



**SCHOOL *of* ENGINEERING & APPLIED SCIENCE  
UNIVERSITY *of* VIRGINIA**

# Thermal transport processes in functional organic and organic/inorganic nanomaterials: The role of the chemical bond



**Patrick E. Hopkins**

Associate Professor

Dept. Mech. & Aero. Eng.

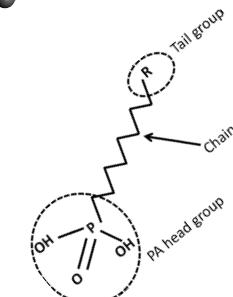
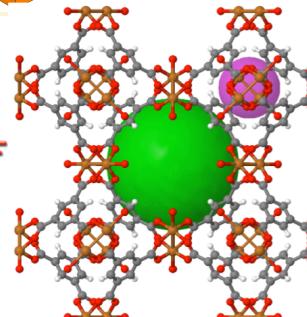
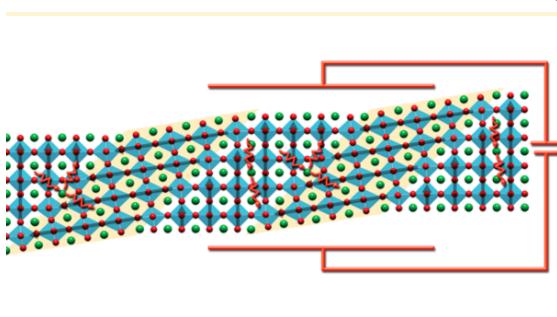
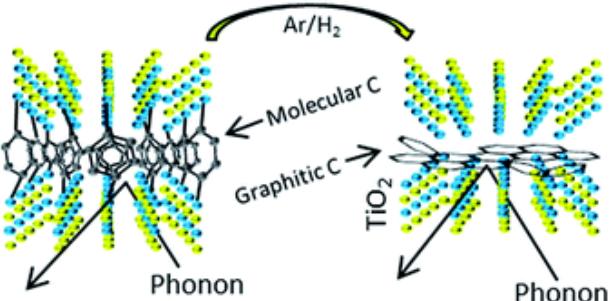
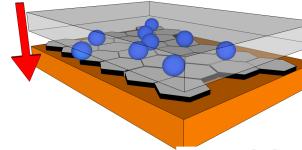
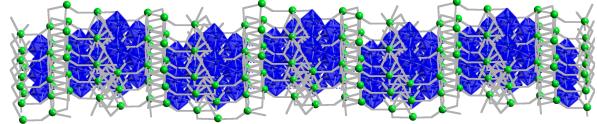
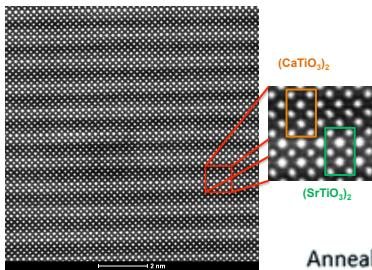
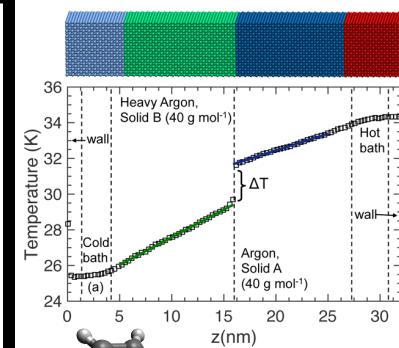
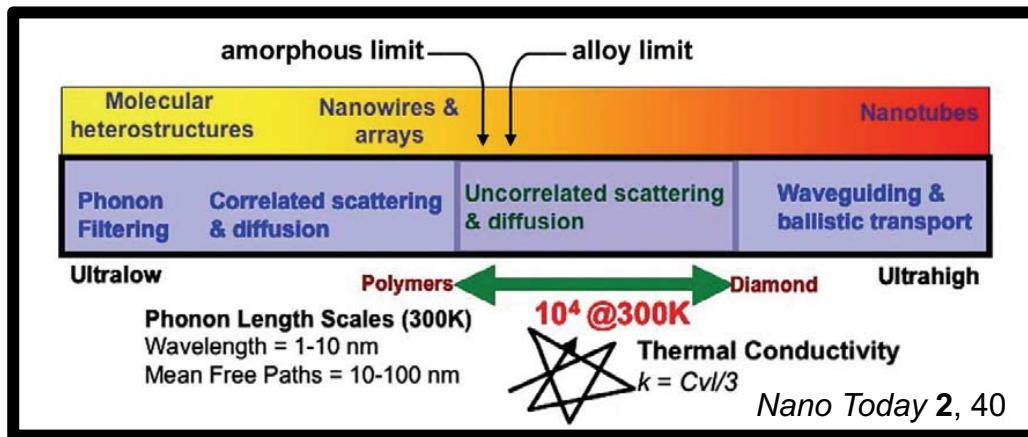
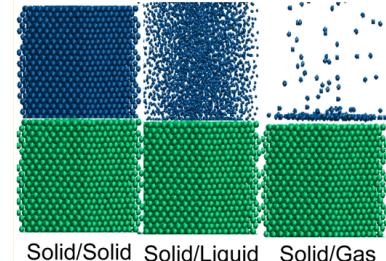
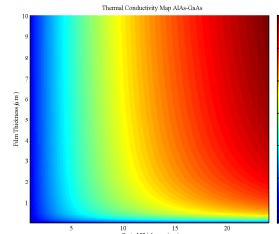
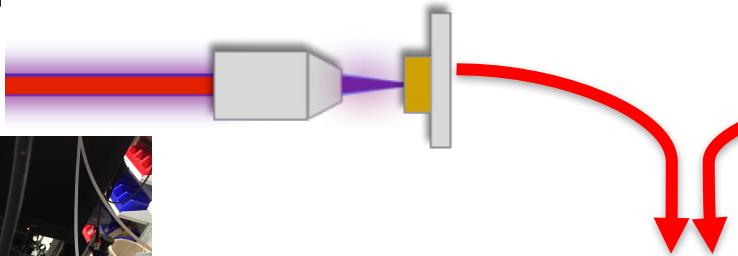
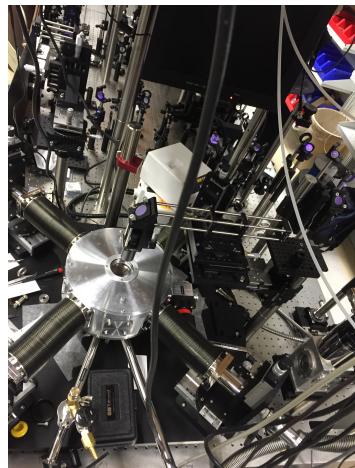
University of Virginia

[phopkins@virginia.edu](mailto:phopkins@virginia.edu)

[patrickehopkins.com](http://patrickehopkins.com)

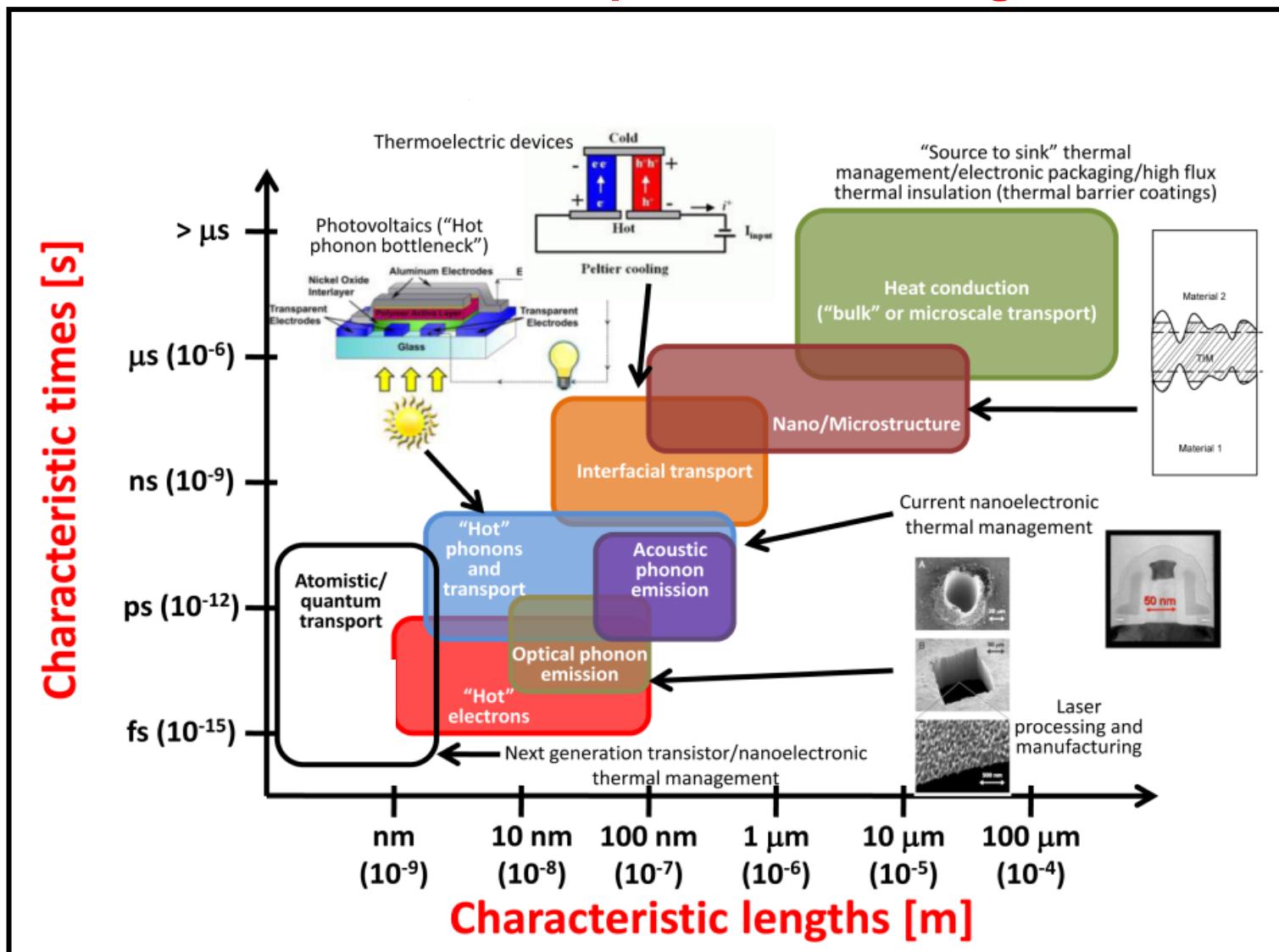
# My group – ExSiTE Lab

## Experiments and Simulations in Thermal Engineering

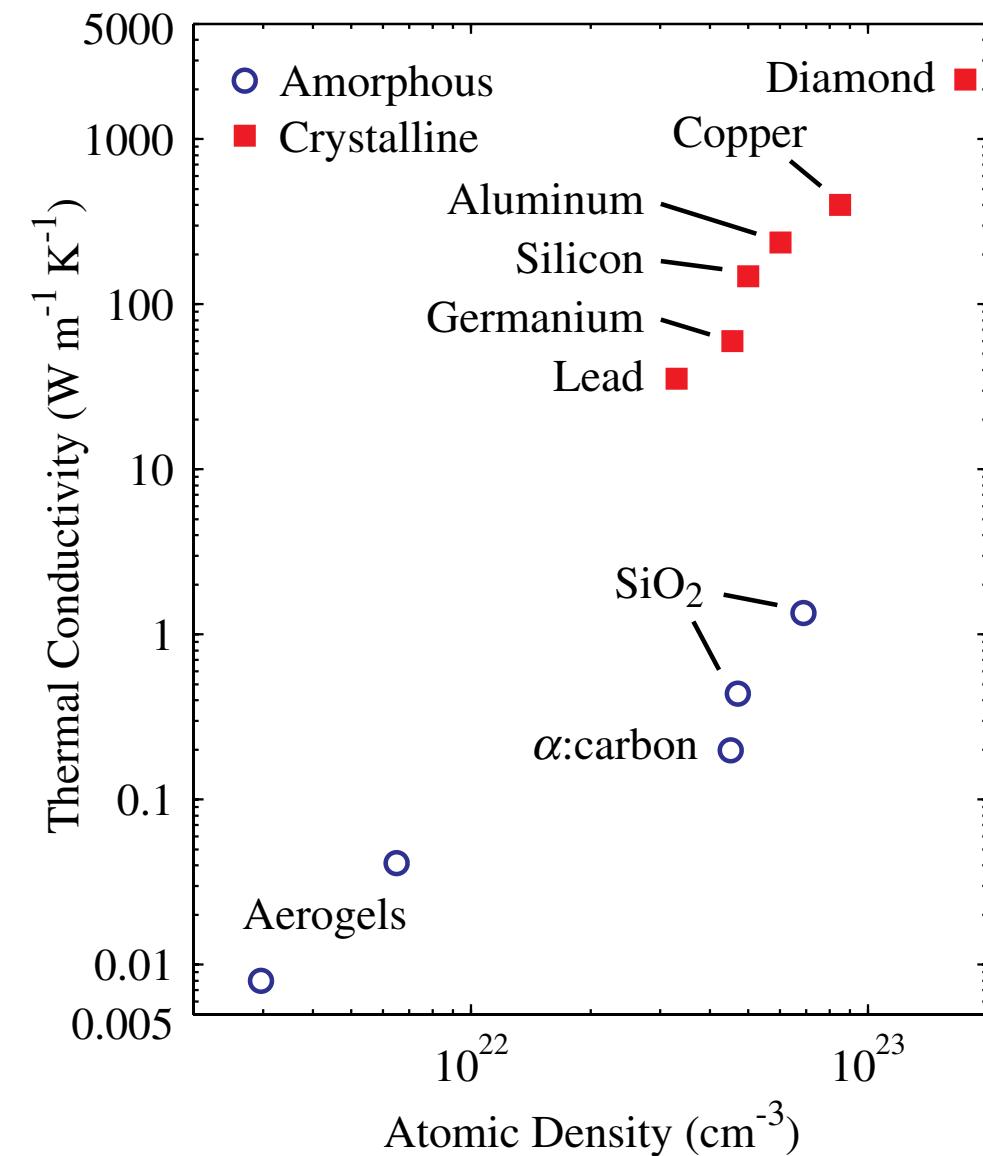


# My group – ExSiTE Lab

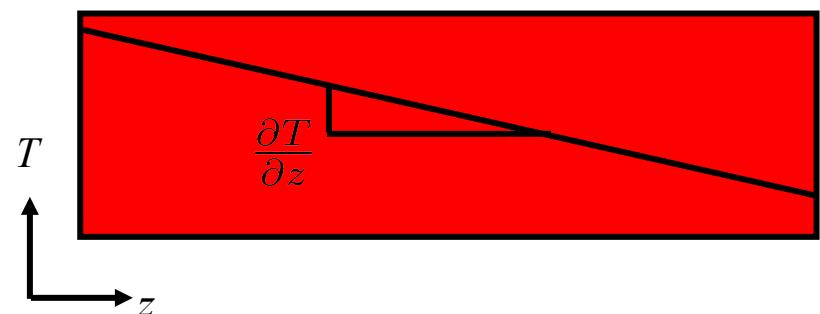
## Engineering energy transport, conversion and storage in materials over multiple time and length scales



# Thermal conductivity of materials – Traditional picture

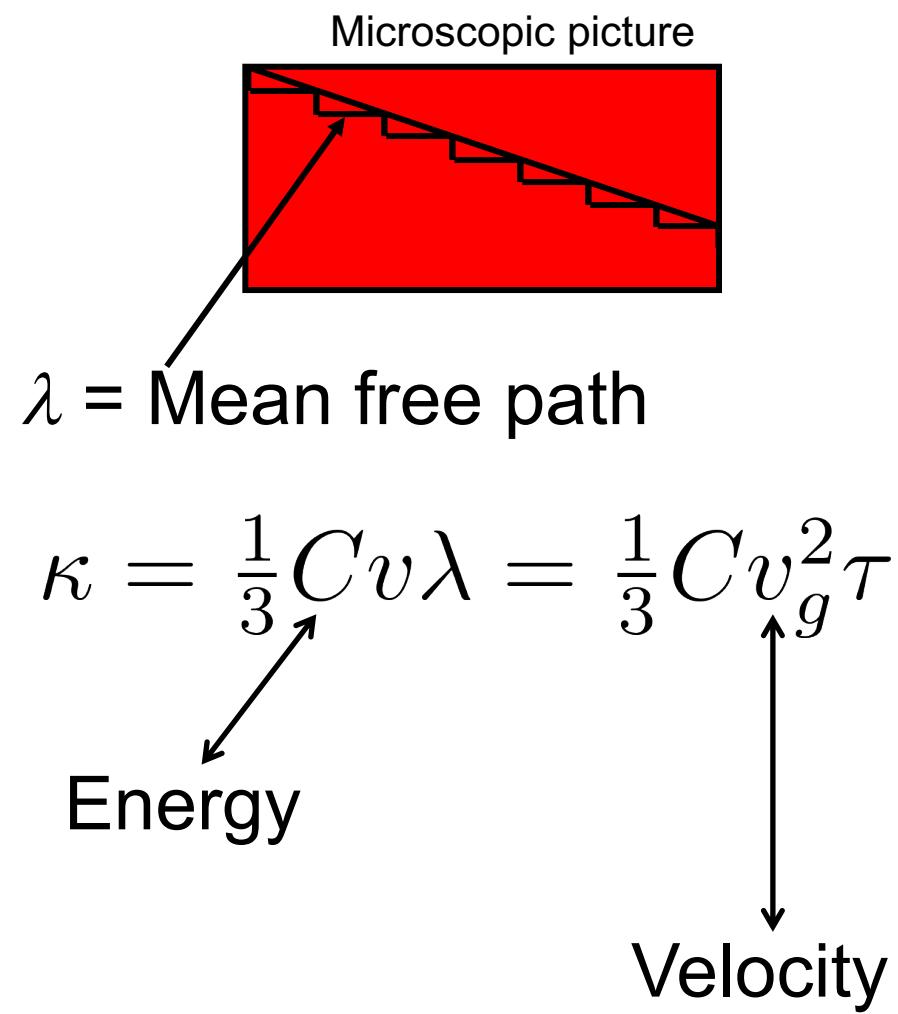
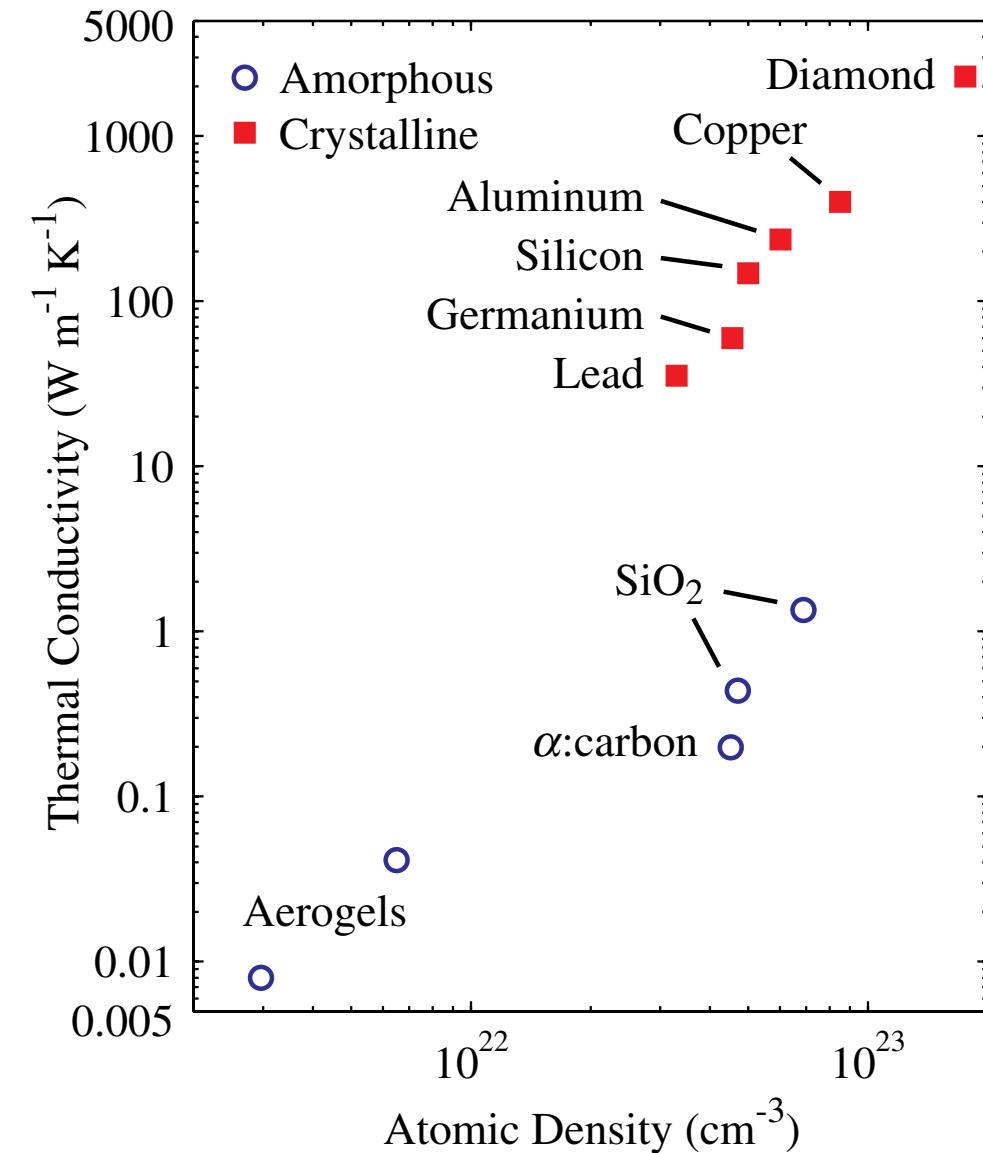


