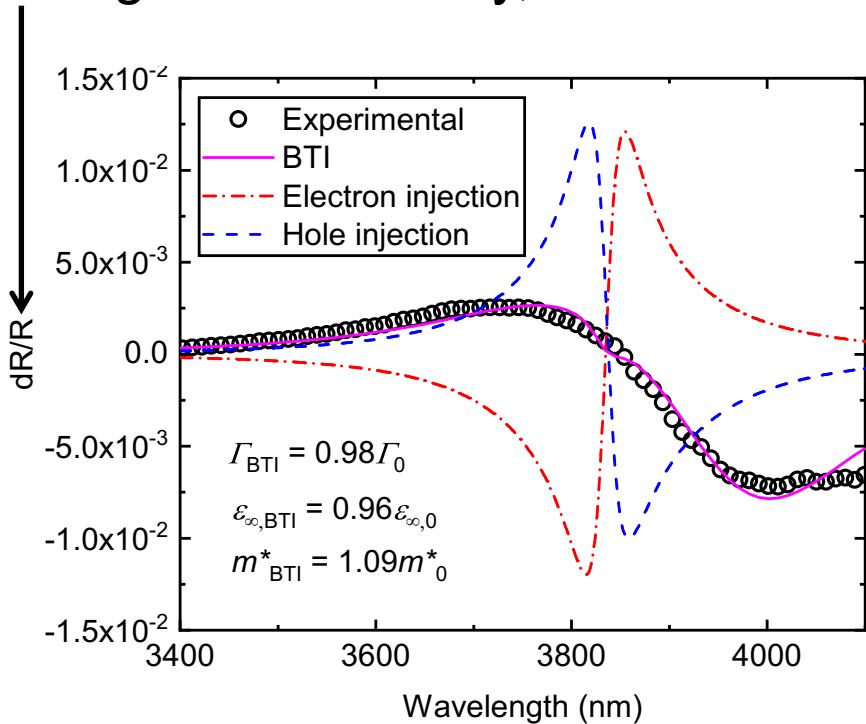
Tomko *et al.* *Nature Nano.* **16**, 47

Collaboration: Caldwell (Vanderbilt), Maria (PSU), Prezhdo (USC) ⁴⁷

Pump electrons in Au, probe plasmon in CdO

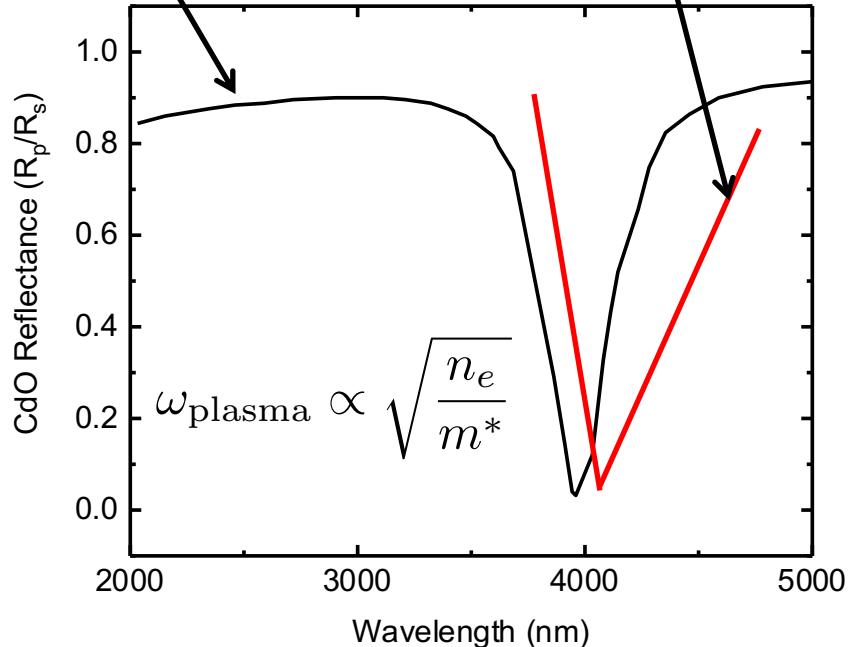
Asymmetric red shift in ENZ plasmon mode due to BTI

Note we are measuring
change in reflectivity, dR



Au/CdO/sapphire
absorption response

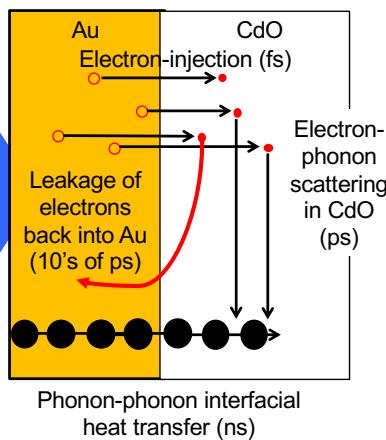
Thermo-modulated
response from indirect
heating form Au film
(qualitative)



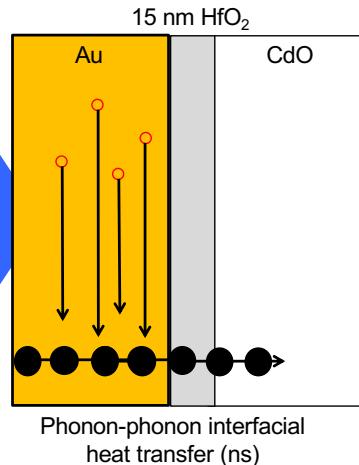
Tomko *et al.* *Nature Nano.* **16**, 47

Collaboration: Caldwell (Vanderbilt), Maria (PSU), Prezhdo (USC) 48

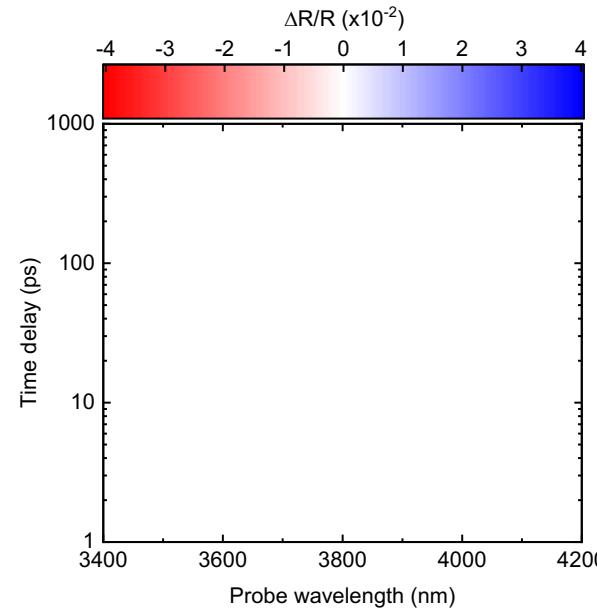
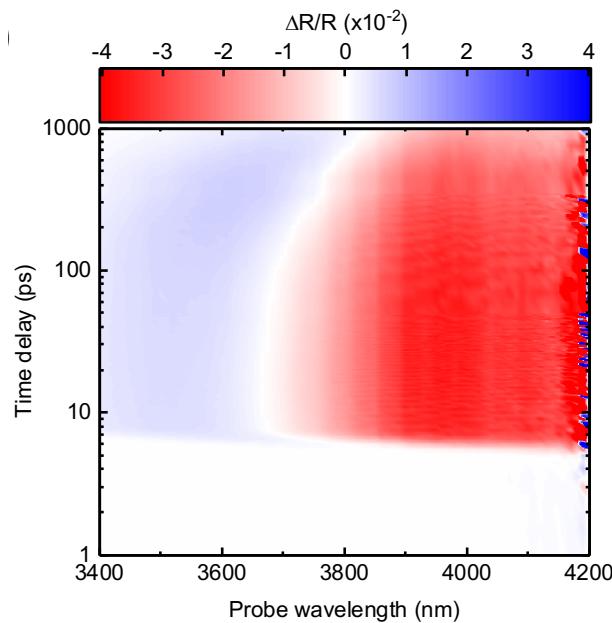
And it's not an optical artifact



15 nm HfO_2 layer prevents any electron energy from moving from **Au** to **CdO**, resulting in no measurable response