Bulk picture (Fourier Law)



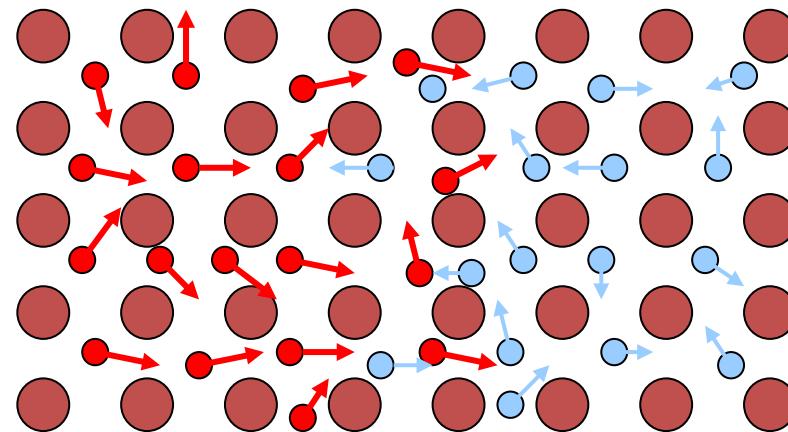
$$q_z = -\kappa \frac{\partial T}{\partial z}$$

# Thermal conductivity of materials – Kinetic Theory picture



# Thermal conductivity of materials - nanoscopic

Diffusion of “hot” electrons



Metals:

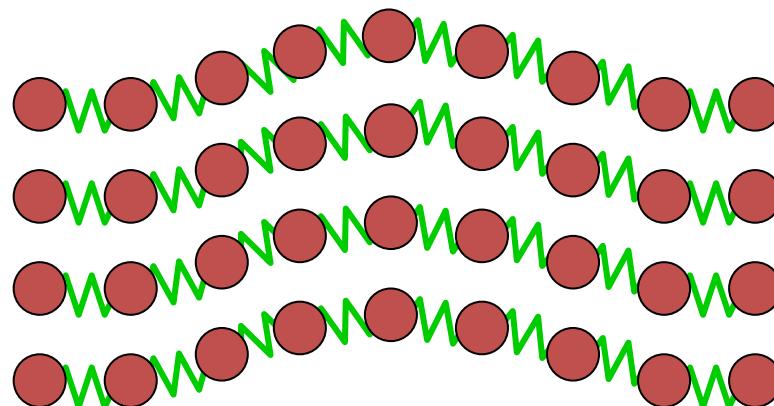
Free electrons are the dominant energy carriers in metals, ballistic velocity  $\sim 10^6$  m/s

- atom
- “hot” free electron
- “cold” free electron

Electron carrier density:

in metals  $\sim 10^{23}$  cm<sup>-3</sup>

in semiconductors  $\sim 10^{18}$  cm<sup>-3</sup>



Phonon propagation

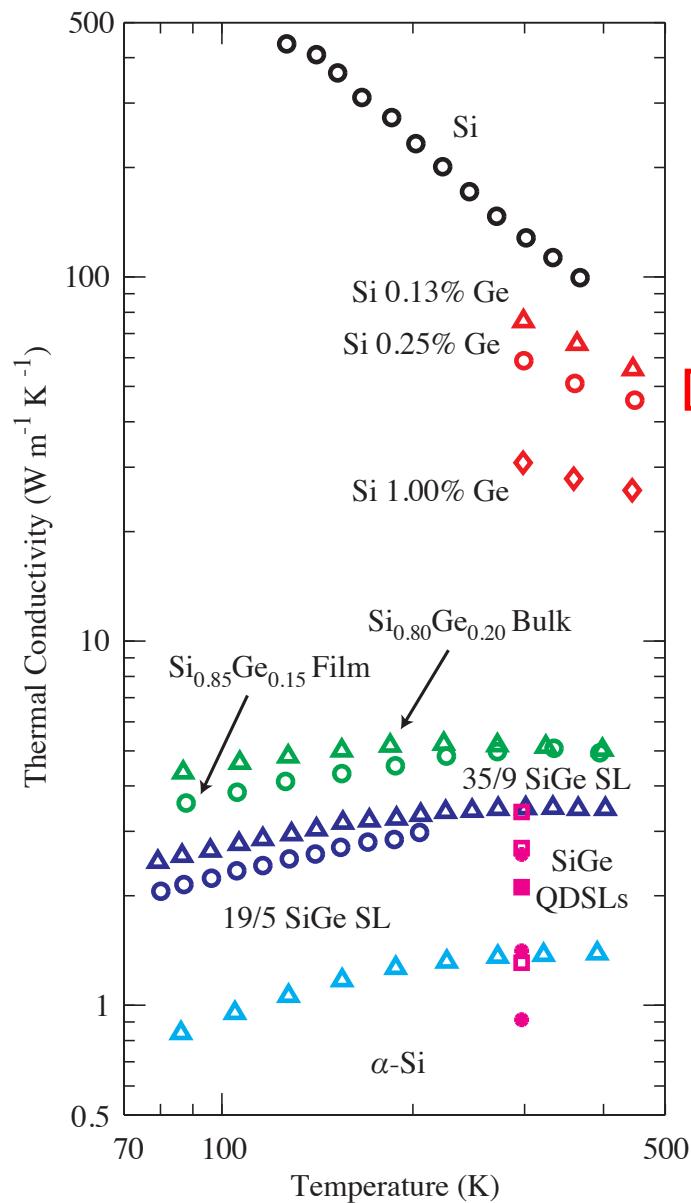


Semiconductors:

Phonons (lattice vibrations) are the dominant energy carriers in semiconductors, velocity  $\sim 10^3$  m/s

$$\kappa = \frac{1}{3} Cv\lambda = \frac{1}{3} Cv_g^2 \tau$$

# Thermal conductivity of materials - nanoscopic

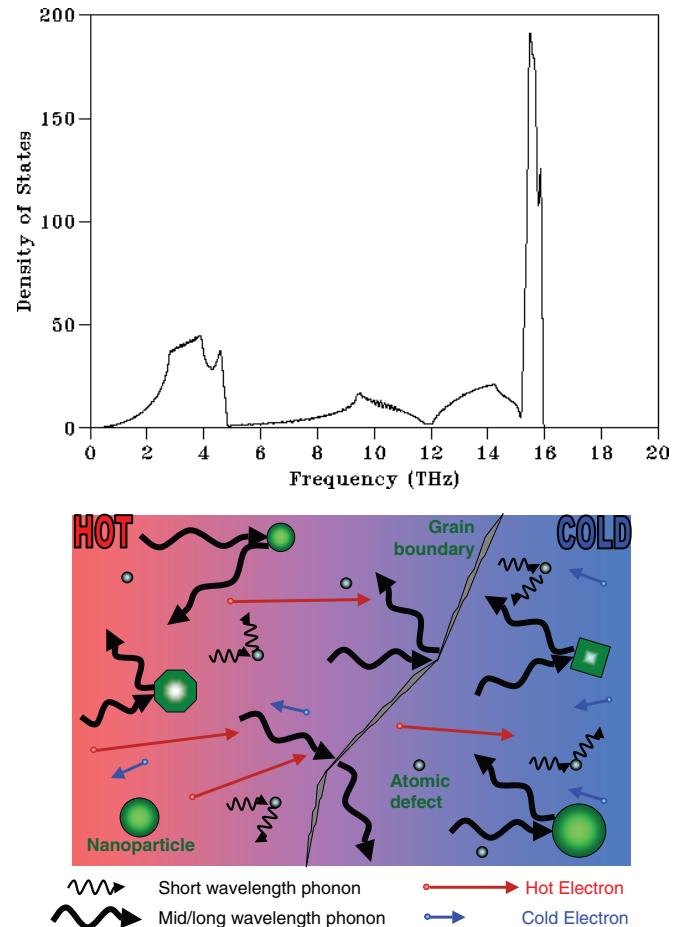


$$\kappa = \frac{1}{3} Cv\lambda = \frac{1}{3} Cv_g^2 \tau$$

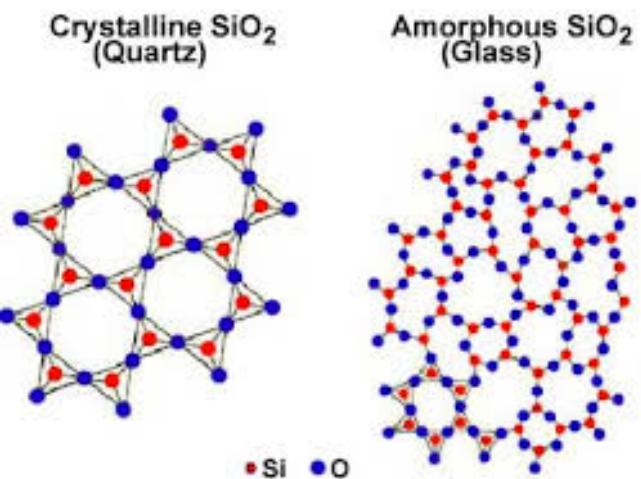
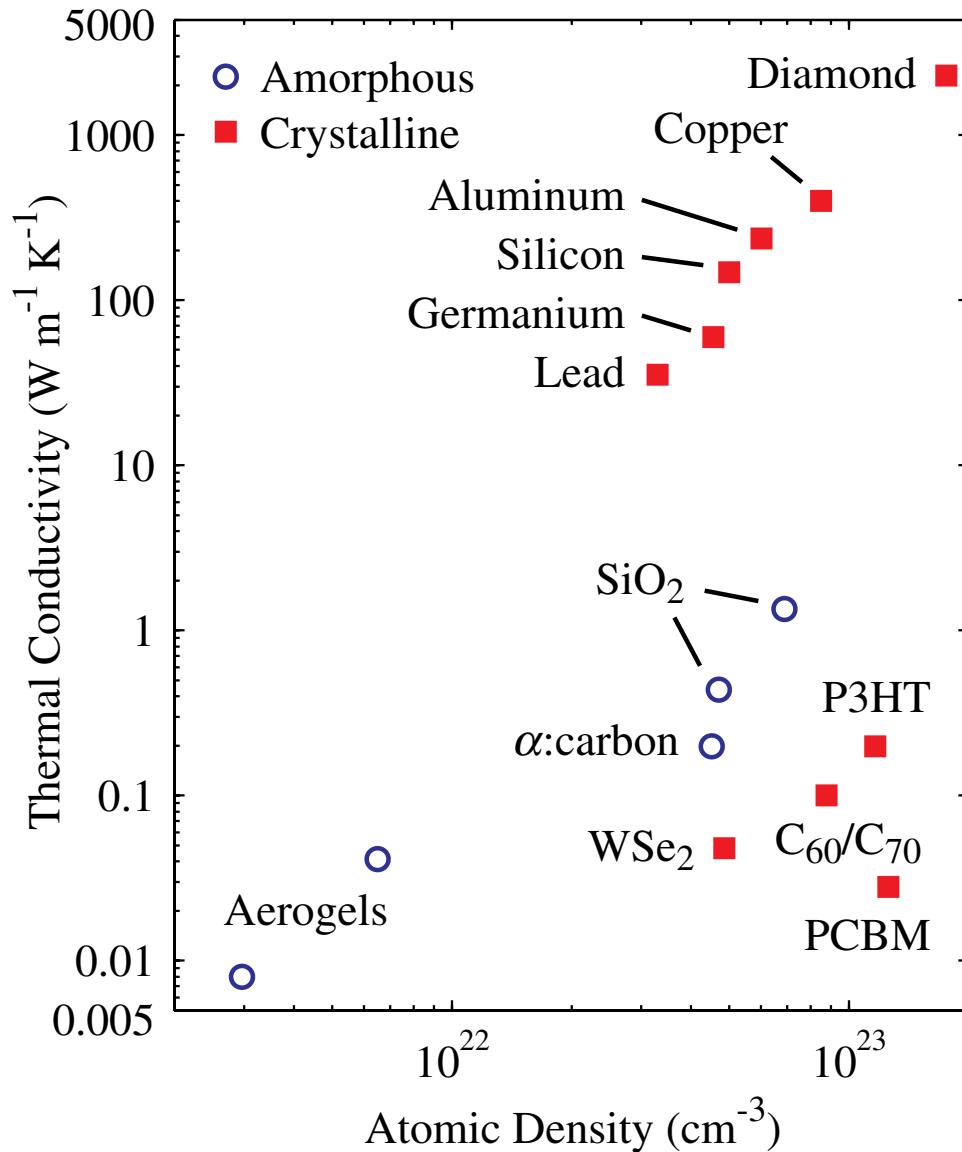
$$\tau = f(\tau_{\text{intrinsic}}, \tau_{\text{impurity}}, \tau_{\text{boundary}})$$

Dilute alloy

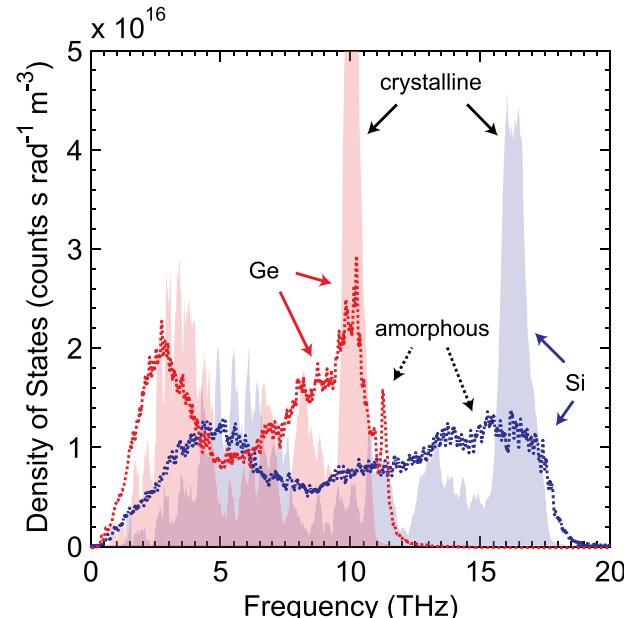
Alloy  
Interface  
Disorder



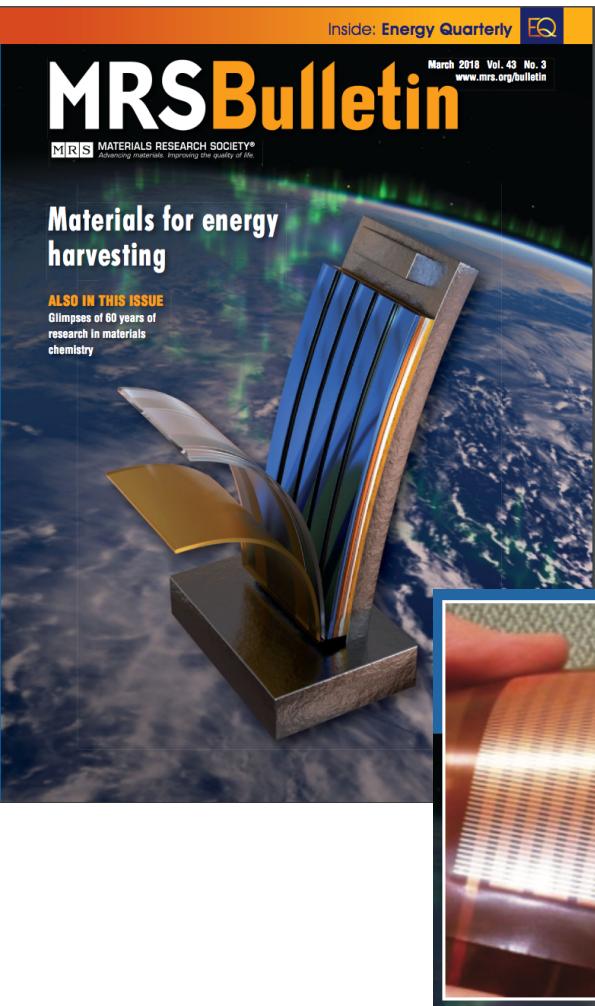
# Thermal conductivity of non-crystalline solids



*NDT Resource Center*



# A future for energy recovery: wearable thermoelectrics



March 2018 MRS Bulletin

Figure of merit  
[unitless]

## Wearable and flexible thermoelectrics for energy harvesting

Ruoming Tian, Chunlei Wan, Naoyuki Hayashi, Toshiaki Aoai,  
and Kunihito Koumoto

$$ZT = \frac{S^2 \sigma T}{\kappa}$$

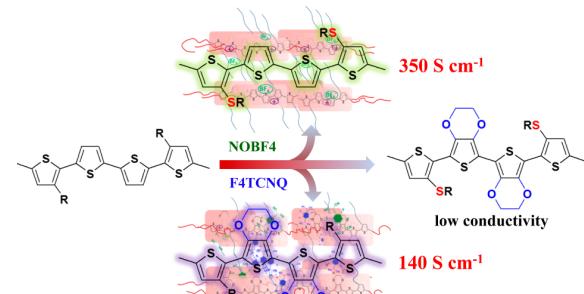
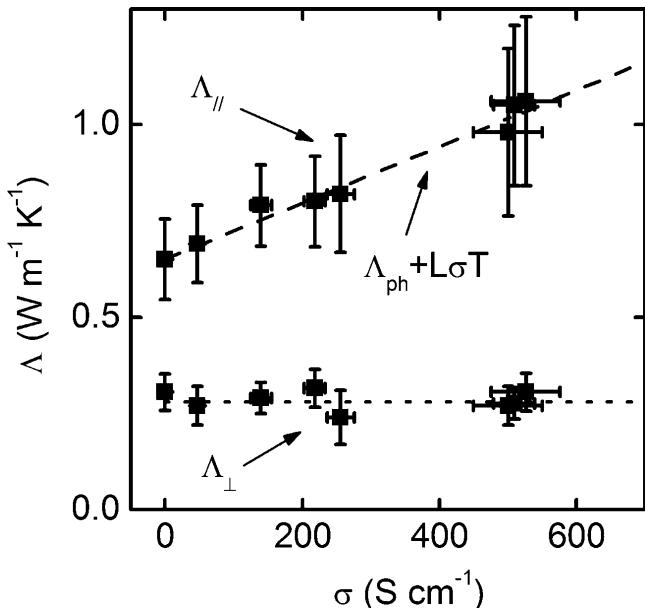
Diagram illustrating the Figure of Merit (ZT) for a thermoelectric material. The equation  $ZT = \frac{S^2 \sigma T}{\kappa}$  is shown, where  $S$  is the Seebeck coefficient [ $\text{V K}^{-1}$ ],  $\sigma$  is the Electrical conductivity [ $\text{S m}^{-1}$ ],  $T$  is Temperature [ $\text{K}$ ], and  $\kappa$  is the Thermal conductivity [ $\text{W m}^{-1} \text{K}^{-1}$ ]. The diagram (a) shows a cross-section of a thermoelectric device in 'Power generation mode'. It consists of a 'Heat source' (P-N junction) with a 'Heat sink' at the bottom. Red wavy arrows indicate heat flow ( $Q$ ) from the heat source to the heat sink. Positive charges ( $+$ ) are shown moving from the P-type side to the N-type side, and negative charges ( $-$ ) are shown moving from the N-type side to the P-type side. A lightbulb at the bottom indicates current flow ( $I$ ).