Au/CdO

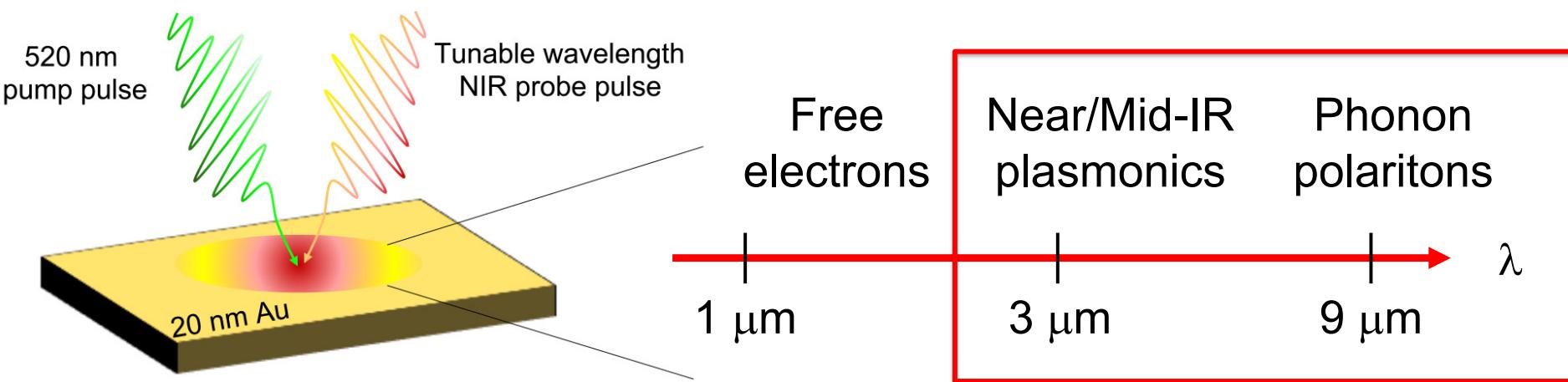


**Au/HfO₂/CdO
(same scale)**

Tomko *et al.* *Nature Nano.* **16**, 47

Collaboration: Caldwell (Vanderbilt), Maria (PSU), Prezhdo (USC) **49**

Pumping with heat, probing in IR away from e^- transitions



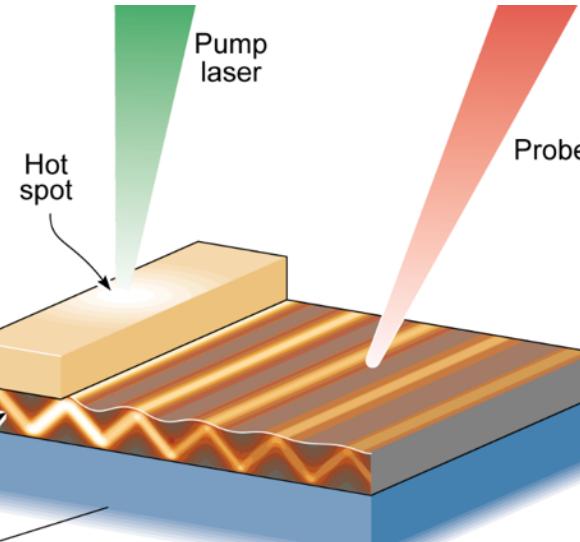
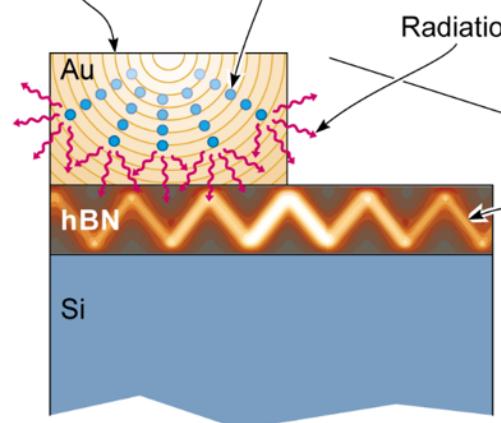
- **Thermally modulated plasmon-polariton in CdO**
 - Electron TBC controls plasmon response
- **Thermally modulated phonon-polariton in h-BN**
 - Broadband radiative flux from Au heat PhP

Can we modulate and control PhP lifetimes with heat?

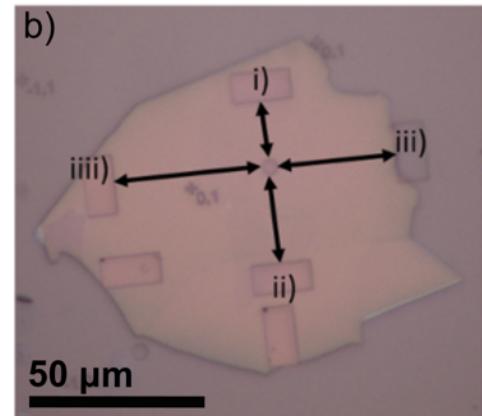
- MIR time delayed probe of ΔR across upper Reststrahlen band
- How does heat from Au indirectly change PhP in h-BN?

a)

Phonons from hot spot
Ballistic electrons



b)

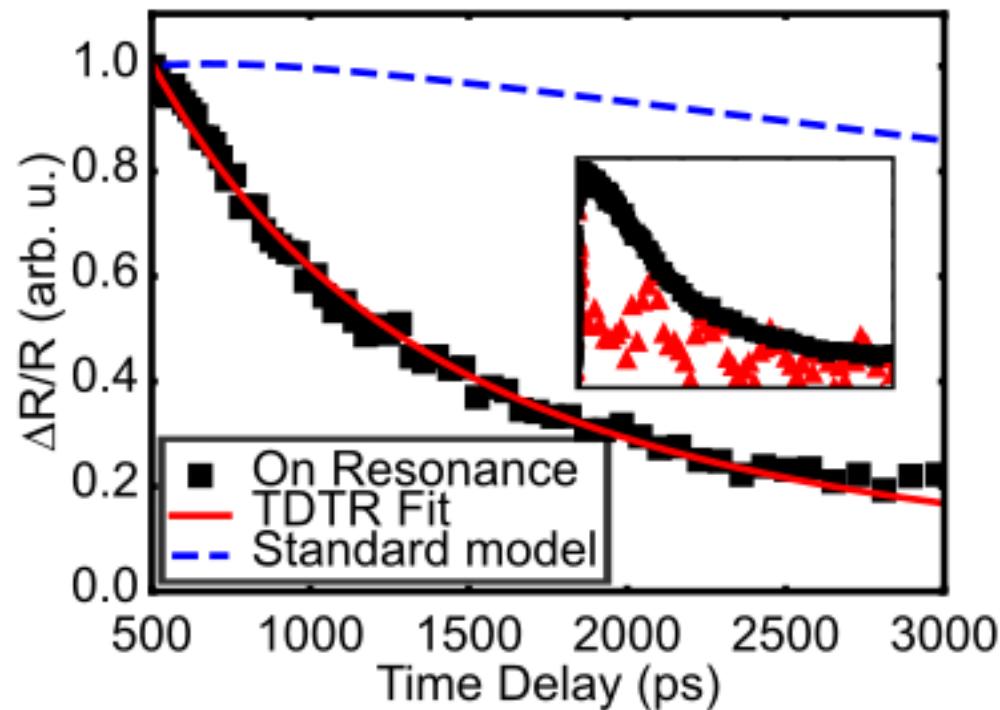
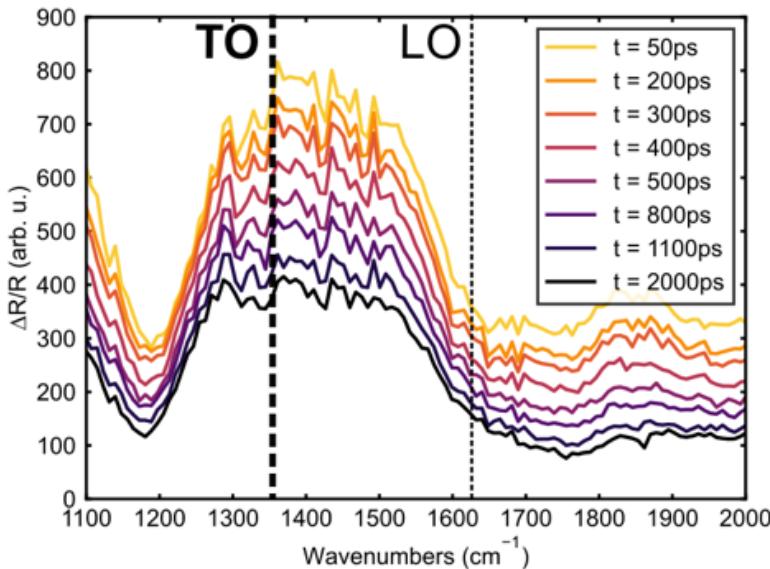


Ultrafast evanescent heat transfer across solid interfaces via hyperbolic phonon-polaritons in hexagonal boron nitride

William Hutchins,¹ John A. Tomko,¹ Dan M. Hirt,¹ Saman Zare,¹ Joseph R. Matson,² Katja Diaz-Granados,² Mingze He,³ Thomas Pfeifer,¹ Jiahua Li,⁴ James Edgar,⁴ Jon-Paul Maria,⁵ Joshua D. Caldwell,^{2,3,&} Patrick E. Hopkins^{1,6,7,%}

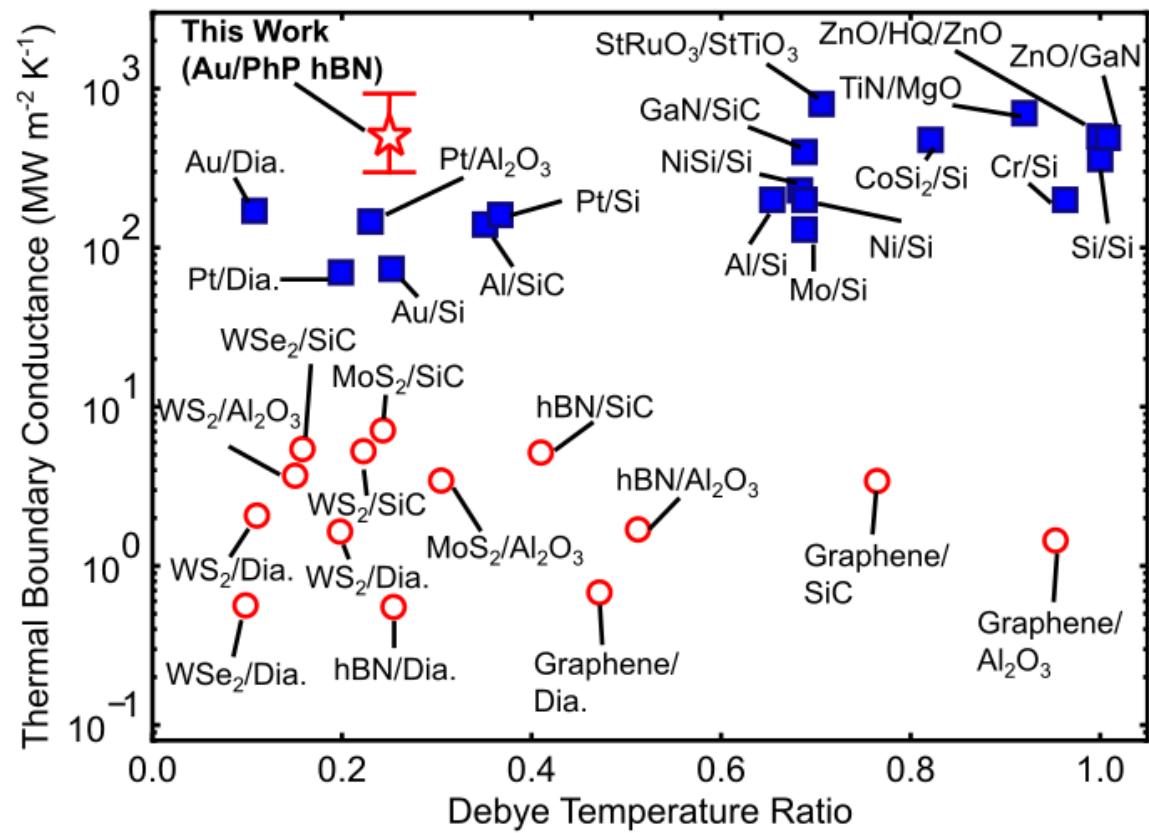
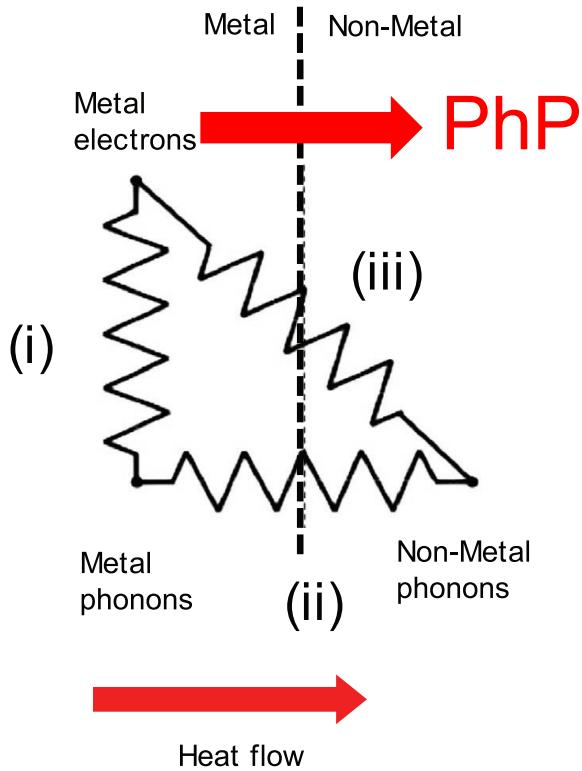
Enhanced heat sinking with PhP coupling

- Temperatures of PhP modes decay an order of magnitude faster than non-PhP modes



Enhanced heat sinking with PhP coupling

- Intensely heated nonequilibrium electrons emit energy in near field to HPhP's in hBN
- Provide additional heat transfer mechanism across Au/hBN interface to enhance TBC

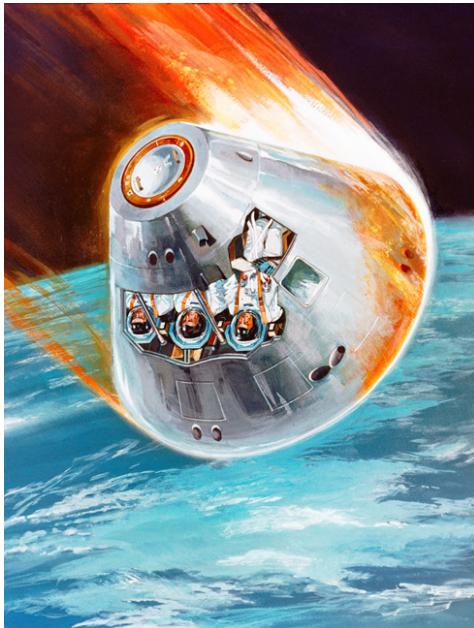


Outline

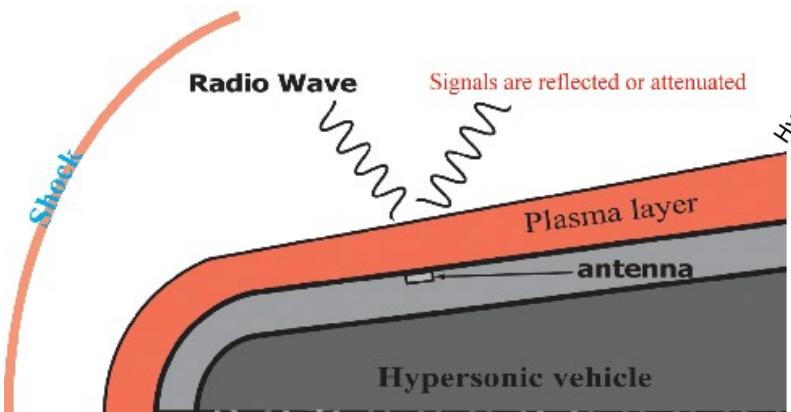
Theme: Coupling of carriers across interfaces lead to unique energy states and modes of heat transfer

- How do we measure the thermal conductivity of materials and the thermal boundary resistance at their interfaces?
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 - Minimum thermal conductivity and creating “crystals of interfaces”
- Interfacial heat transfer control of the IR properties of solids
 - Near field radiative interfacial heat transfer with plasmon-polaritons and phonon-polaritons
- **Transient temperature changes during plasma-surface interactions**
 - **“Plasma cooling”**

Plasma surface interactions

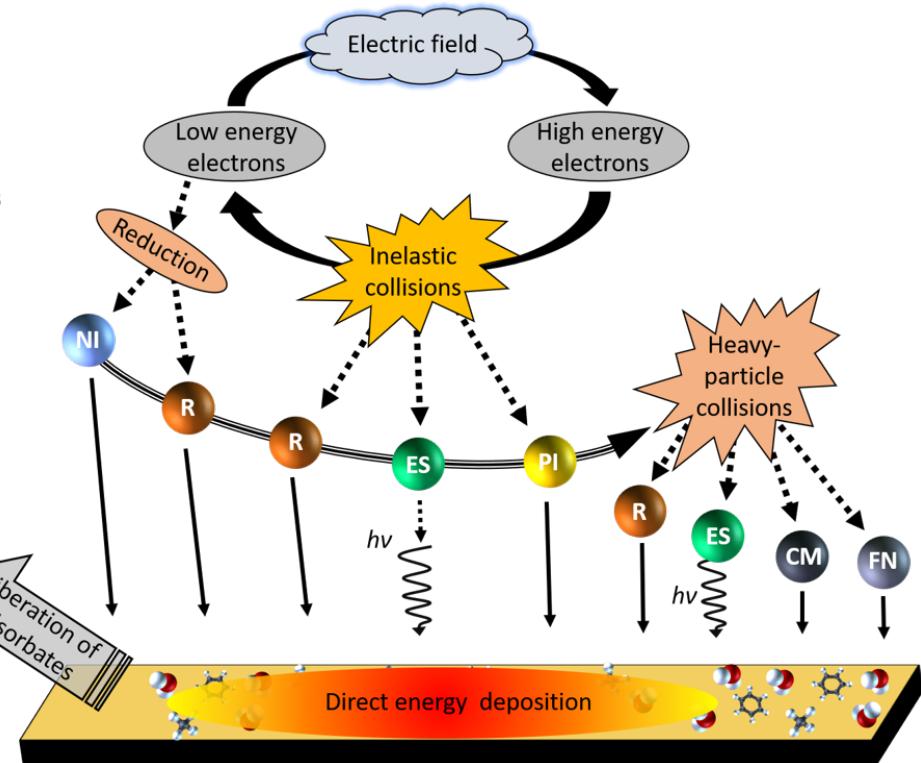


Re-entry vehicle: “All you can do is pray for it”
J. App. Phys. news referencing: 117 233301 (2015)