renewableenergy  
focus.com

# Engineering thermal conductivity in polymeric TE materials

$$ZT = \frac{S^2 \sigma T}{\kappa}$$

Macromolecules 48, 585

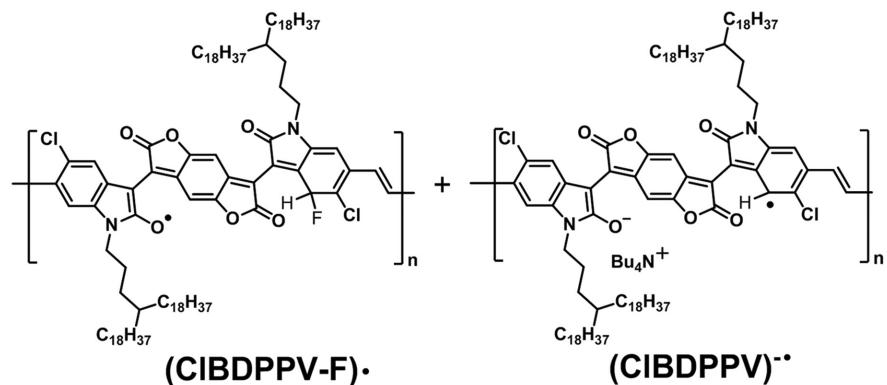


J | A | C | S  
JOURNAL OF THE AMERICAN CHEMICAL SOCIETY

Article  
pubs.acs.org/JACS

## Modification of the Poly(bisbiododecylquaterthiophene) Structure for High and Predominantly Nonionic Conductivity with Matched Dopants

Hui Li,<sup>†</sup> Mallory E. DeCoster,<sup>‡</sup> Robert M. Ireland,<sup>†</sup> Jian Song,<sup>†</sup> Patrick E. Hopkins,<sup>‡</sup> and Howard E. Katz<sup>\*,†</sup>



## COMMUNICATION

Thermoelectrics

ADVANCED MATERIALS  
www.advmat.de

## High Conductivity and Electron-Transfer Validation in an n-Type Fluoride-Anion-Doped Polymer for Thermoelectrics in Air

Xingang Zhao, Deepa Madan,\* Yan Cheng, Jiawang Zhou, Hui Li, Susanna M. Thon, Arthur E. Bragg, Mallory E. DeCoster, Patrick E. Hopkins, and Howard E. Katz\*

# Outline

- **TDTR: Measurement of thermal conductivity of thin films and thermal resistance across interfaces**
- Weakly bonded solids: new lower limits to thermal conductivity
- Functionalized interfaces at graphene contacts: tuning heat and electrical transport via the interfacial bond
- Heat transport across single molecule interfaces: when does a molecule become a defect?
- Molecular interfaces in organic/inorganic composites: diffusive scattering via the vibron-phonon interaction

# Measuring heat flow in thin films: TDTR

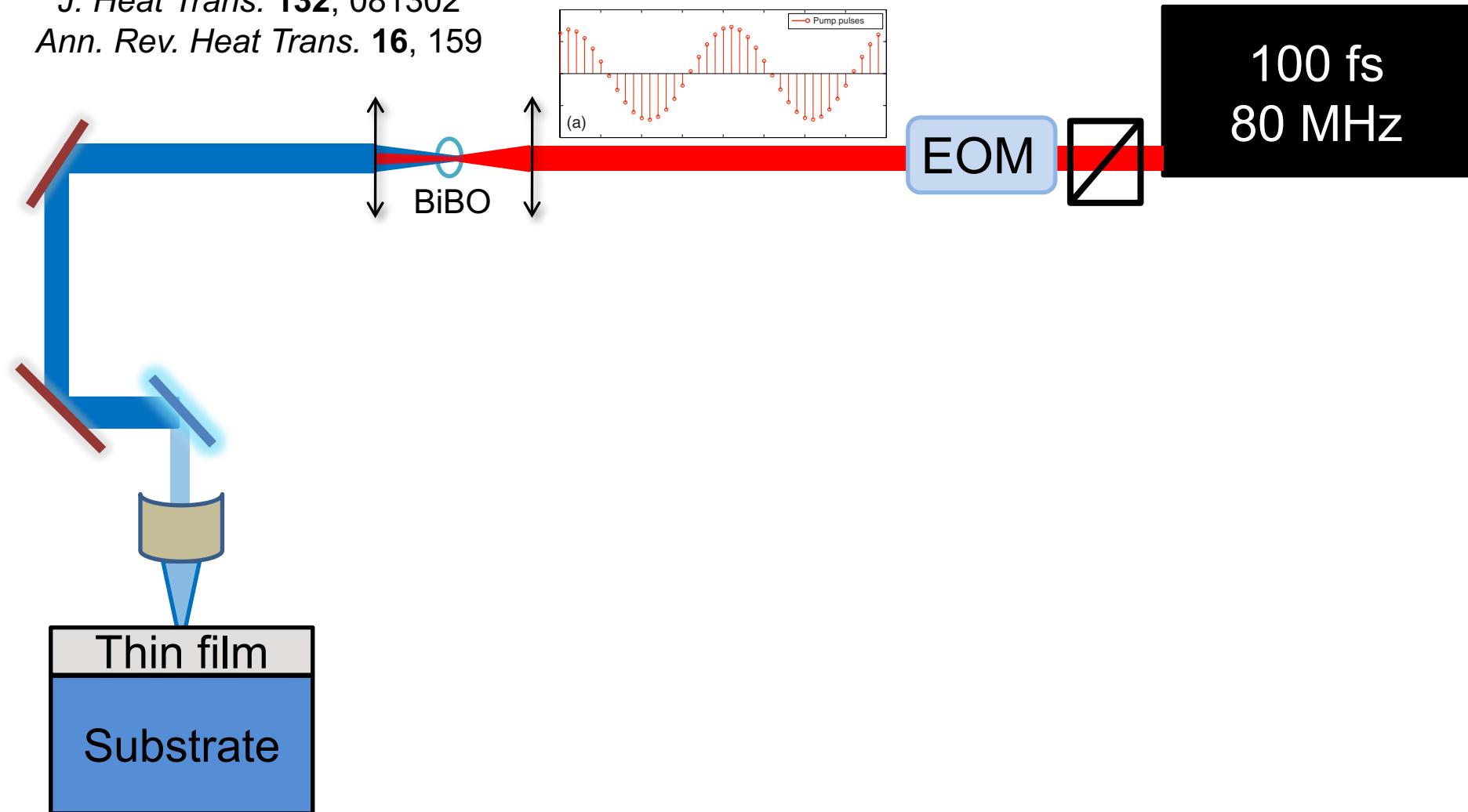
## TDTR Reviews and Analyses

*Rev. Sci. Instr.* **75**, 5119

*Rev. Sci. Instr.* **79**, 114902

*J. Heat Trans.* **132**, 081302

*Ann. Rev. Heat Trans.* **16**, 159



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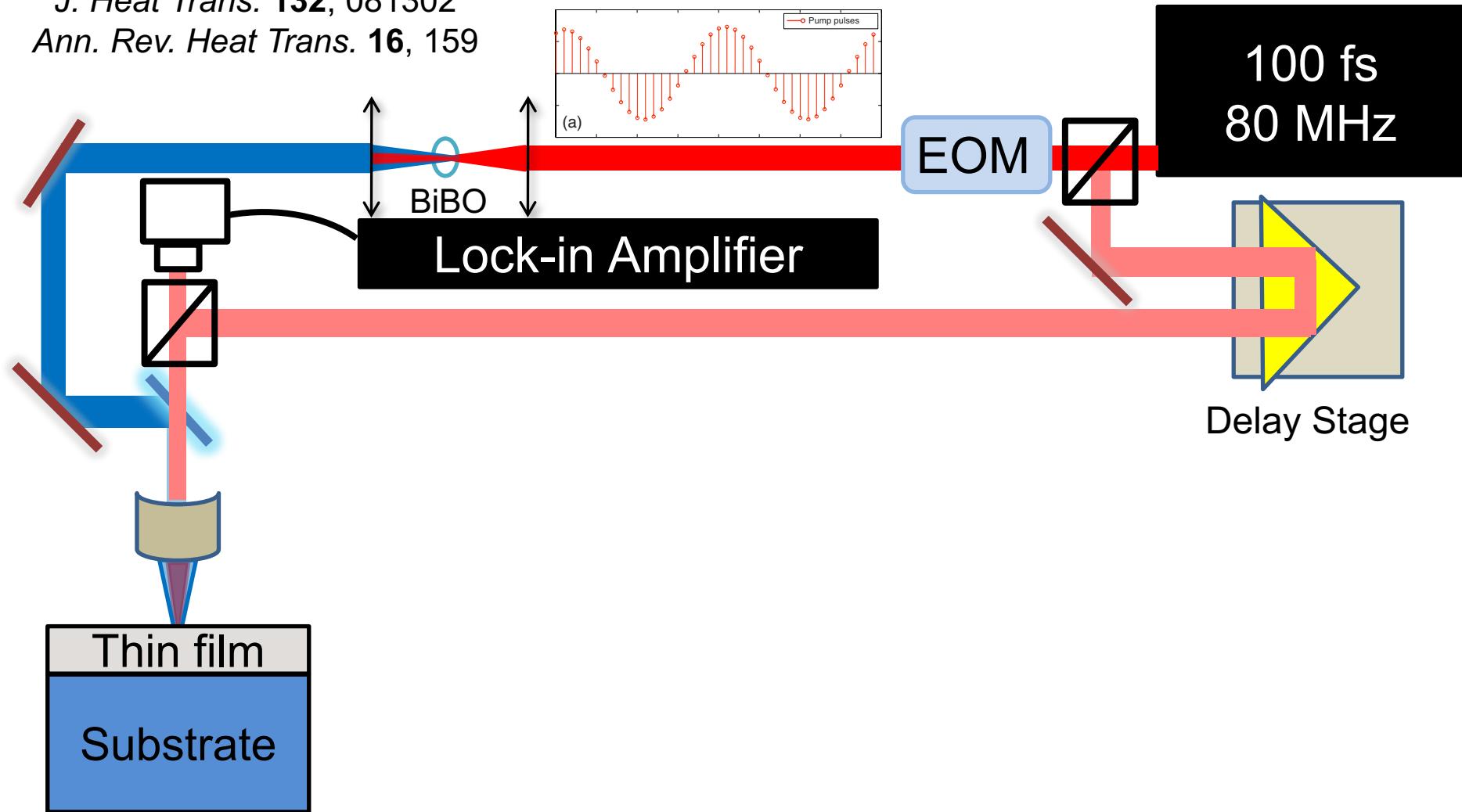
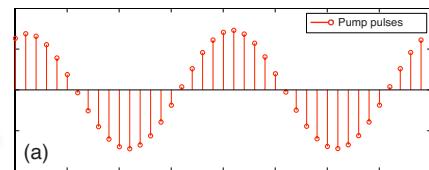
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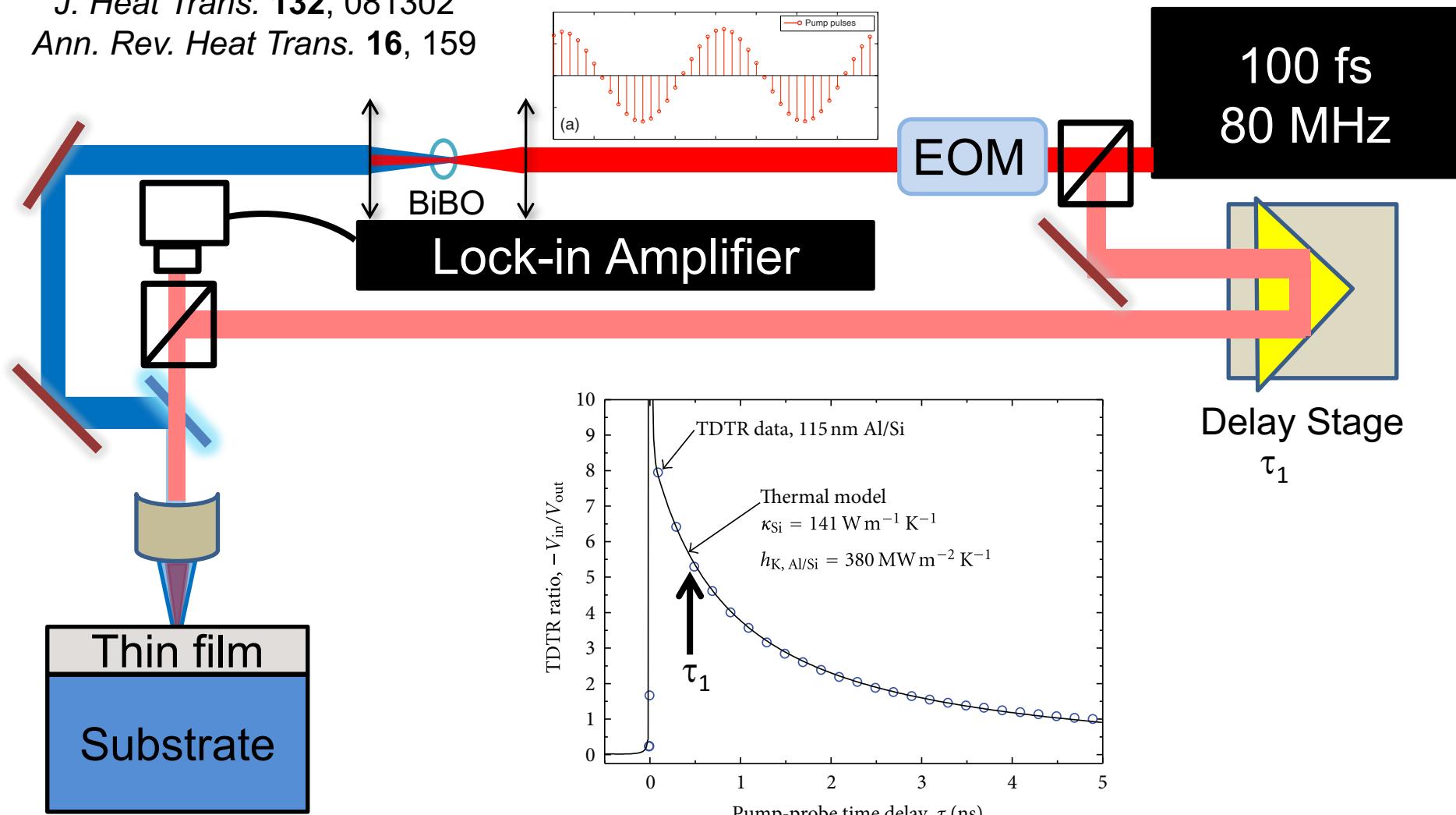
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*Rev. Sci. Instr.* **75**, 5119

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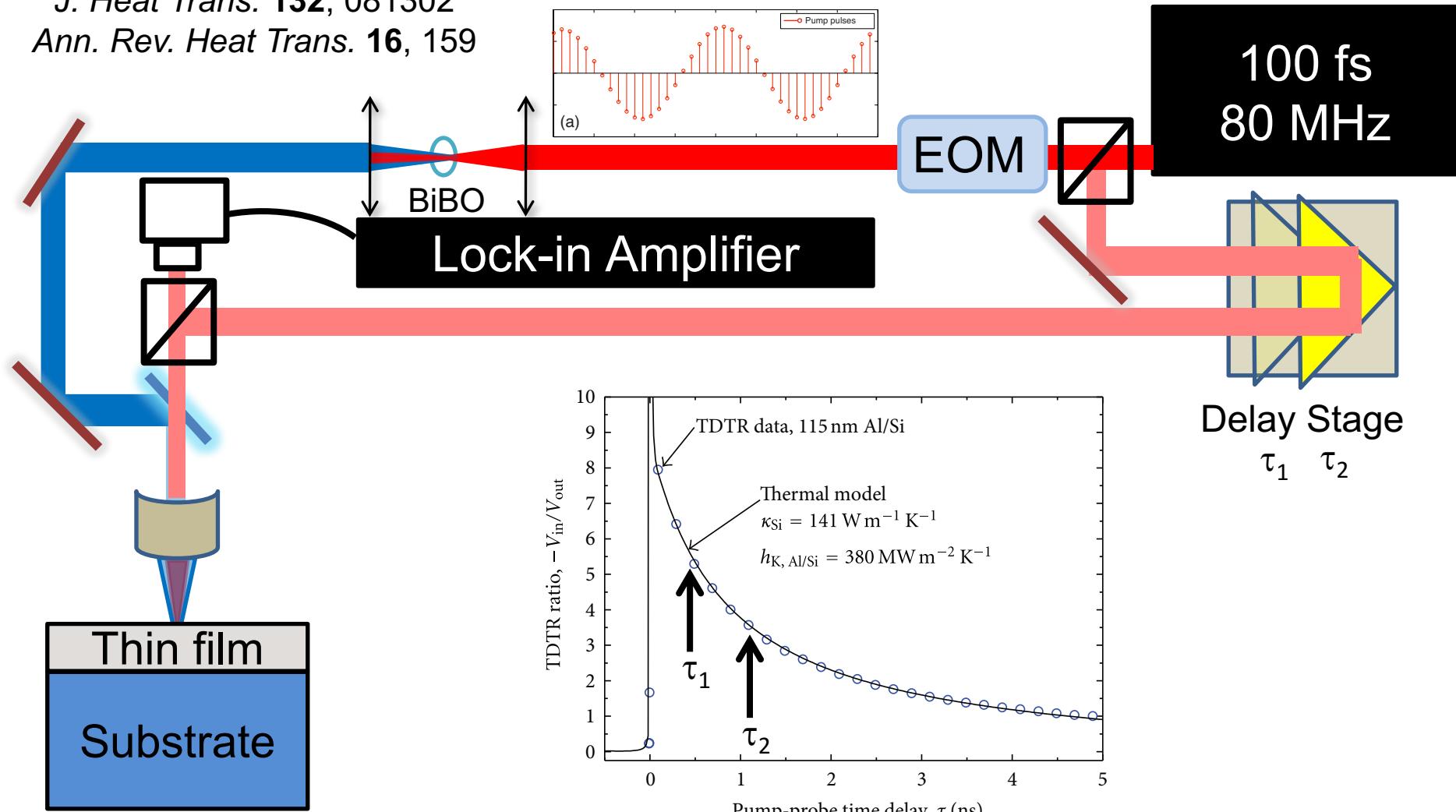
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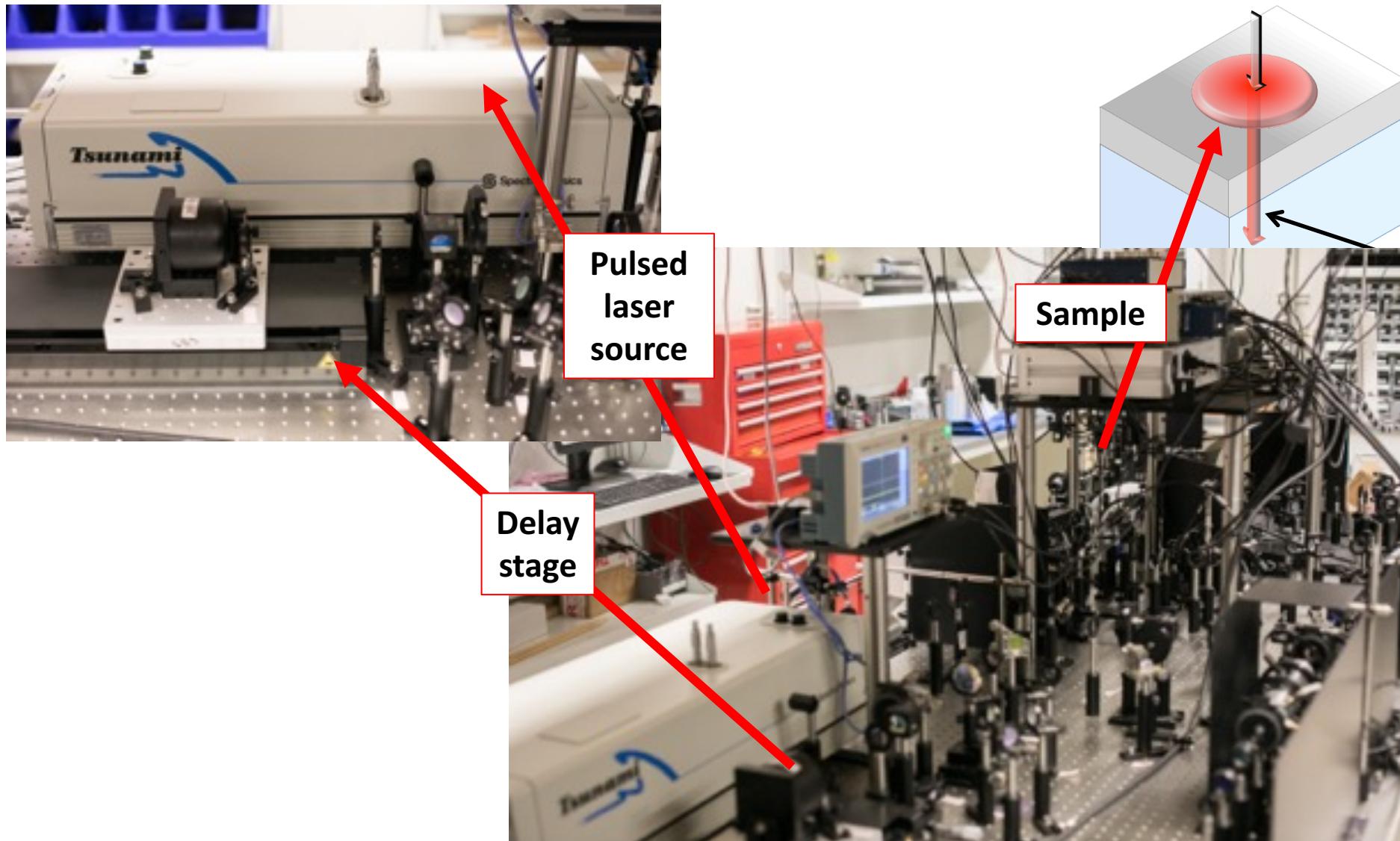
*Rev. Sci. Instr.* **79**, 114902