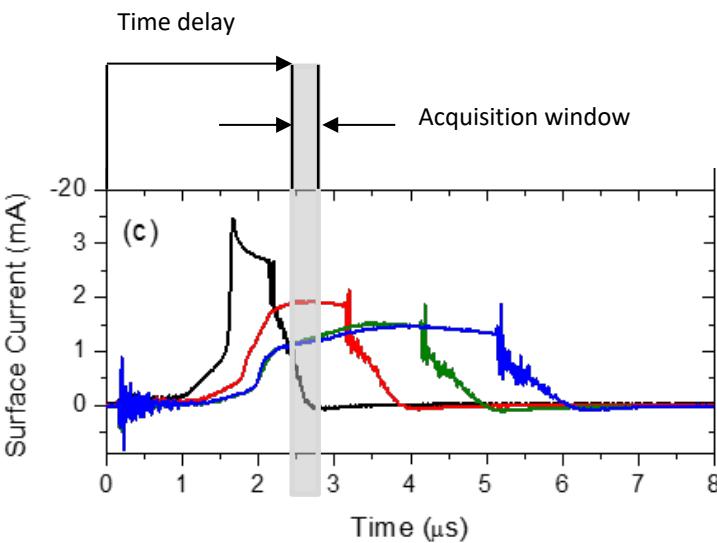
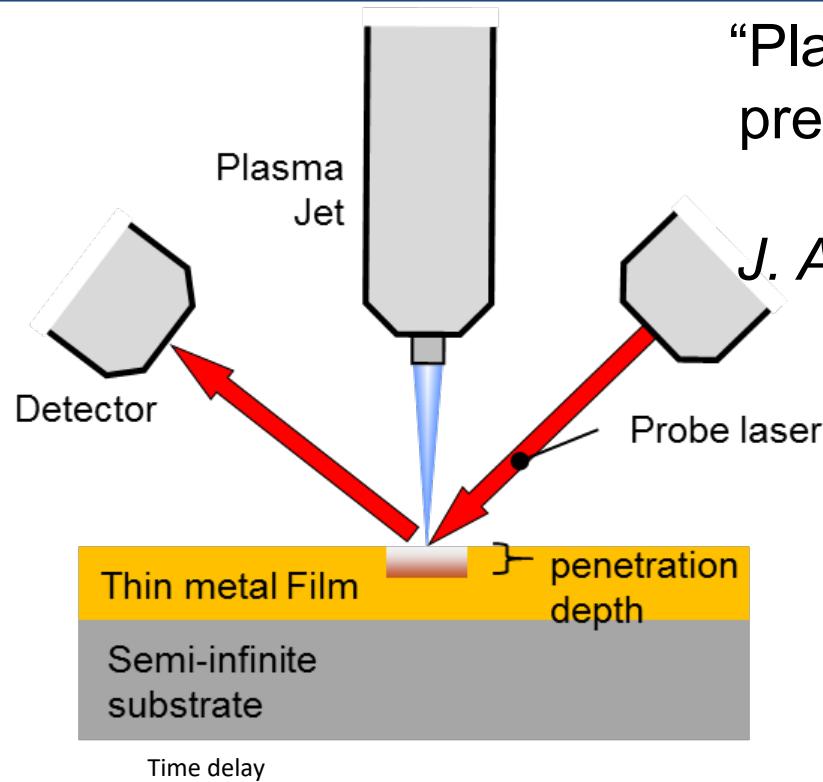
Kim: “Plasma communication for radio blackout during hypersonic flight”

- NI Negative ions
- R Radicals
- ES Excited species
- PI Positive ions
- CM Compounds
- FN Fast neutrals



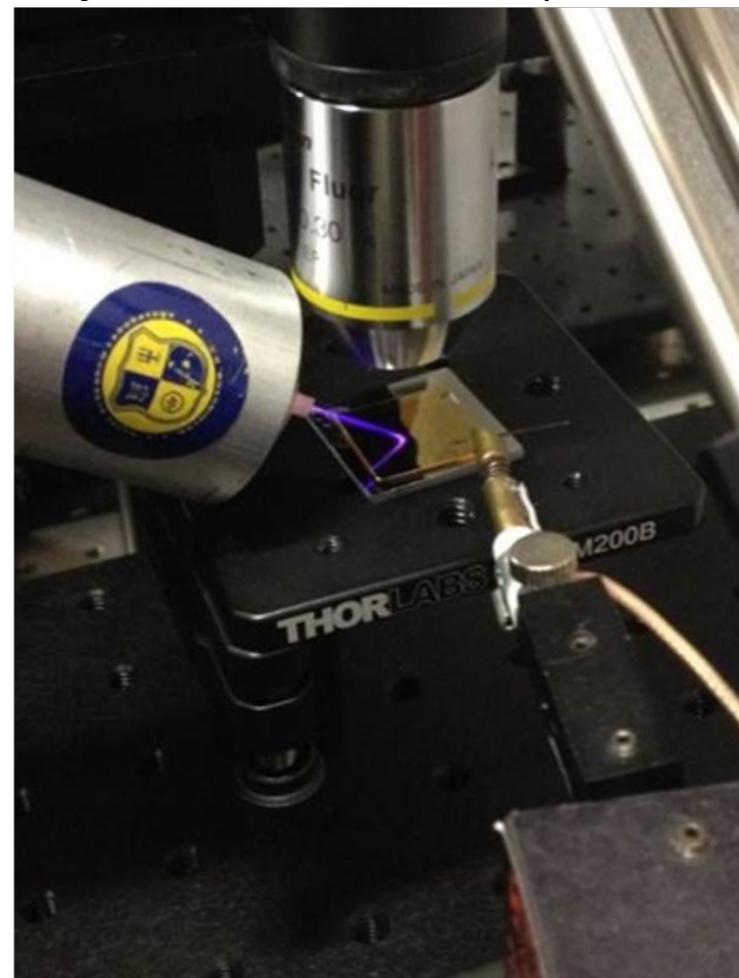
“Plasma-induced surface cooling,” *Nature Comm.* 13, 2623 (2022)

Thermoreflectance for plasma diagnostics



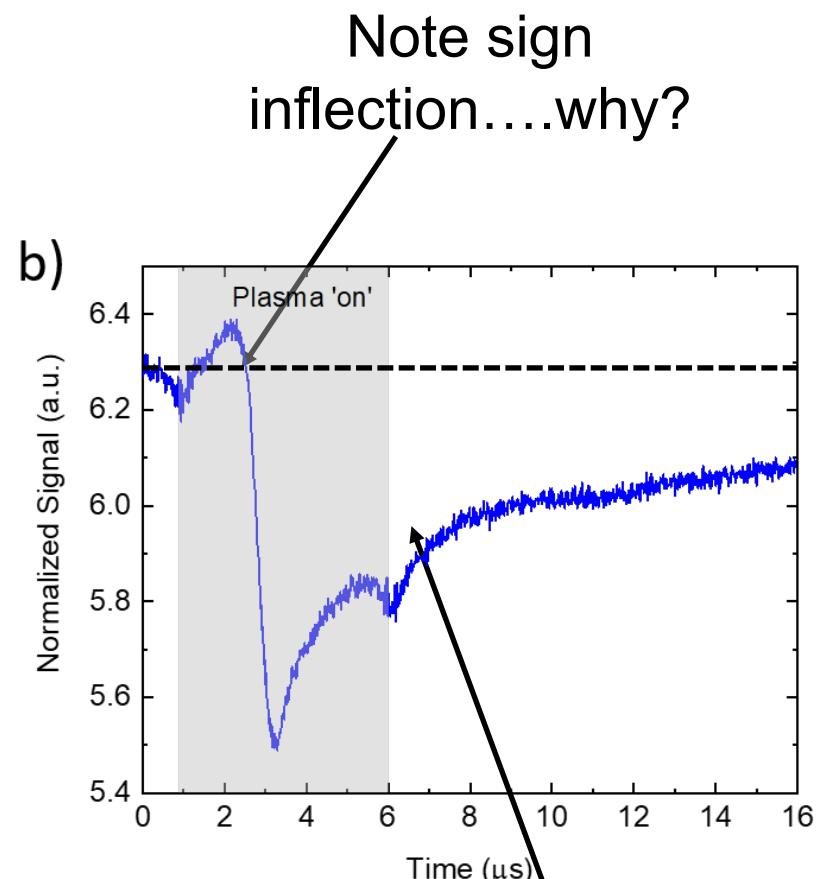
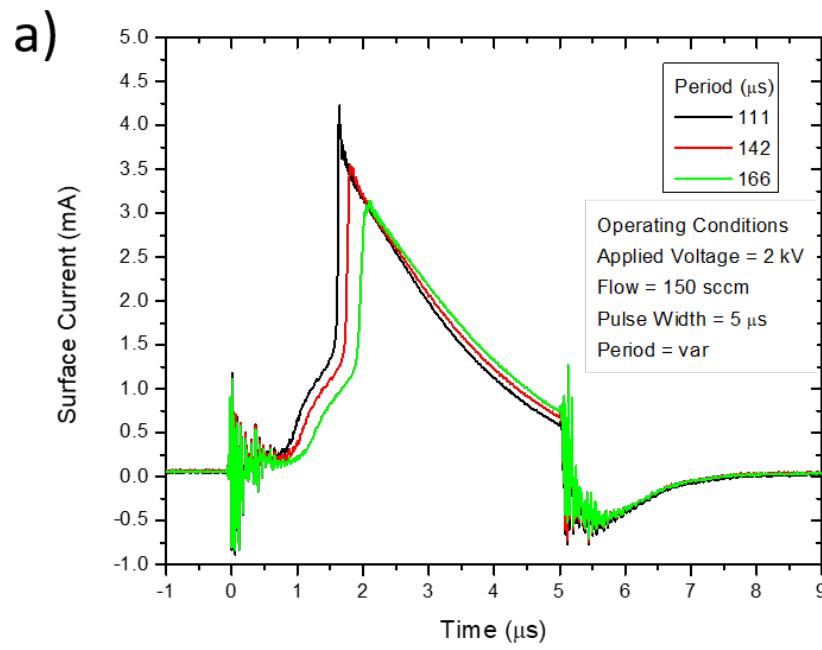
“Plasma-surface interactions in atmospheric pressure plasmas: *In situ* measurements of electron heating in materials”

J. Appl. Phys. **124**, 043301 (Editor's Pick)



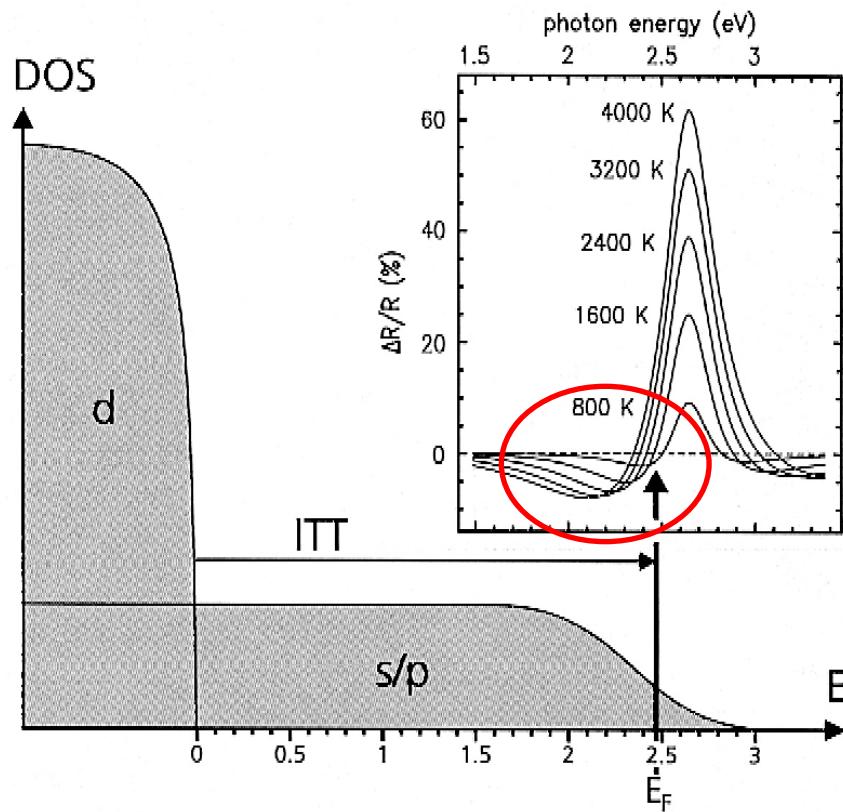
Tracking surface temperature during plasma jet irradiation

Surface current temporal profile
for different DC pulse widths



“Accommodation” from neutrals?

Thermoreflectance of gold



Chemical Physics 251 (2000) 237–258

www.elsevier.nl/locate/chemphys

Chemical
Physics

Electron and lattice dynamics following optical excitation of metals

J. Hohlfeld, S.-S. Wellershoff, J. Güdde, U. Conrad, V. Jähnke, E. Matthias

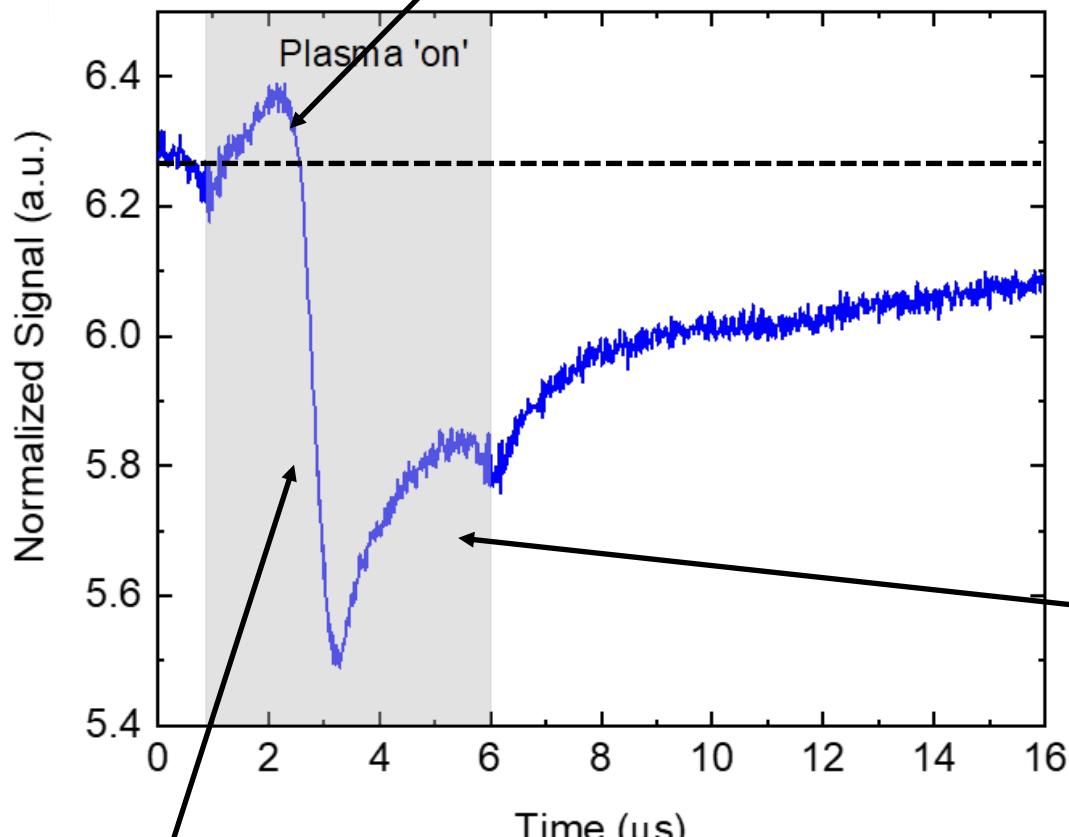
Fachbereich Physik, Freie Universität Berlin, Arnimallee 14, D-14195 Berlin, Germany

Received 13 May 1999

- Probe energies below d-band to Fermi level transition (ITB)
 - Increase ΔT
 - Decrease $\Delta R/R$