*J. Heat Trans.* **132**, 081302

*Ann. Rev. Heat Trans.* **16**, 159

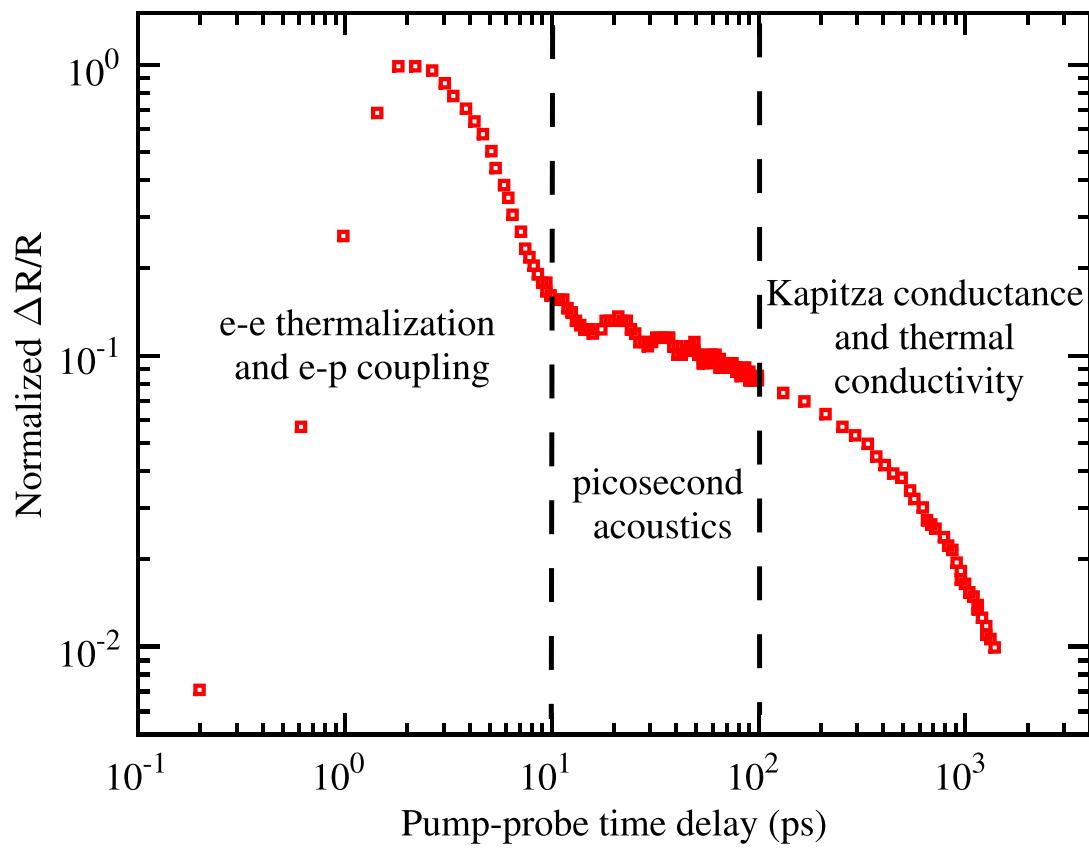


# Measuring heat flow in thin films: TDTR



# How can we measure nanoscale heat transport processes?

Need time scale resolution < picoseconds



Pulse absorption ( $\sim 100$  fs)

↓

Fermi relaxation and ballistic transport (few hundred fs)

↓

Electron-phonon coupling (a few ps)

↓

Strain propagation in film (10's of ps)

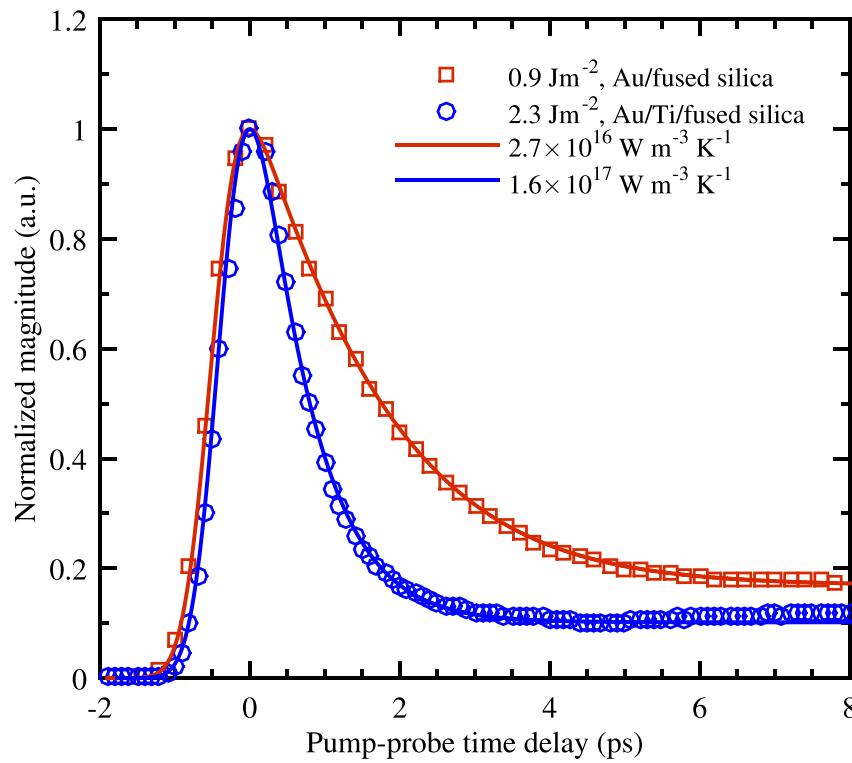
↓

Thermal diffusion (hundreds of ps to ns)

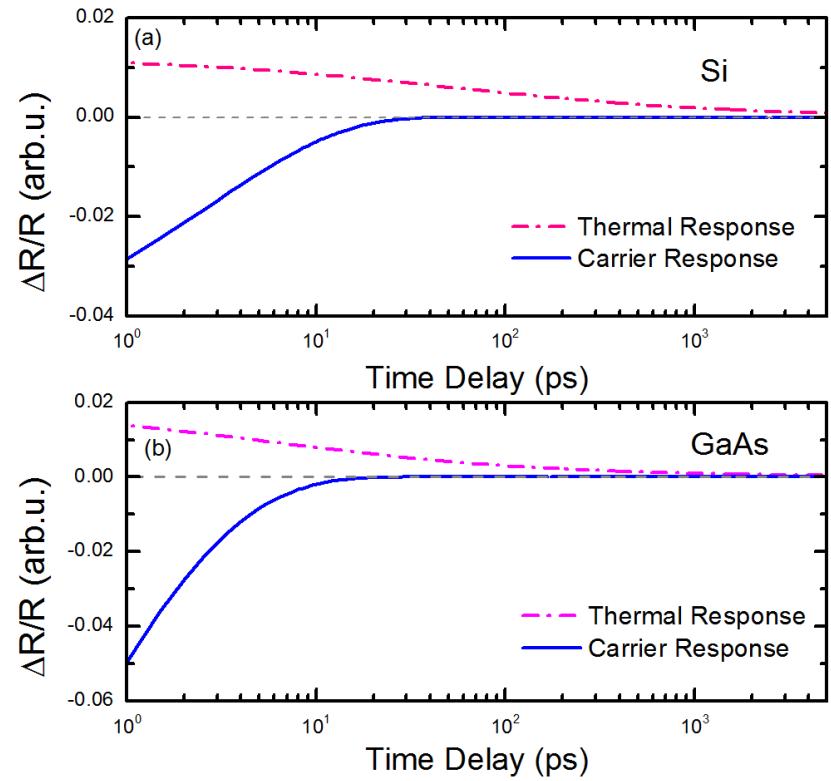
# What can we measure with TDTR?

## Hot electron relaxation and recombination

### Electron-phonon coupling in metals



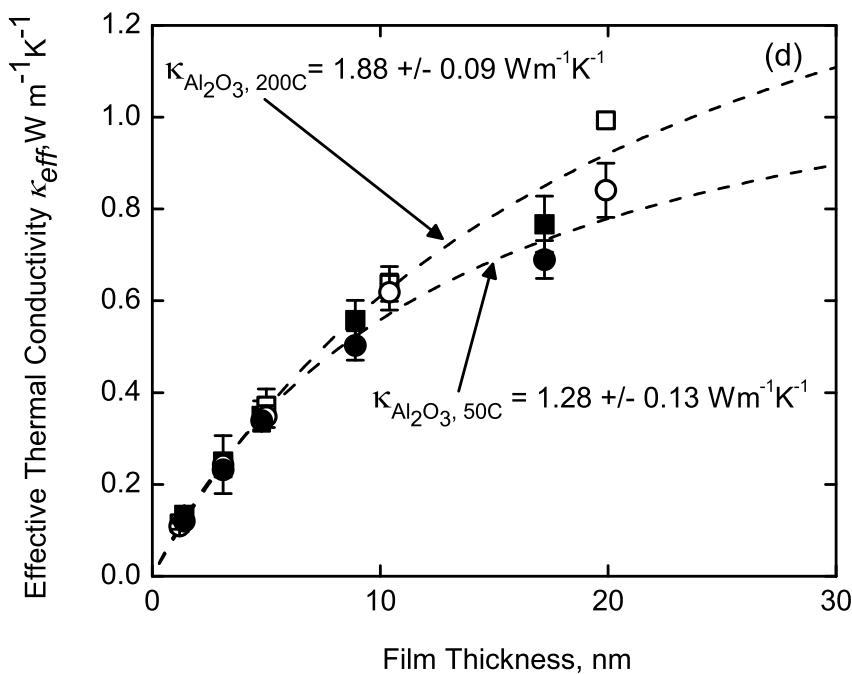
### Excited state recombination in semiconductors



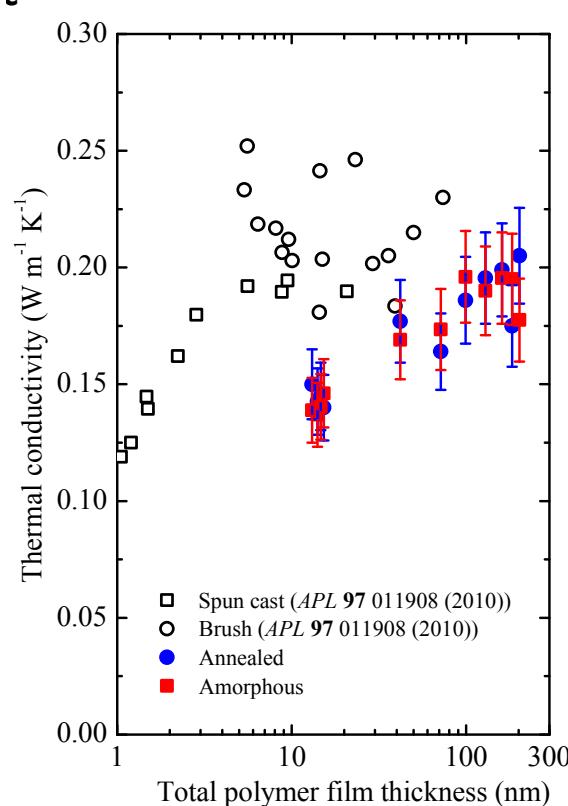
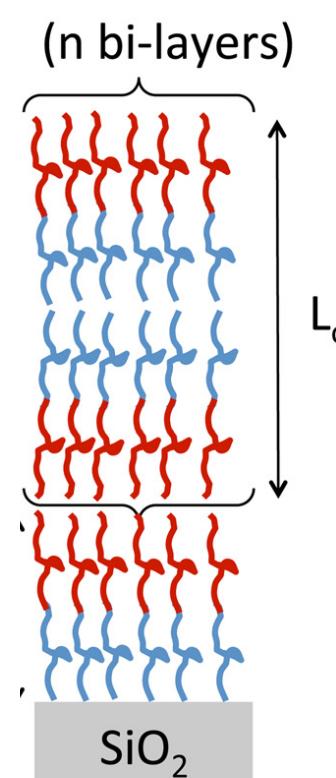
# What can we measure with TDTR?

## Thermal conductivity of extremely thin films/interfaces

### ALD-grown thin films



### Block co-polymer thin films



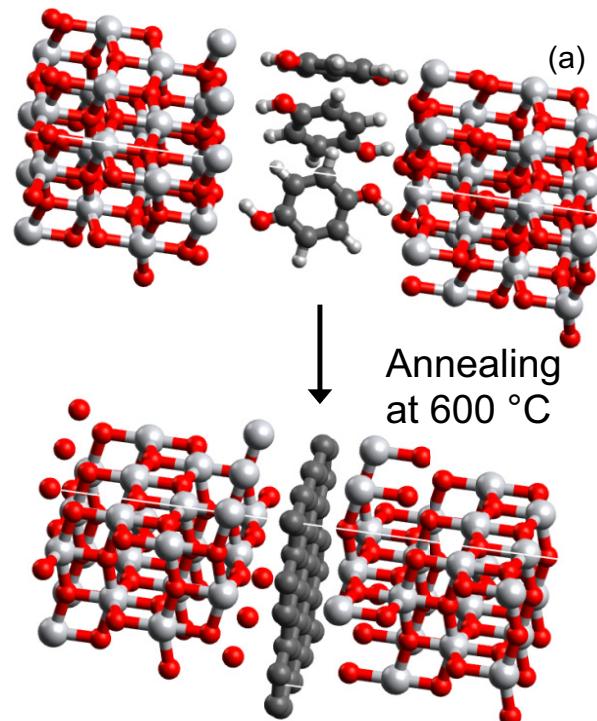
*Thin Solid Films* 650, 71

*J. Heat Trans.* 138 024505

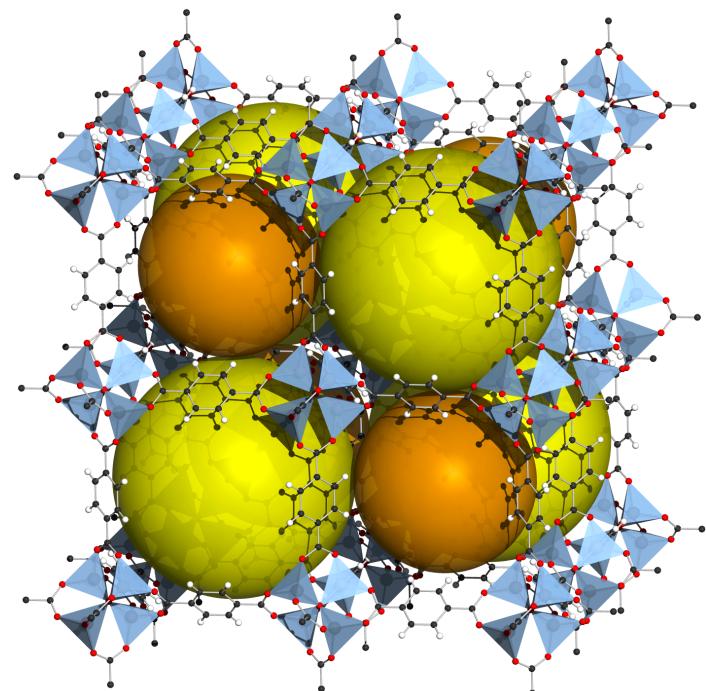
# What can we measure with TDTR?