Plasma cooling of surface

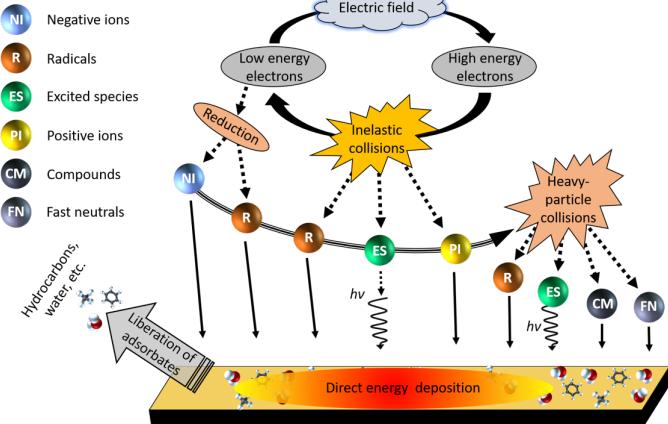
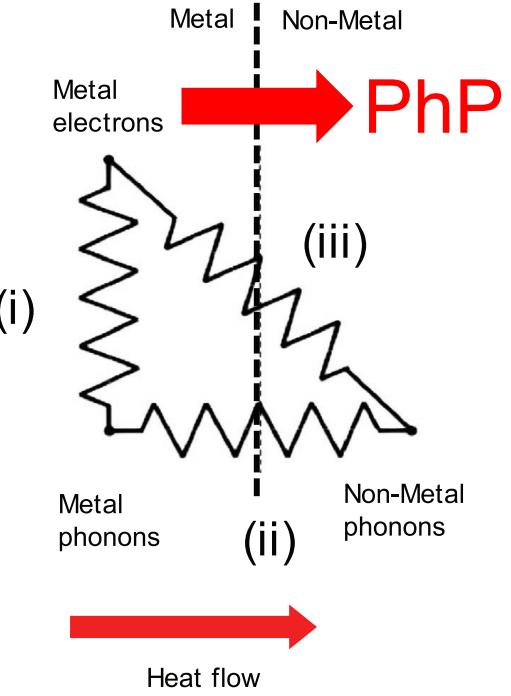
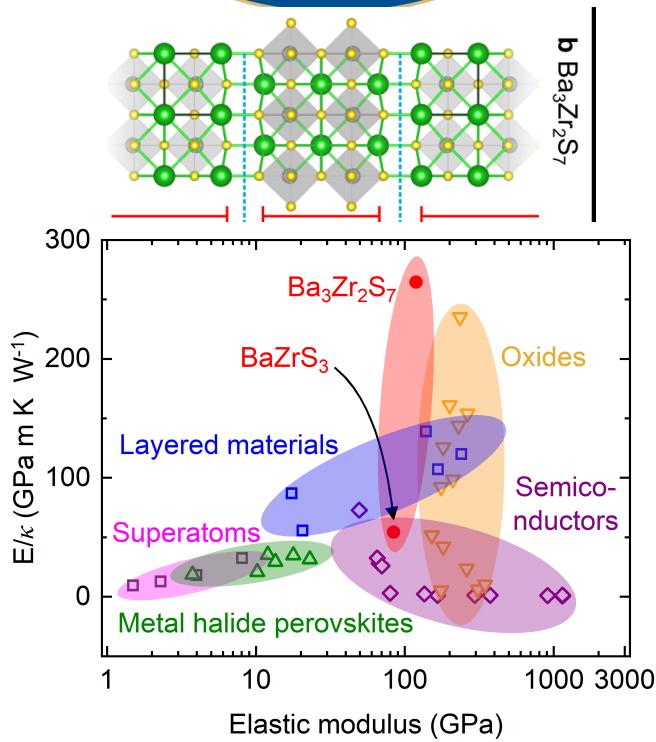
Plasma cooling: either mass removal or electron ejection from high eV photons (can not rule out either)



Cooling of surface from decrease in surface current or pulse turning off

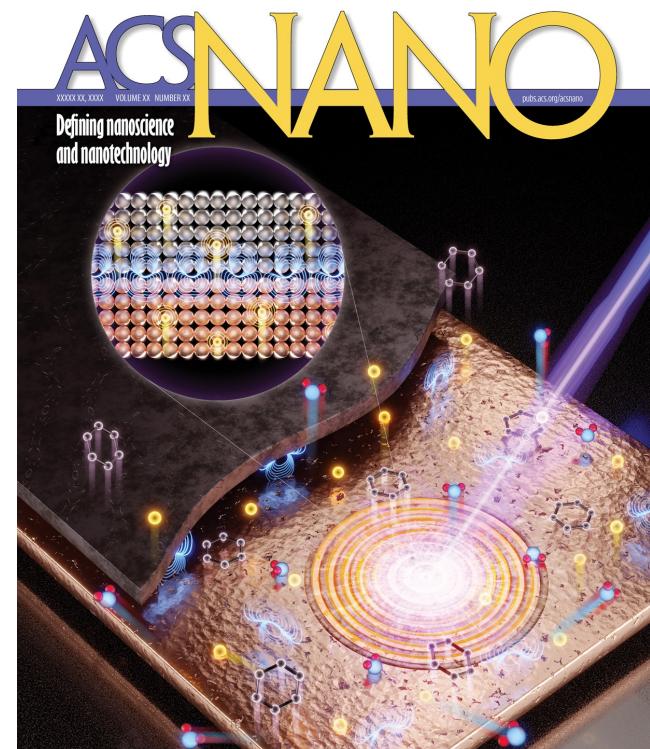
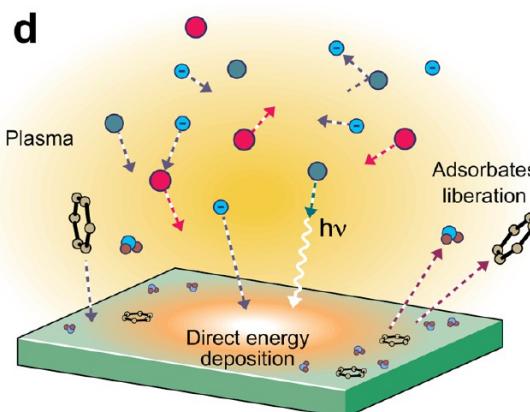
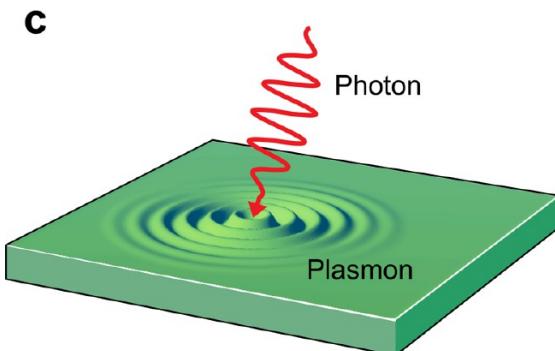
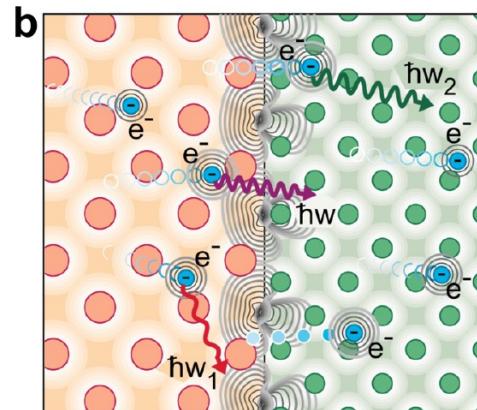
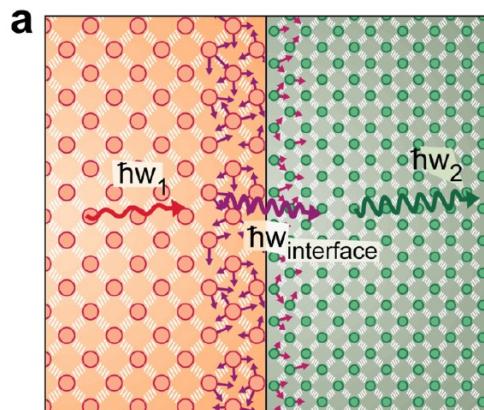
$-\Delta R$ = increase in temperature from energy transfer to gold

Theme: Coupling of carriers across interfaces lead to unique energy states and modes of heat transfer

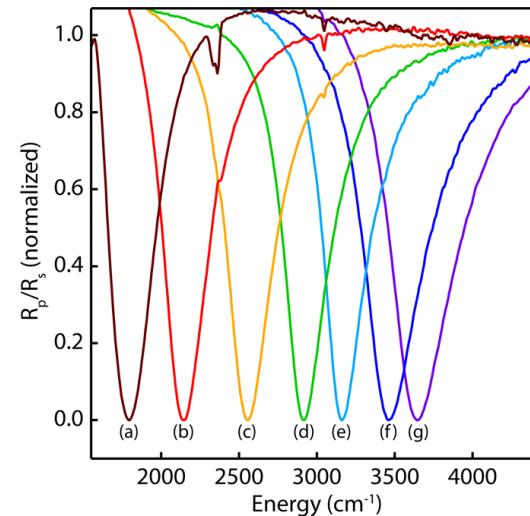
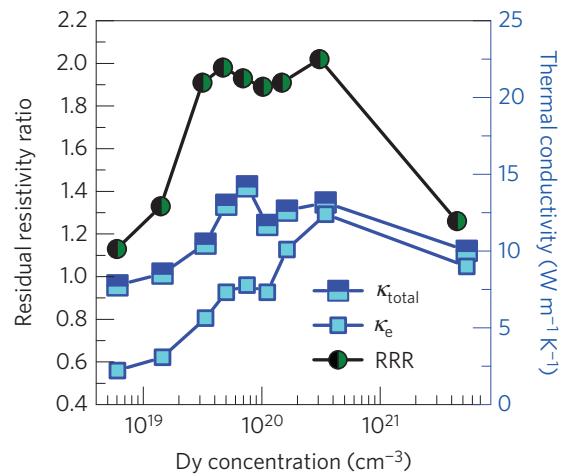
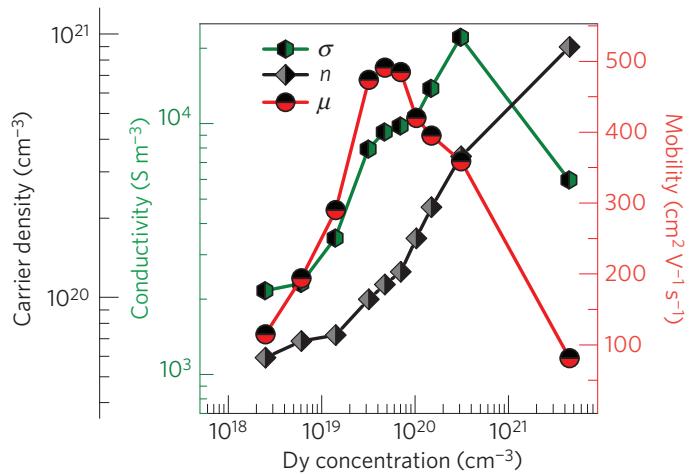


Ultrafast and Nanoscale Energy Transduction Mechanisms and Coupled Thermal Transport across Interfaces

Ashutosh Giri,* Scott G. Walton, John Tomko, Niraj Bhatt, Michael J. Johnson, David R. Boris, Guanyu Lu, Joshua D. Caldwell, Oleg V. Prezhdo, and Patrick E. Hopkins*



CdO – a gateway for mid-IR plasmonics



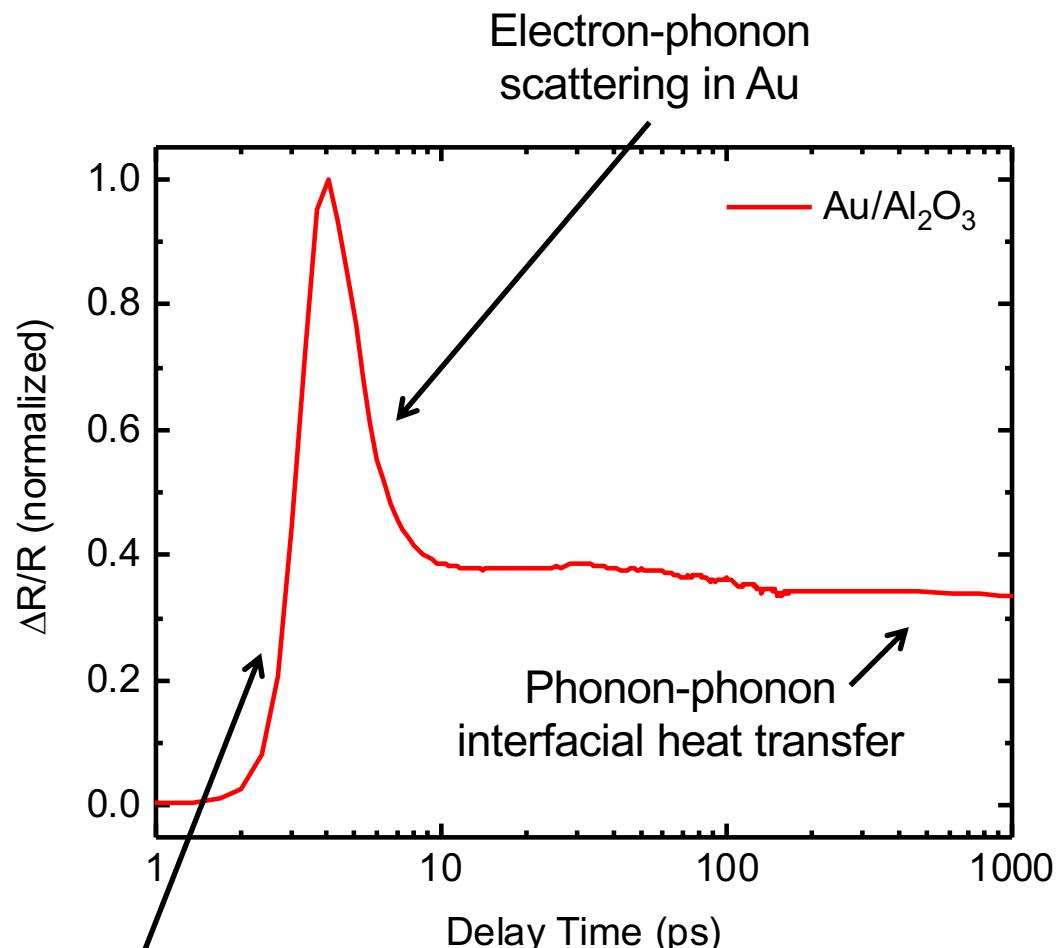
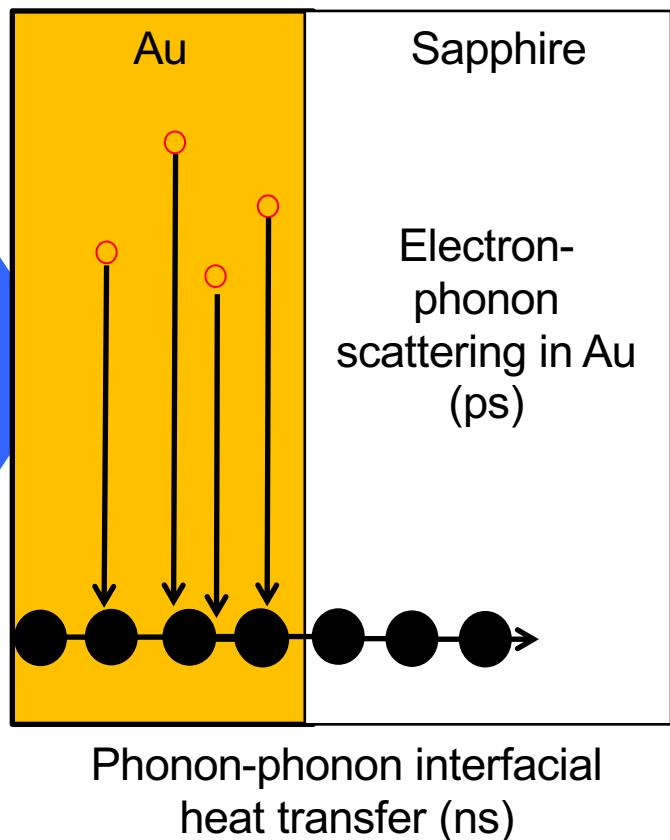
ACS Photonics 4, 1885

- Large electron mobility in CdO results in large electronic thermal conductivity
- Doping concentration tunes electronic conductivity and IR absorption