Heat capacity of thin films and some bulk systems

## Organic/inorganic hybrid “superlattices”



## Metal-organic frameworks “MOFs”

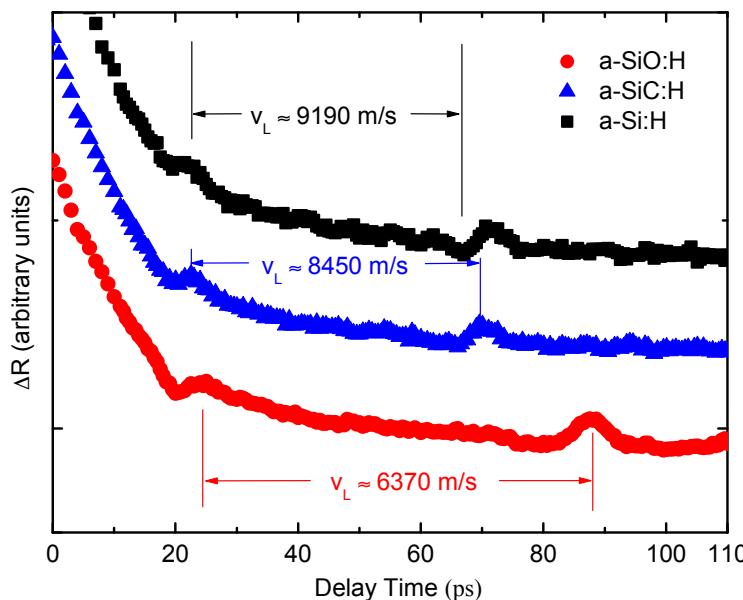
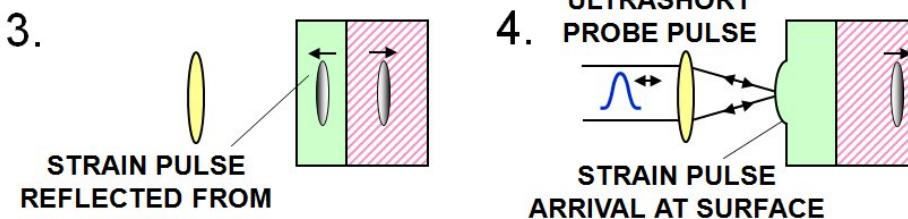
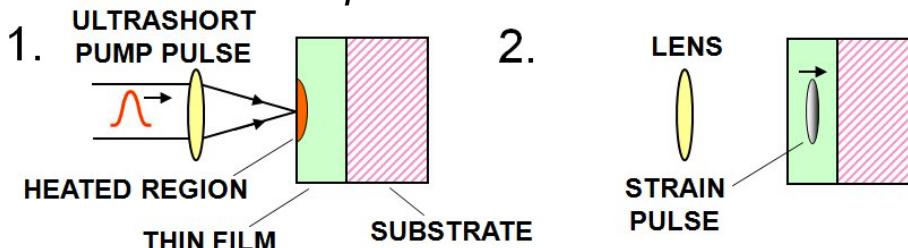


Phys. Rev. B **93**, 024201

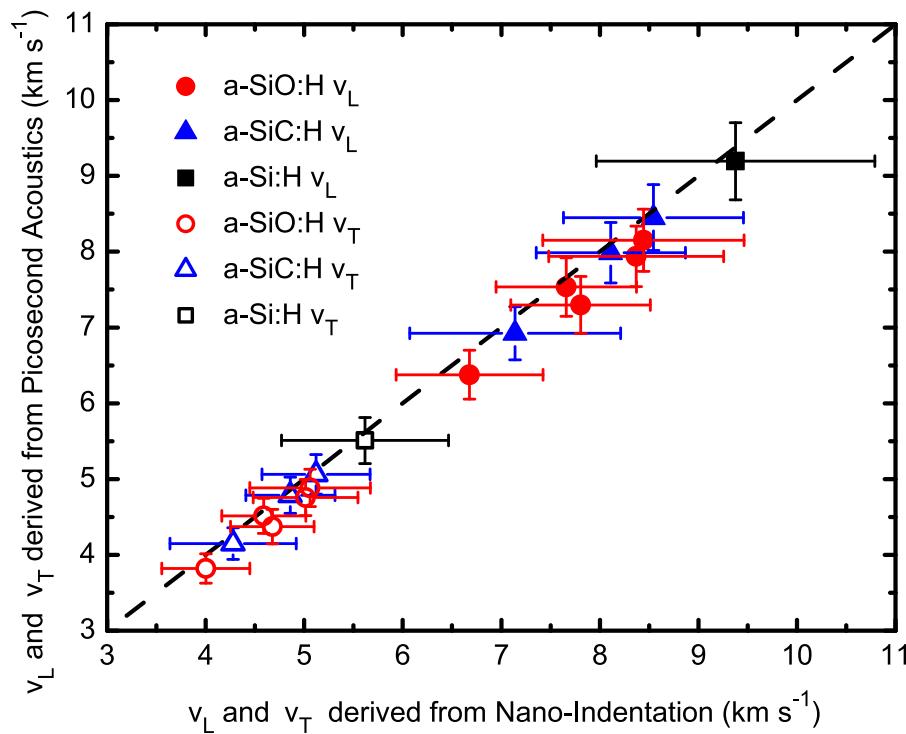
Adv. Mater. **27**, 3453  
MRS Bulletin **41**, 877

# What can we measure with TDTR?

Wikipedia

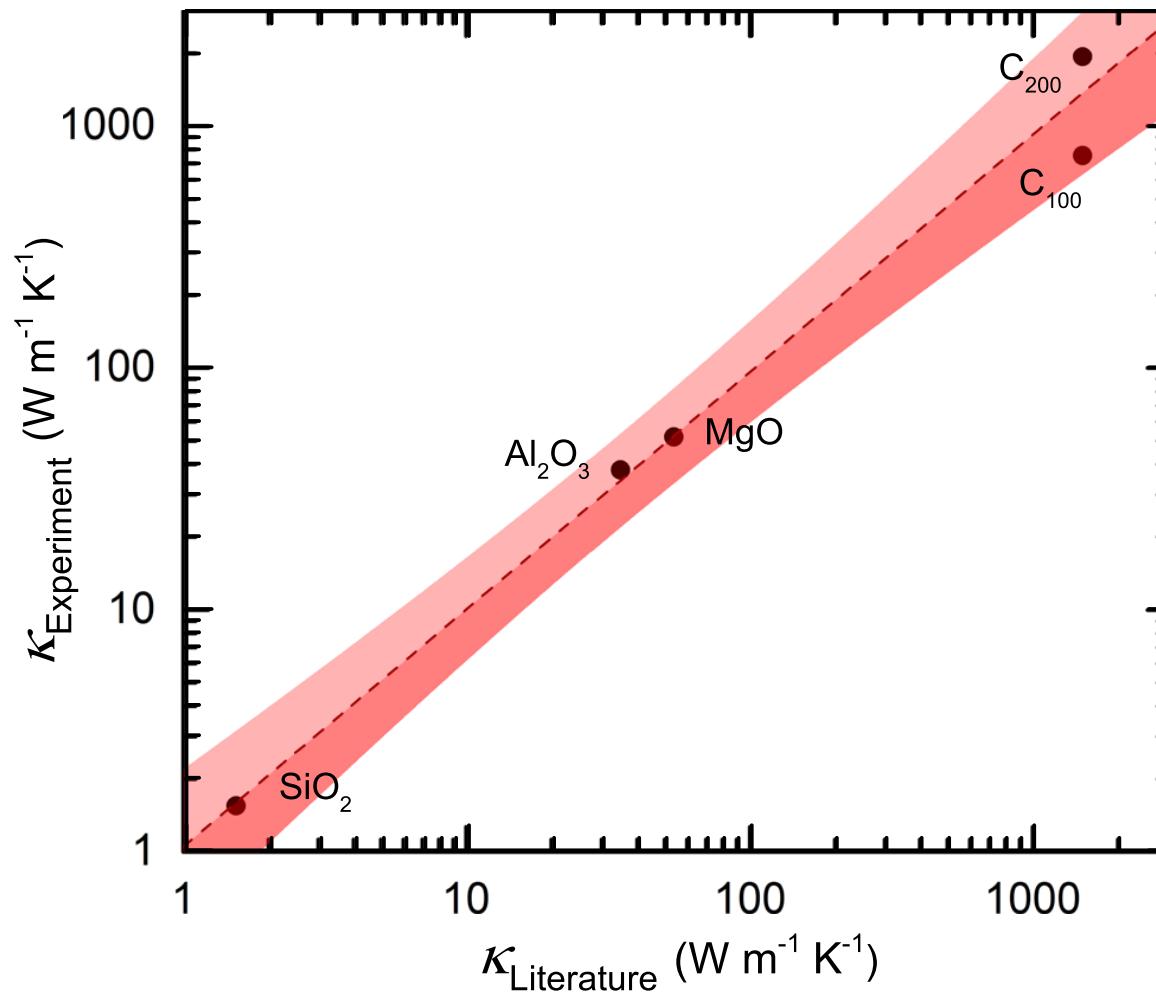


Measurement of strain  
wave/acoustic wave  
propagation in thin films



# What can we measure with TDTR?

## Bulk materials

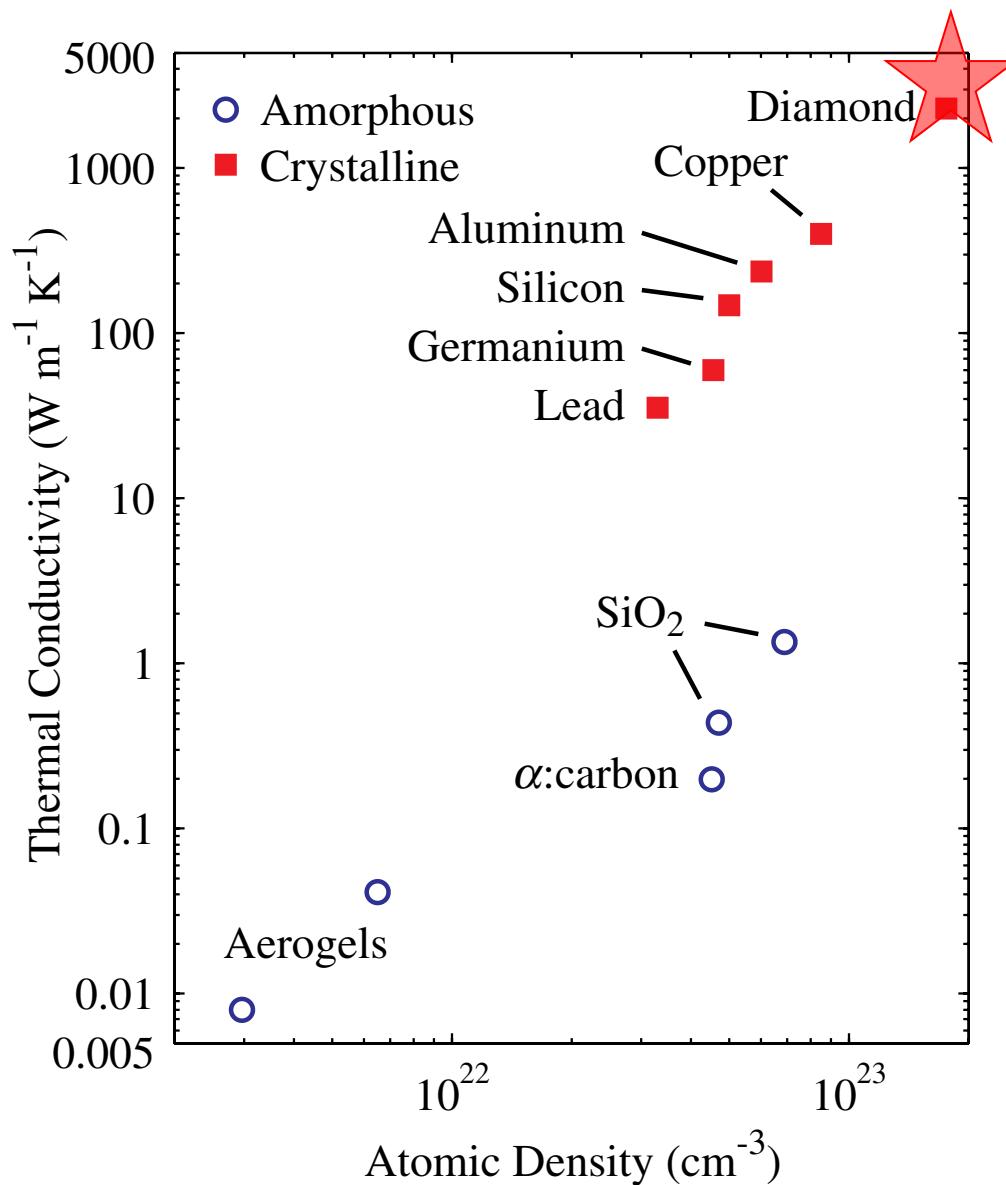


*Rev. Sci. Instrum.* **87**, 094902  
*Appl. Phys. Lett.* **111**, 151902

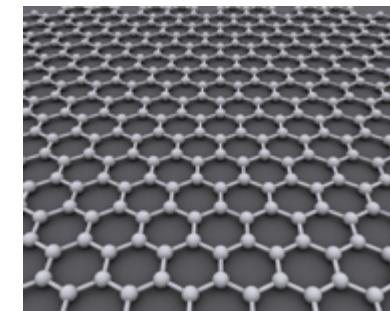
# Outline

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- Heat transport across single molecule interfaces: when does a molecule become a defect?
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# Thermal conductivity of materials – Role of the bond



Ultrahigh thermal conductivities  
(diamond, graphene, BAs, etc)

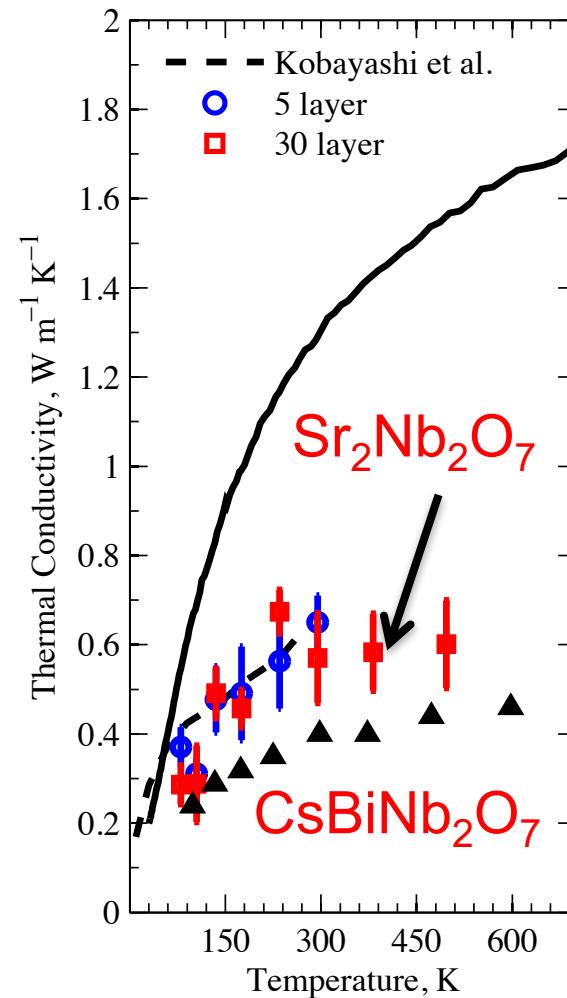
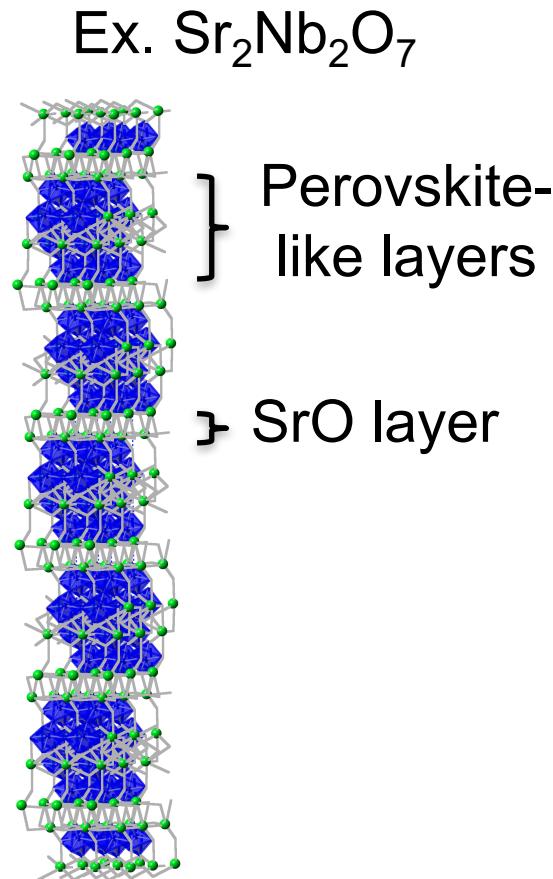


Light masses, stiff bonds

How do we go the other way?  
Heavy masses, weak bonds

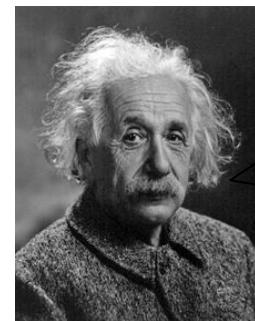
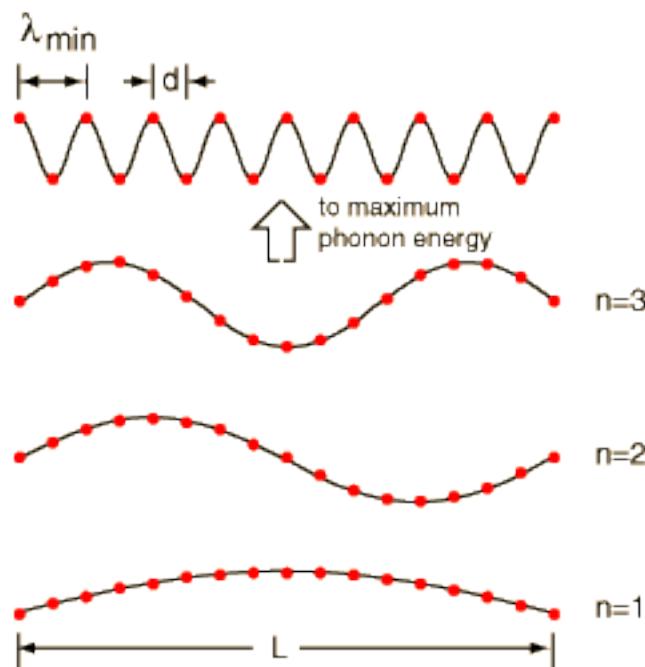
# Thermal conductivity of materials – Role of the bond

Layered structures can exhibit ultralow thermal conductivity



# The Einstein oscillator

The phonon picture  
(coupled oscillators):  
several different wavelength in a  
lattice (many energies)



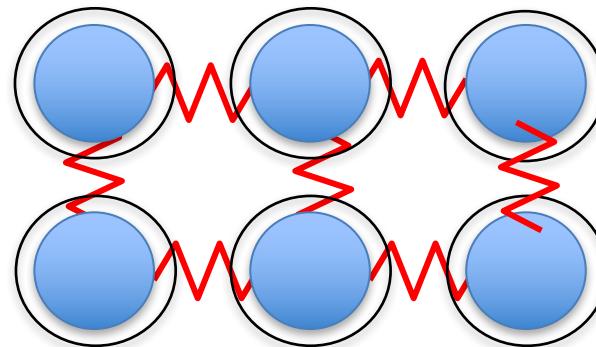
Vibrations of atoms  
are independent  
with random phases

Single frequency of vibration of  
atom and energy “hops” from  
one site to another

$$\lambda_{\min} = 2d$$

$$\lambda_n = \frac{2L}{n}$$

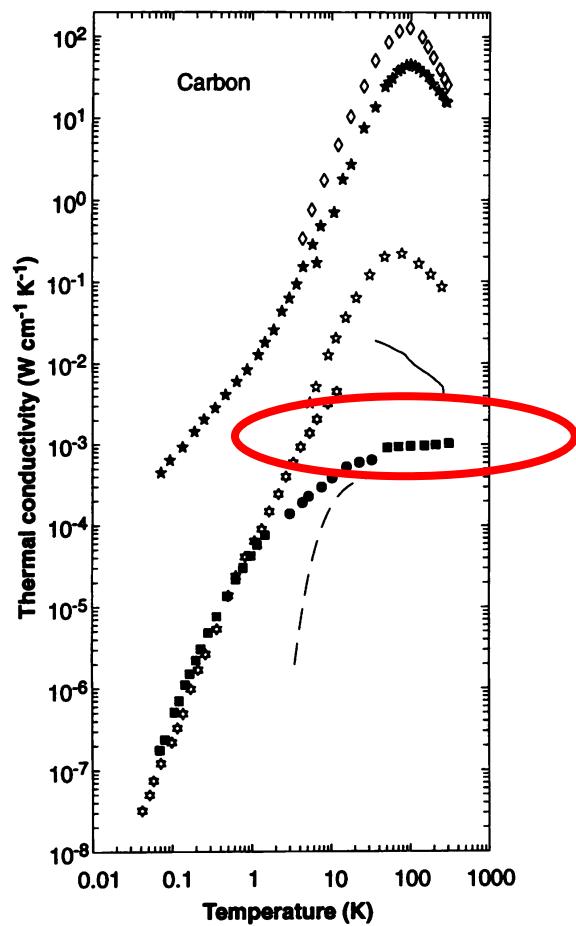
$$\lambda_{\max} = 2L$$



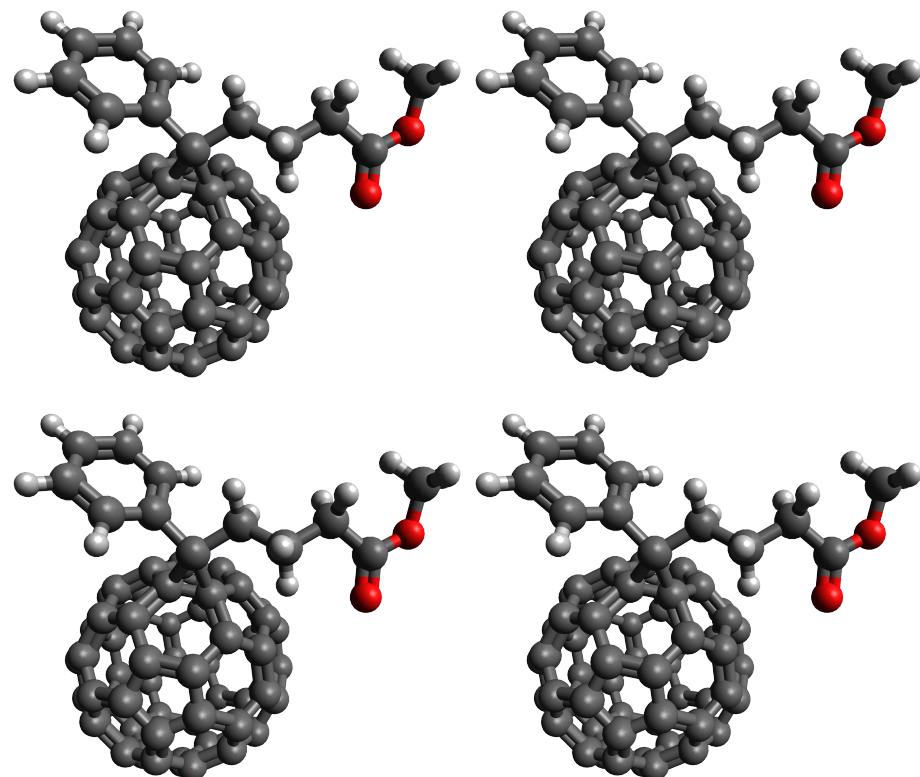
**“Springs” are very very weak**

# The Einstein oscillator – weakly interacting buckyballs

## Fullerene films: Low thermal conductivities



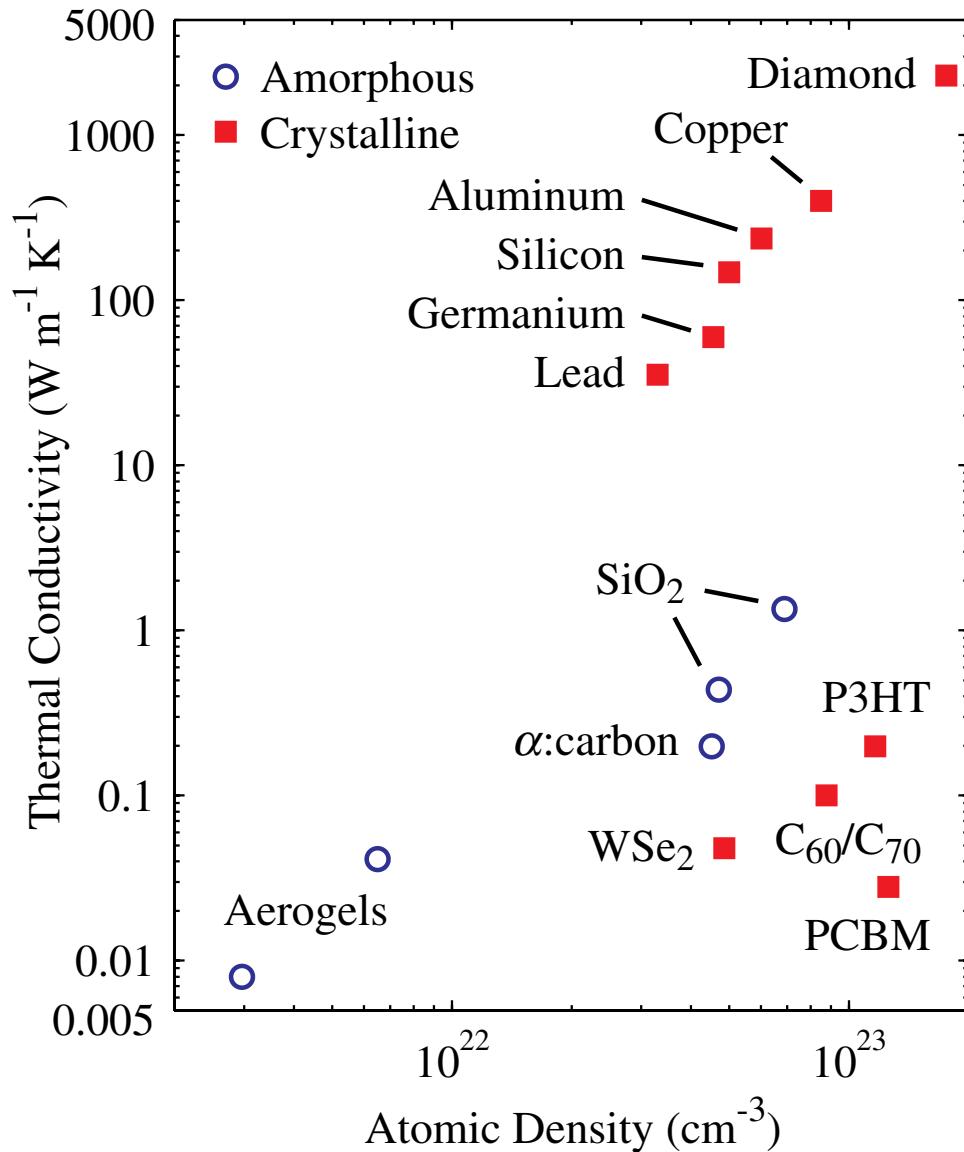
Add “disorder” with molecules:  
[6,6]-phenyl C<sub>61</sub>-butyric acid  
methyl ester (PCBM)



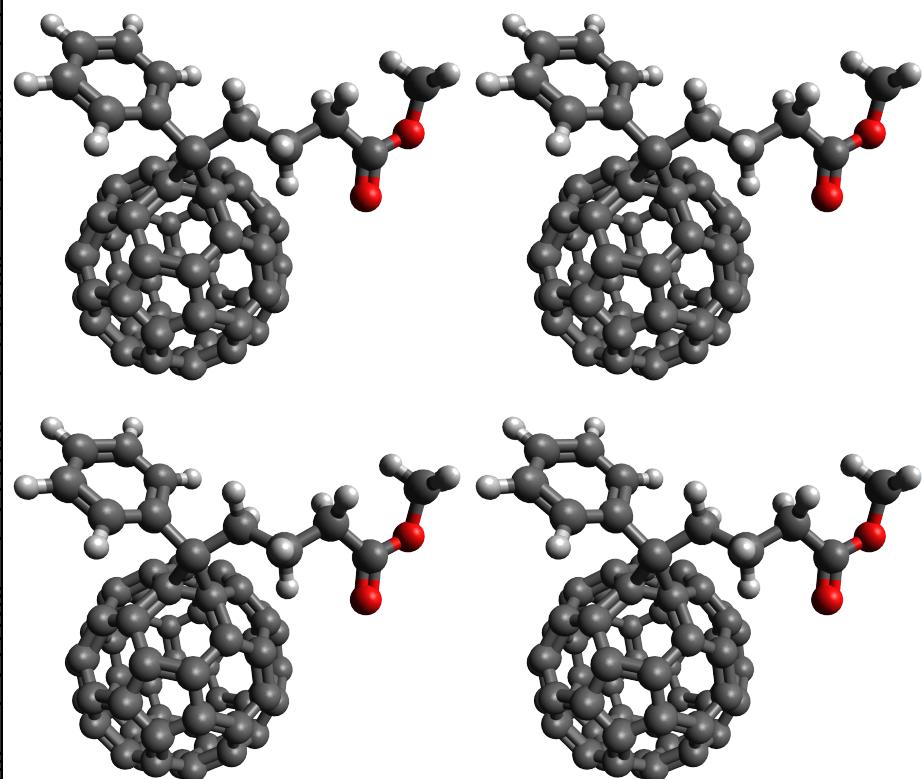
Science 259, 1145

# New lower extreme of thermal conductivity of materials

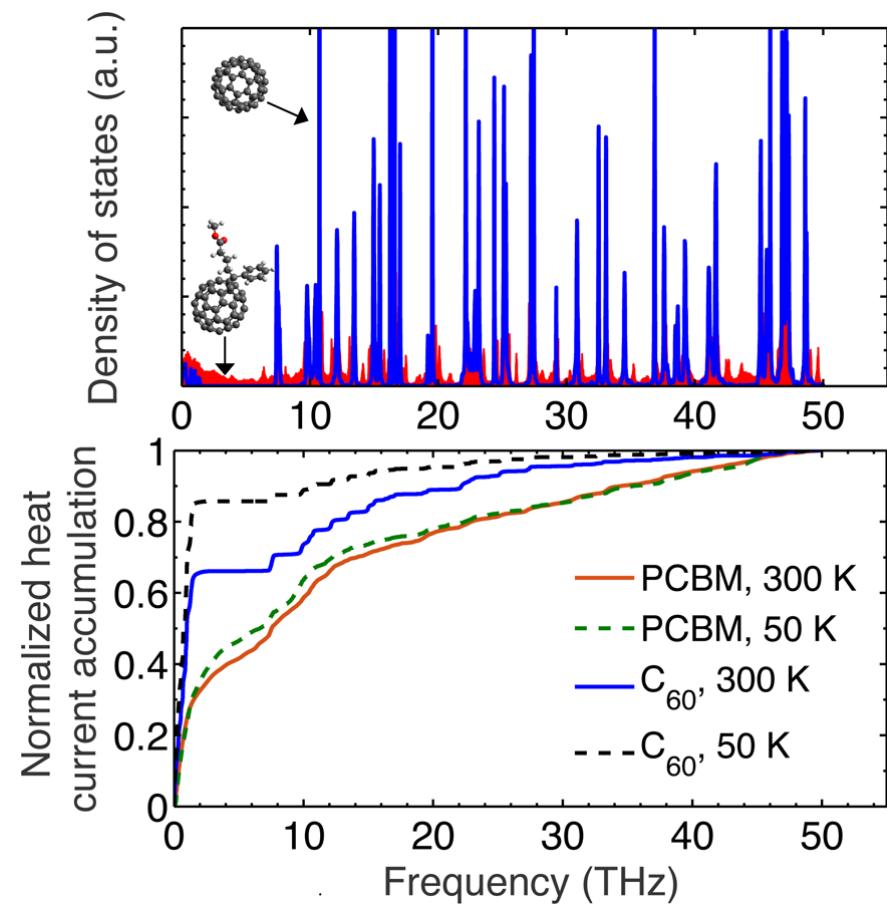
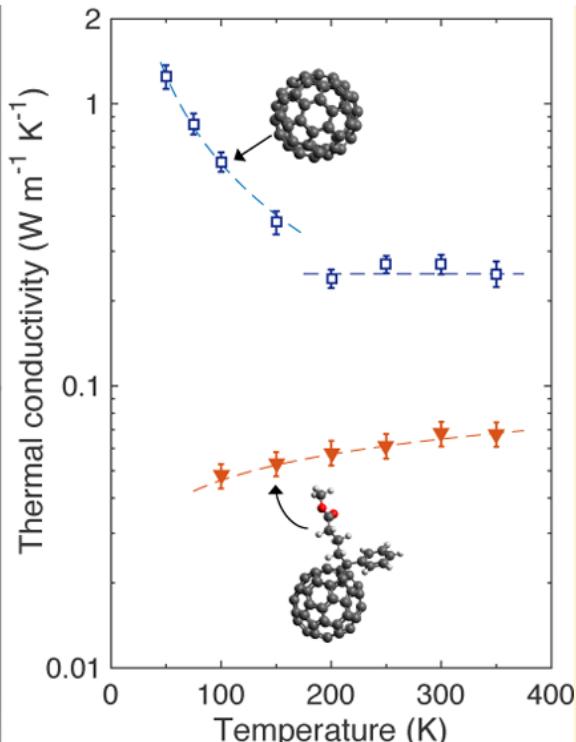
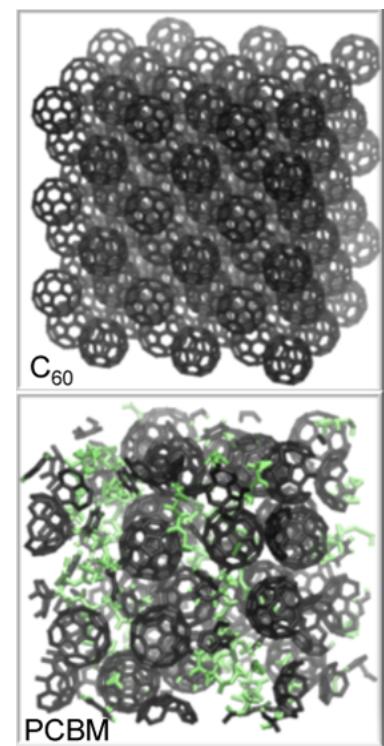
But why???