Nat. Mat. 14, 414

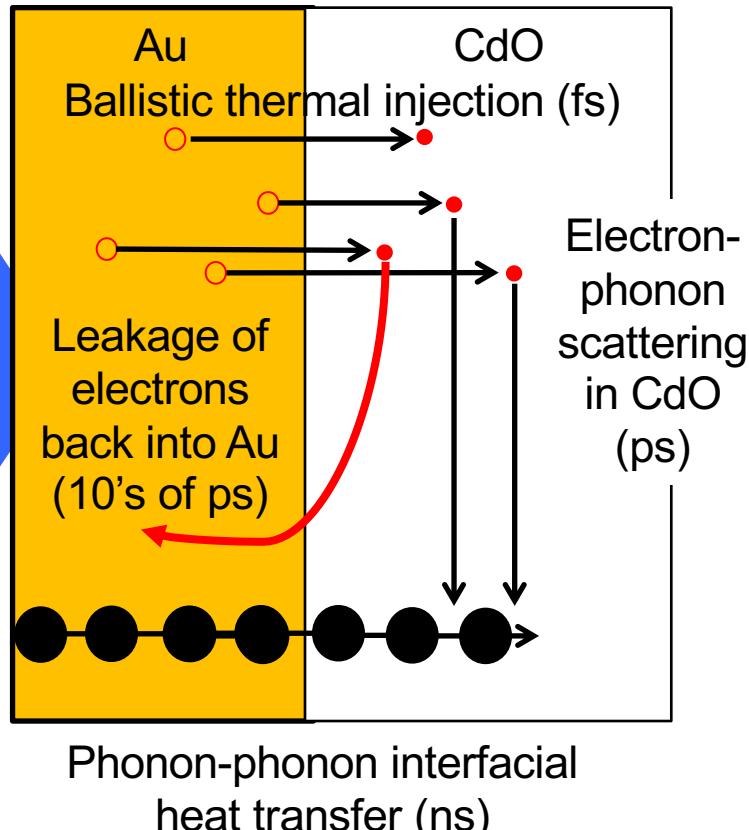
Appl. Phys. Lett. 108, 021901

Nonequilibrium processes at Au/CdO interfaces

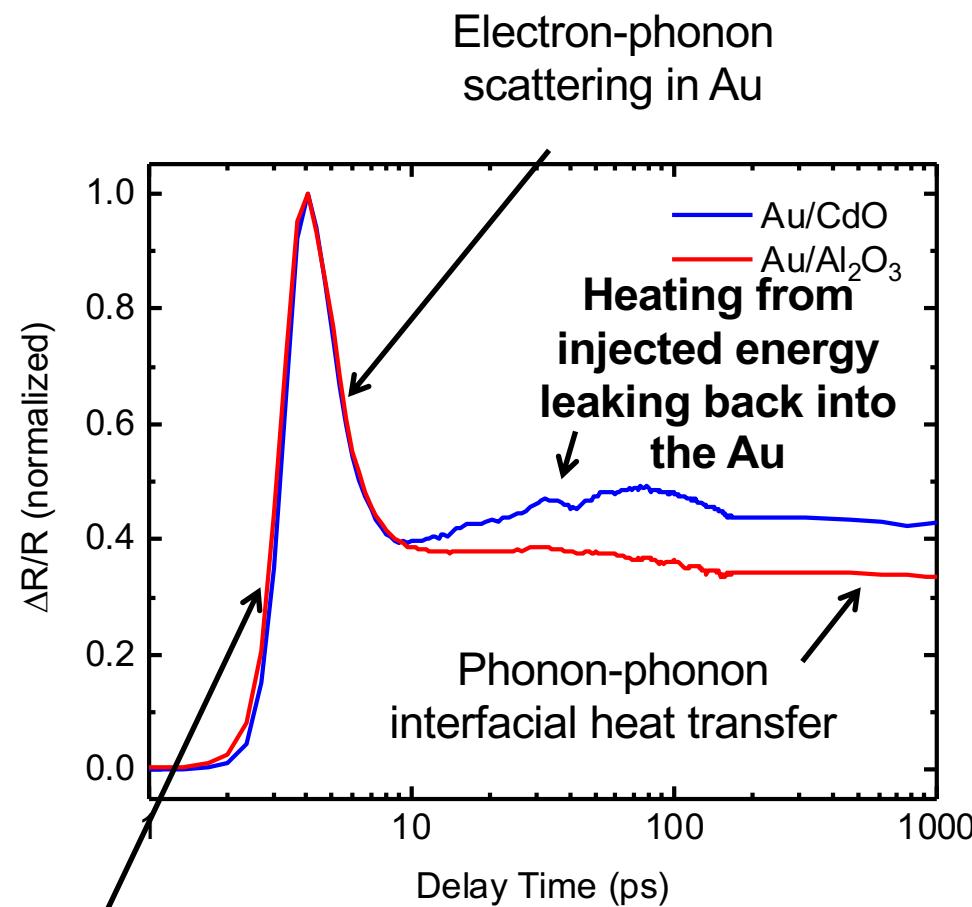


~sub picosecond pulsed laser absorption in Au

Nonequilibrium processes at Au/CdO interfaces

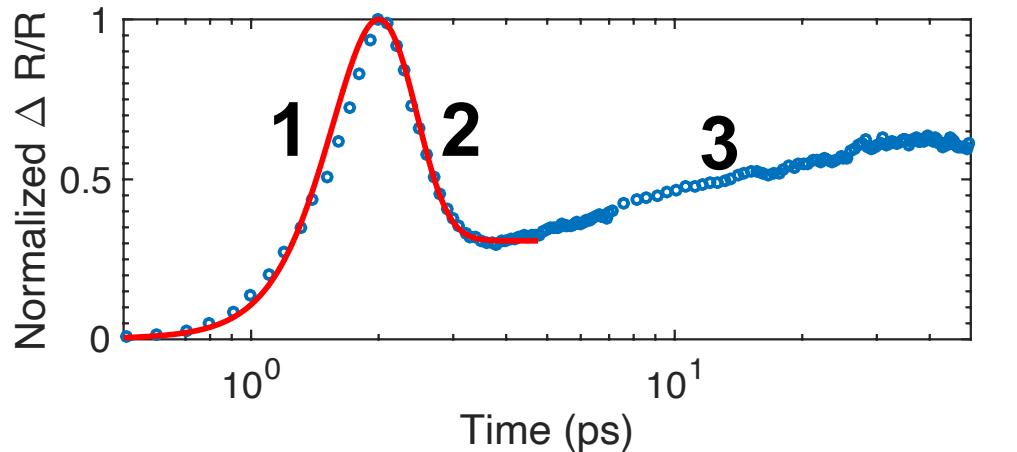


~sub picosecond pulsed laser absorption in Au + charge injection into CdO

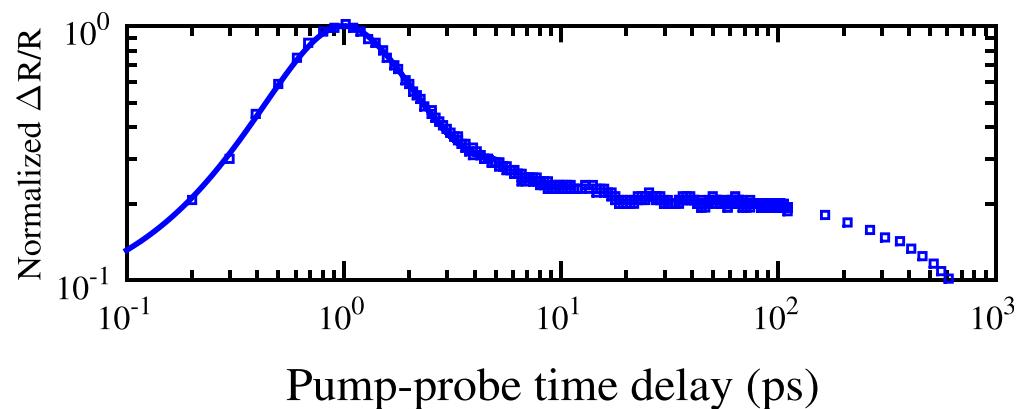
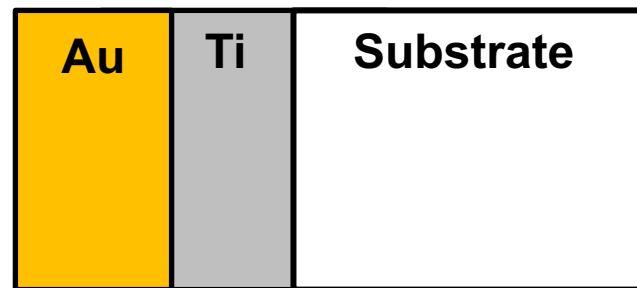


TDTR measurements of time scales of noneq. transport

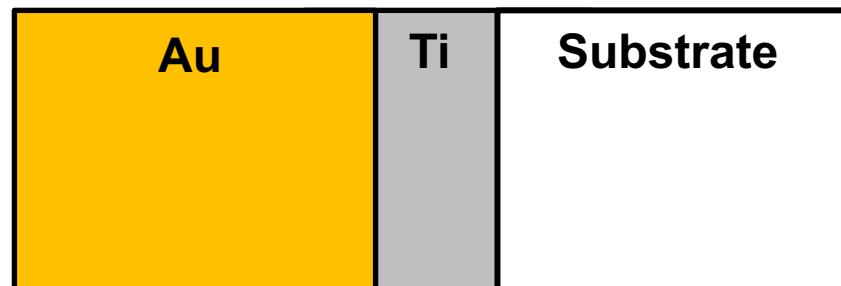
Hypothesis: If Au thickness (d_{Au}) is thicker than electron-phonon mean free path (λ_{ep}), nonequilibrium at interface will be negligible and “back heating” (time regime 3) will not be observed



$$d_{\text{Au}} < \lambda_{\text{ep}}$$



$$d_{\text{Au}} > \lambda_{\text{ep}}$$



Nonequilibrium processes at Au/CdO interfaces

Doping will control electron-electron TBC and electron thermal conductivity in CdO, vary “back heating”