$$\kappa = \frac{1}{3} Cv\lambda = \frac{1}{3} Cv_g^2 \tau$$

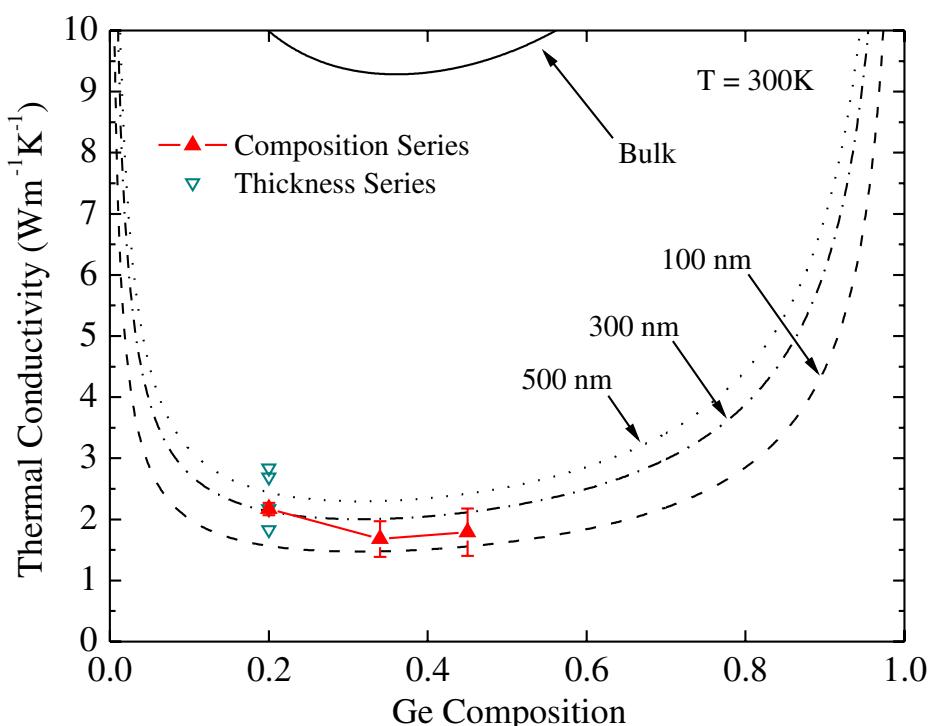


# Turn to molecular dynamics simulations

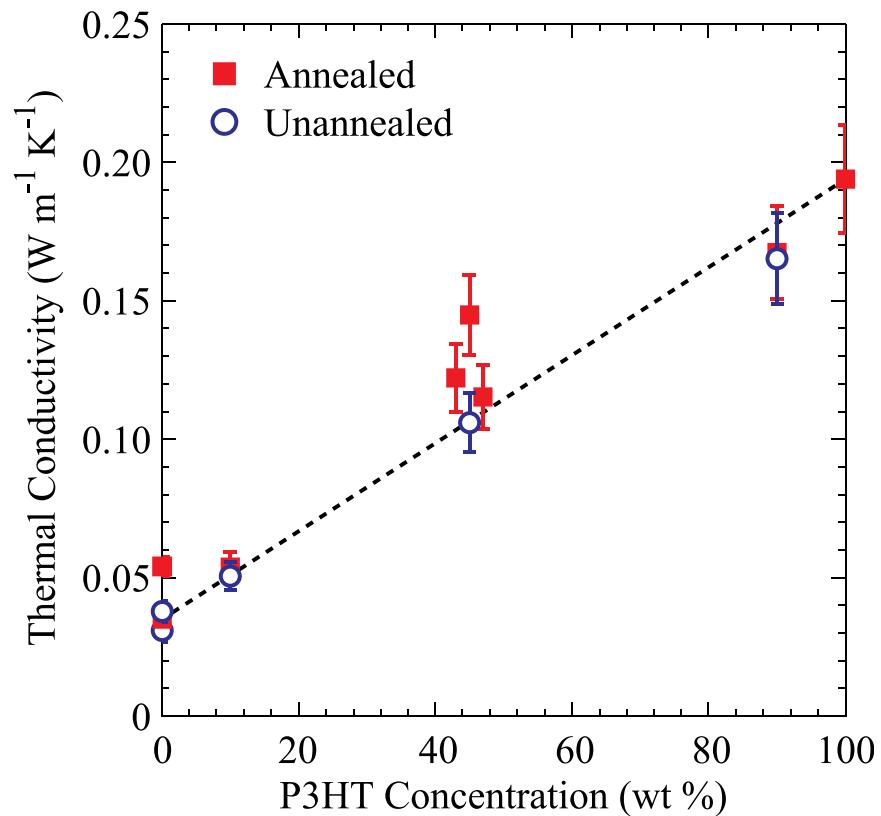


# Tuning the thermal conductivity of molecular films

Covalent bonds  
Crystals/“phonons”  
 $\text{Si}_x\text{Ge}_{1-x}$  alloy



Weak bonds  
Vibrons  
PCBM-P3HT blend

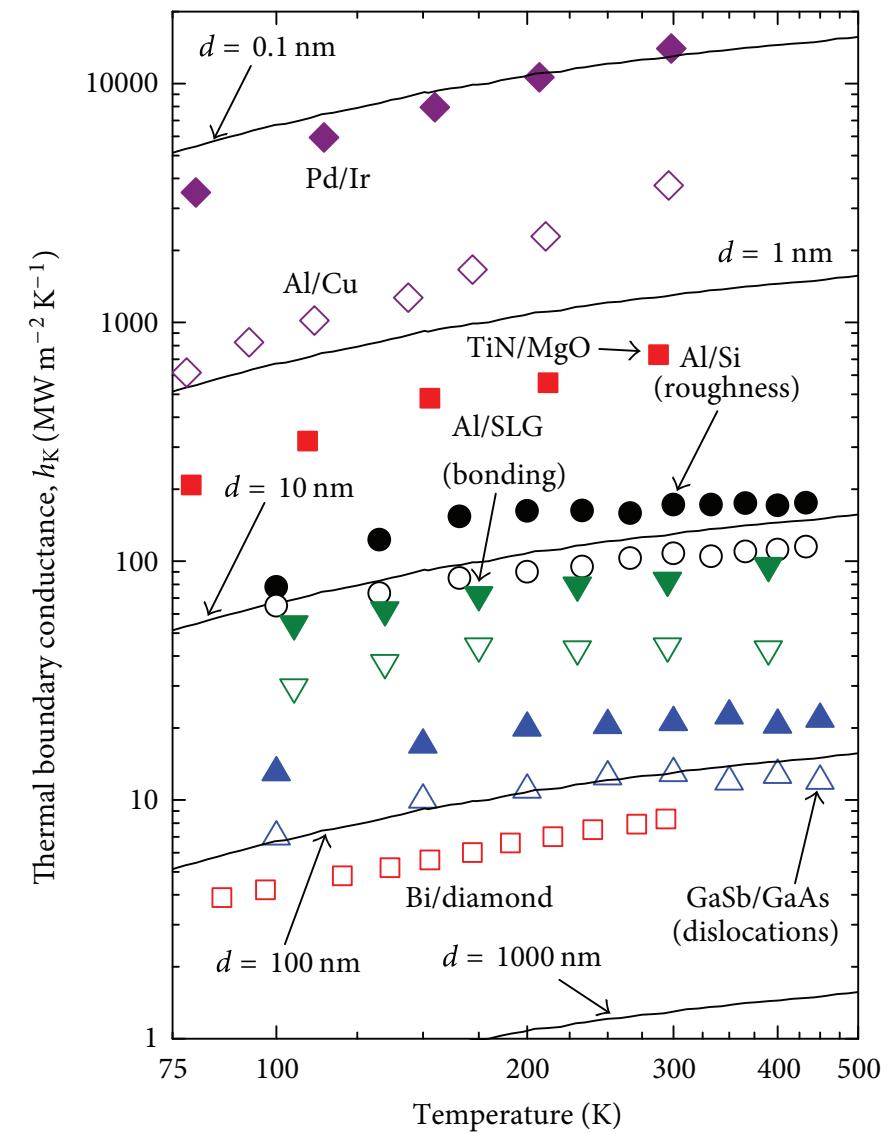
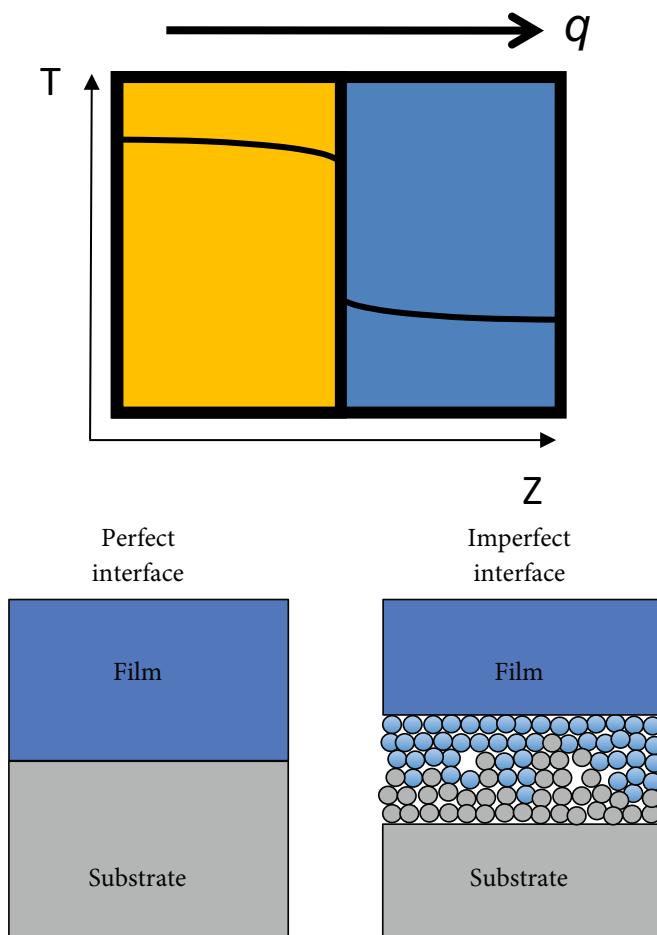


# Outline

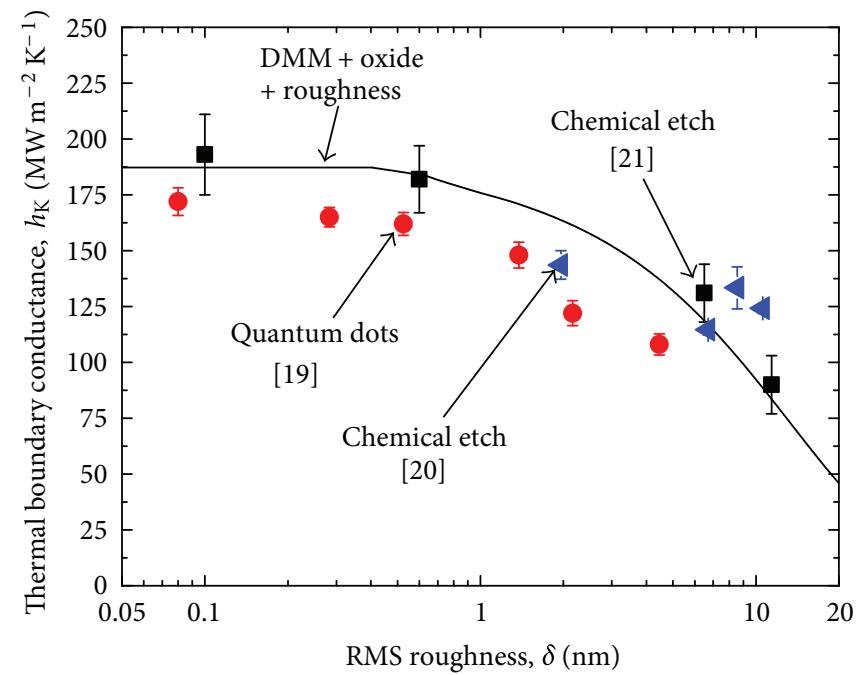
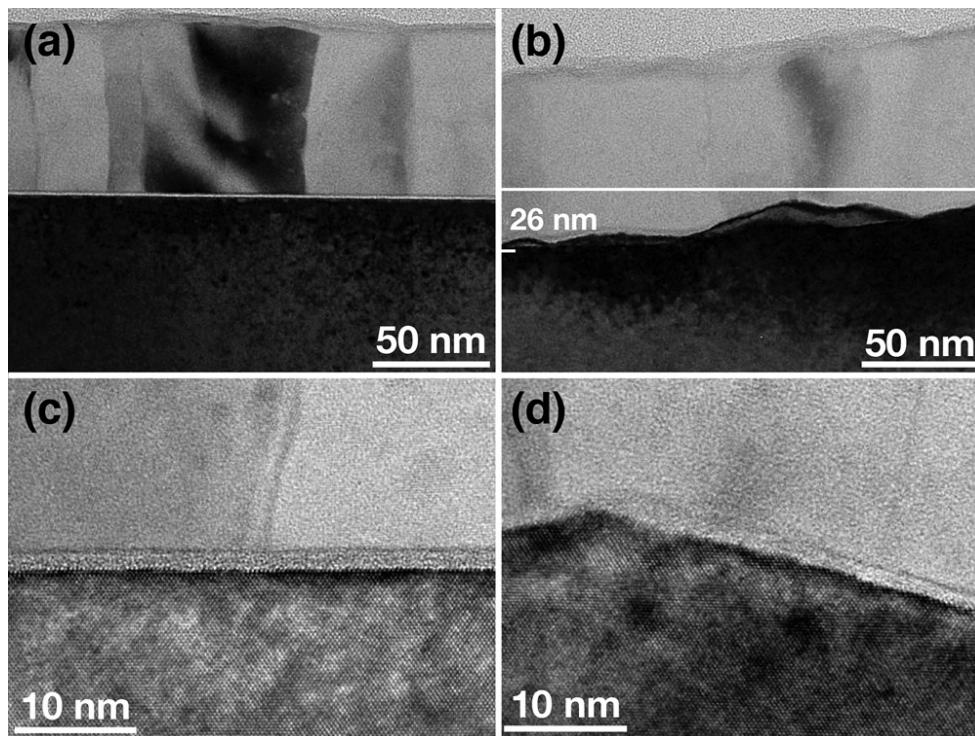
- TDTR: Measurement of thermal conductivity of thin films and thermal resistance across interfaces
- Weakly bonded solids: new lower limits to thermal conductivity
- **Functionalized interfaces at graphene contacts: tuning heat and electrical transport via the interfacial bond**
- Heat transport across single molecule interfaces: when does a molecule become a defect?
- Molecular interfaces in organic/inorganic composites: diffusive scattering via the vibron-phonon interaction

# Thermal boundary conductance – nanoscale resistances

$$q = h_K \Delta T = \frac{1}{R_K} \Delta T$$

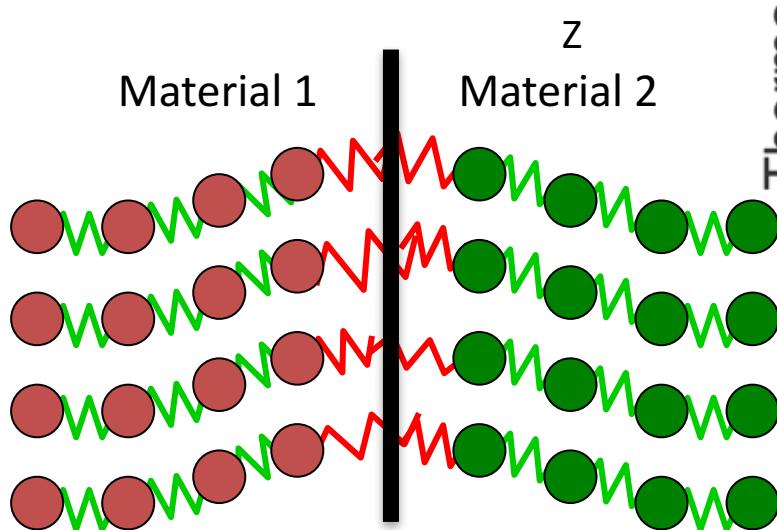
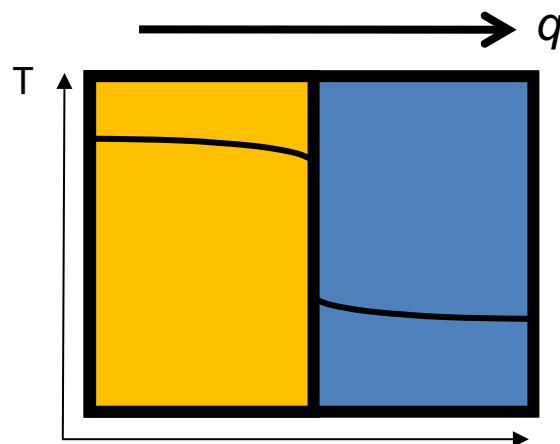


# Thermal boundary conductance – nanoscale resistances

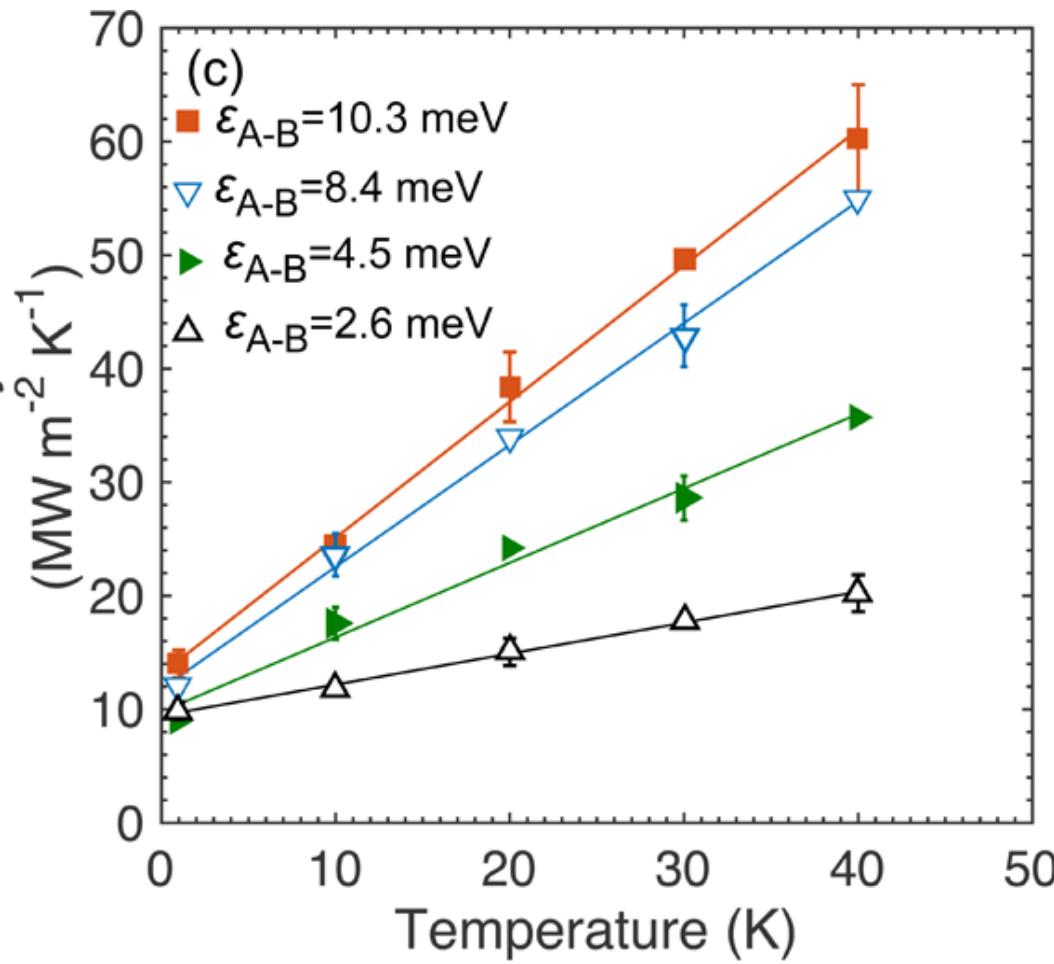


# Increase in bonding increases solid/solid TBC

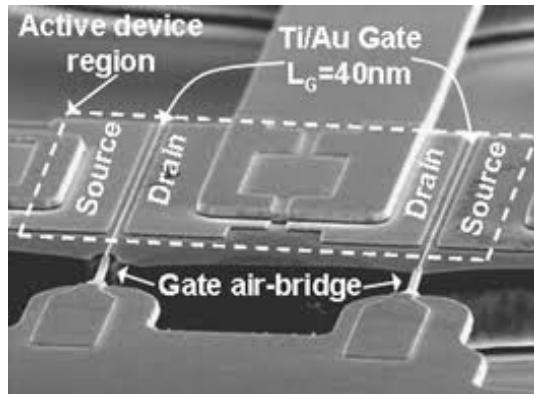
$$q = h_K \Delta T = \frac{1}{R_K} \Delta T$$



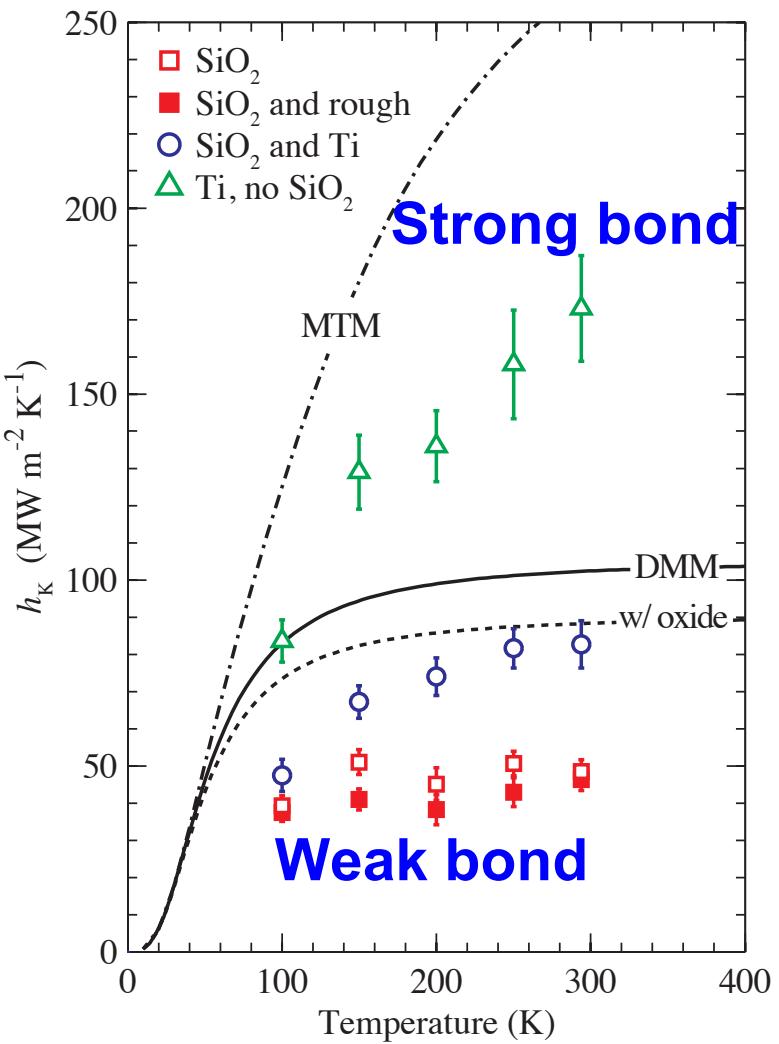
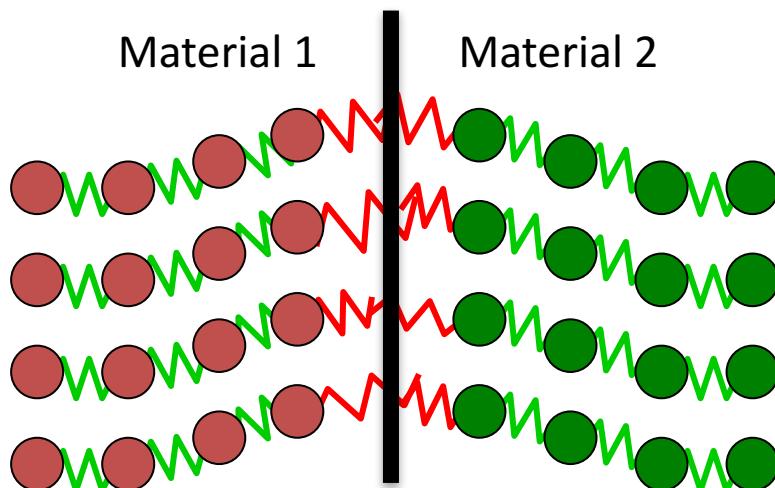
Thermal boundary conductance



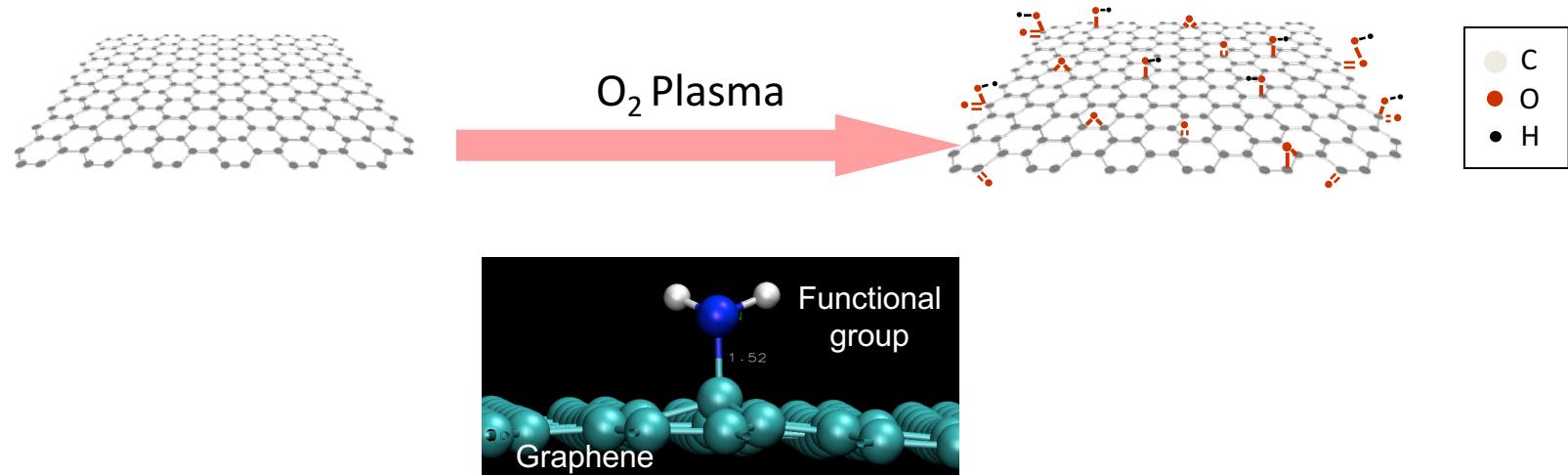
# Thermal boundary conductance – bonding effects



However, bonding plays a **HUGE**  
**ROLE** at the Au/Si interface



# Atmospheric plasma functionalization of graphene surfaces

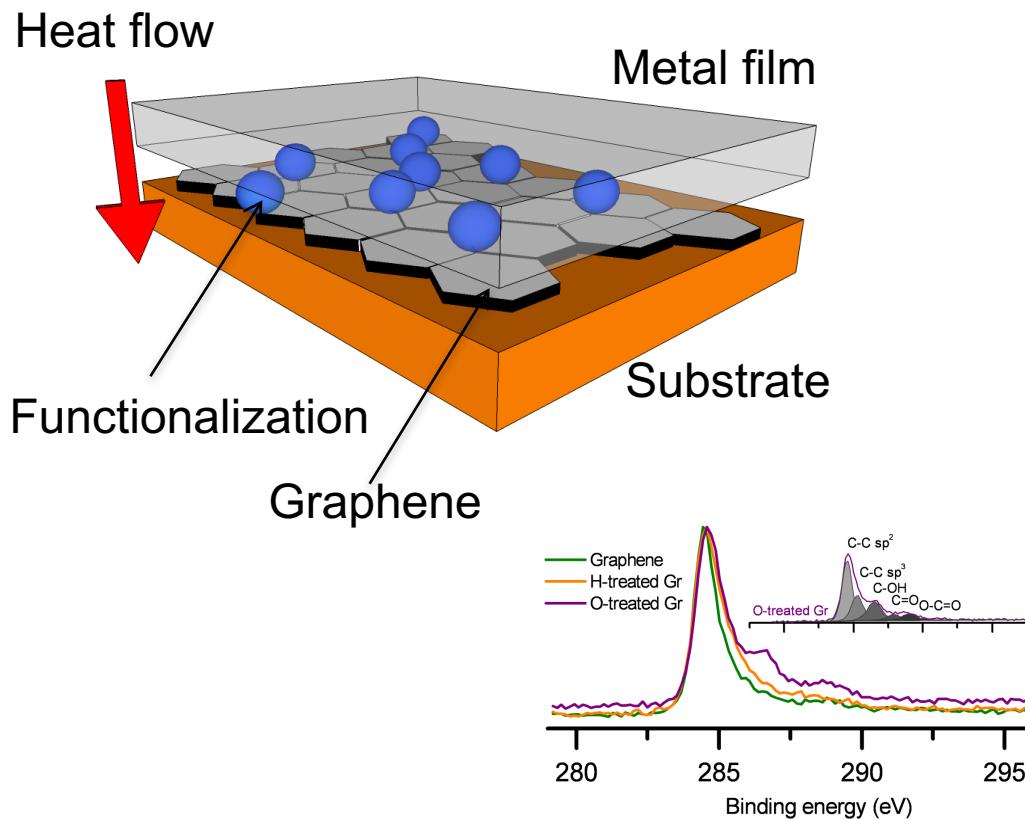


**Functional groups covalently bound to graphene  
Reversible after anneal**

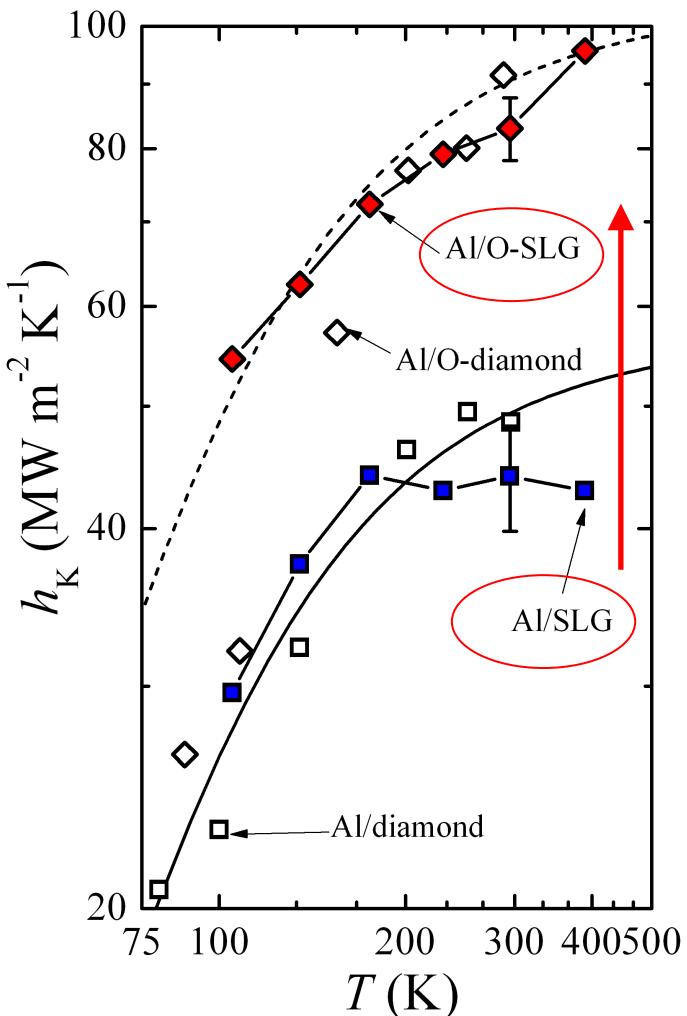
*Appl. Phys. Lett.* **96**, 231501

**Collaboration:** Scott Walton (NRL)

# Atmospheric plasma functionalization of graphene surfaces



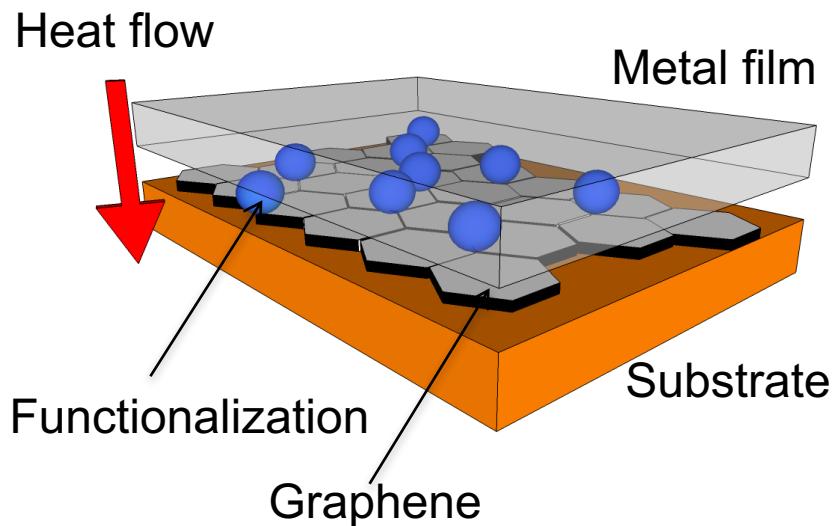
- **Al/graphene interaction increased with oxygen bond (Al-O bond)**
- **But what implications does this have for SLG devices?**



*Nano Lett.* **12**, 590 (2012)

**Collaboration:** Scott Walton (NRL)

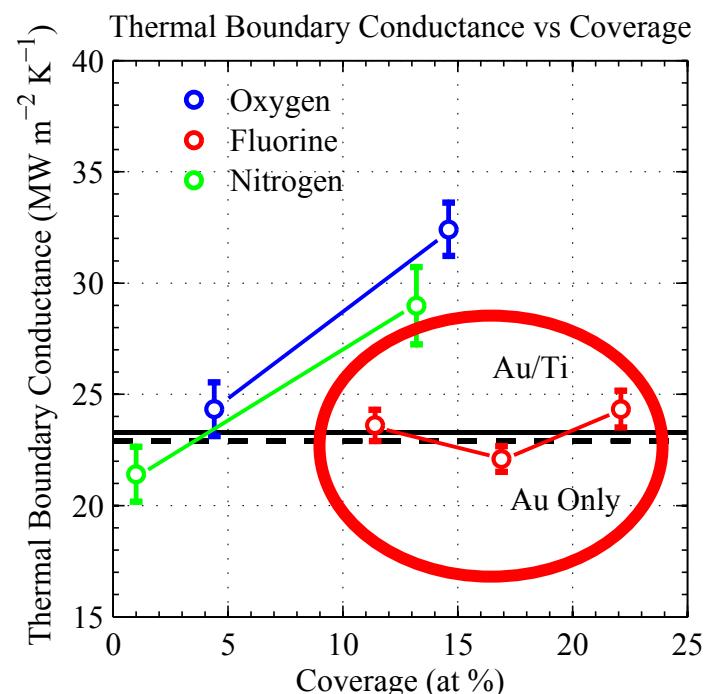
# Au/graphene electronic contacts



**C-F bond inert  
does not want to  
interact!**

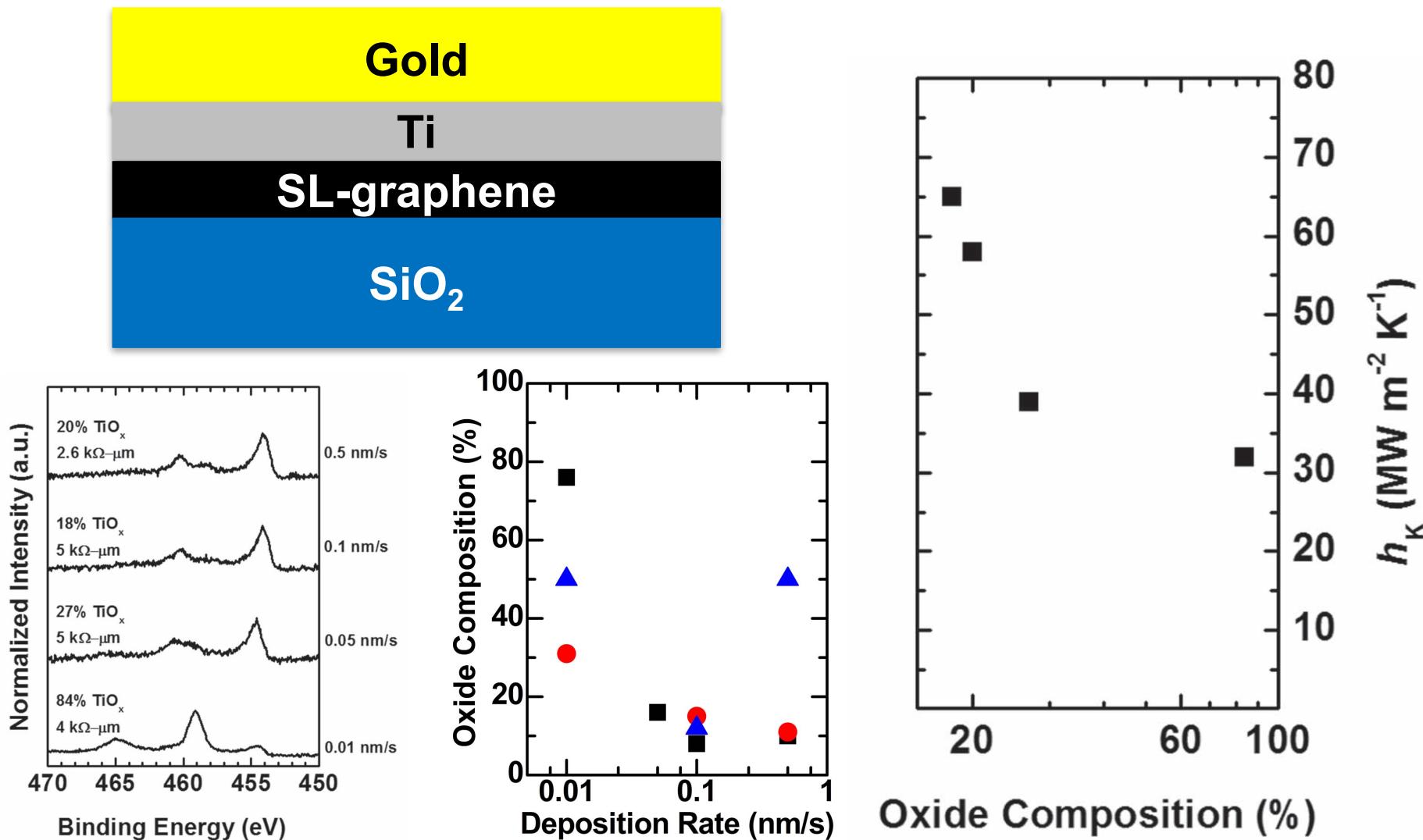
*Nano Lett. 15, 4876 (2015)*

Thermal boundary conductance results  
Similar trends for oxygen as Al  
**Ti adhesion layer does nothing**  
**Fluorine does nothing**



**Collaboration:** Scott Walton (NRL)

# Chemistry effects on the TBC across Au/Ti/graphene

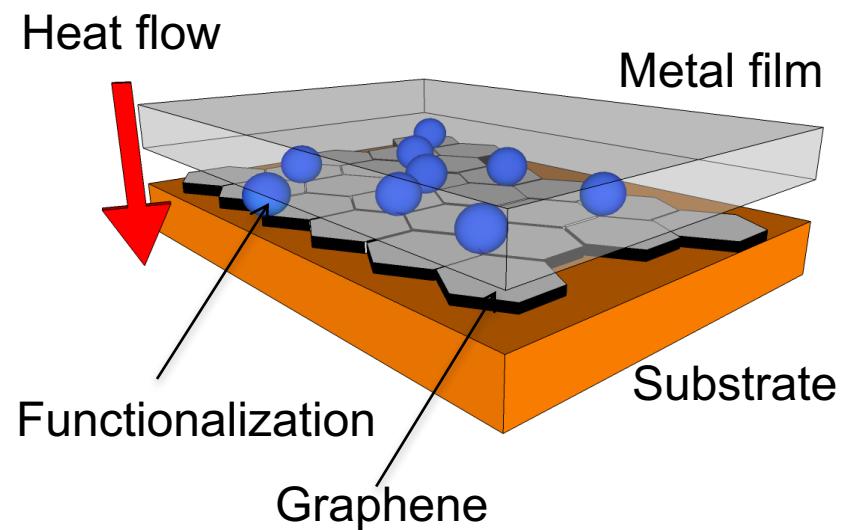


*Nanotechnology* **29**, 145201

**Collaboration:** Stephen McDonnell (UVA)

# Chemistry effects on the TBC across Au/Ti/graphene

## Bonding engineering at interface



## Defect engineering at interface

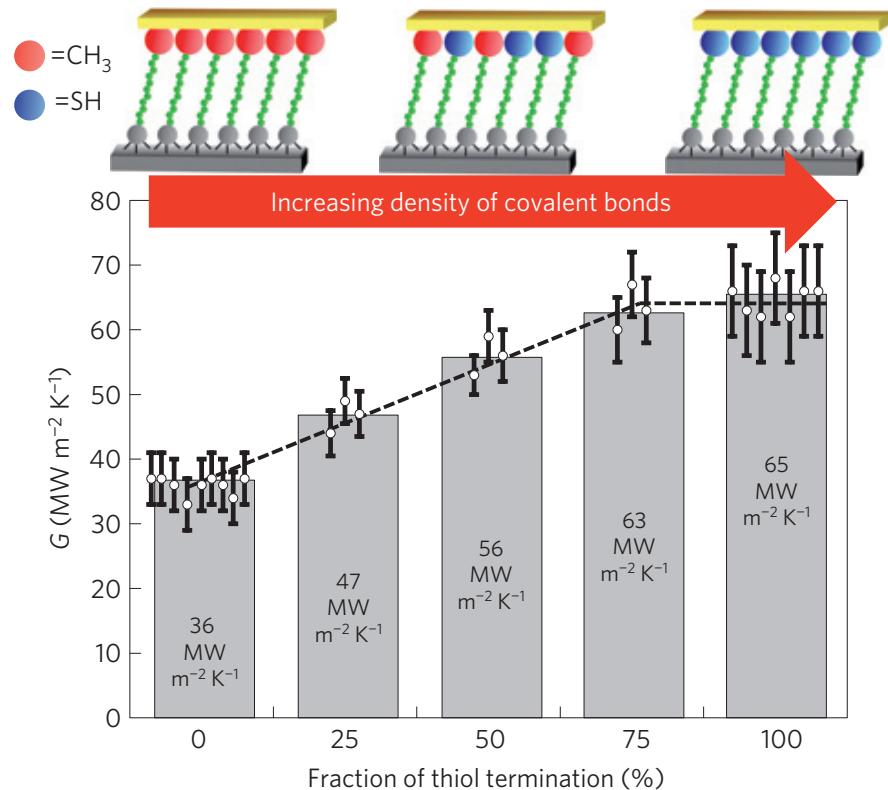


**When does a molecular bond become a defect?**

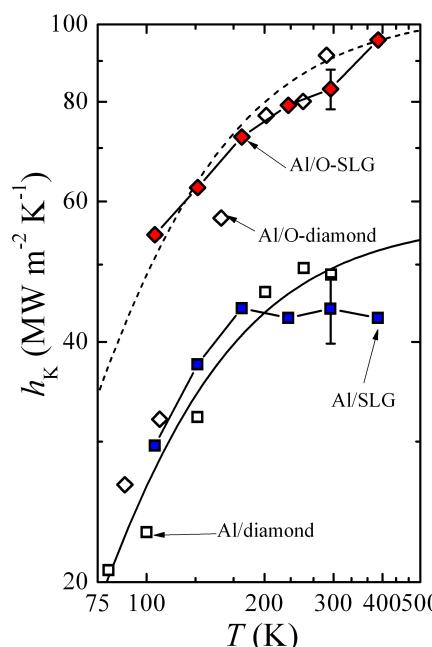
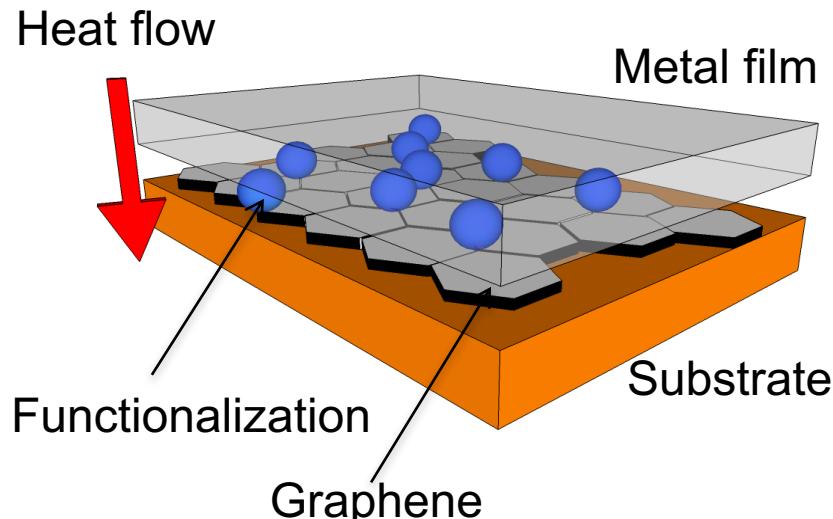
# Outline

- TDTR: Measurement of thermal conductivity of thin films and thermal resistance across interfaces
- Weakly bonded solids: new lower limits to thermal conductivity
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- Molecular interfaces in organic/inorganic composites: diffusive scattering via the vibron-phonon interaction

# Contact chemistry to manipulate TBC

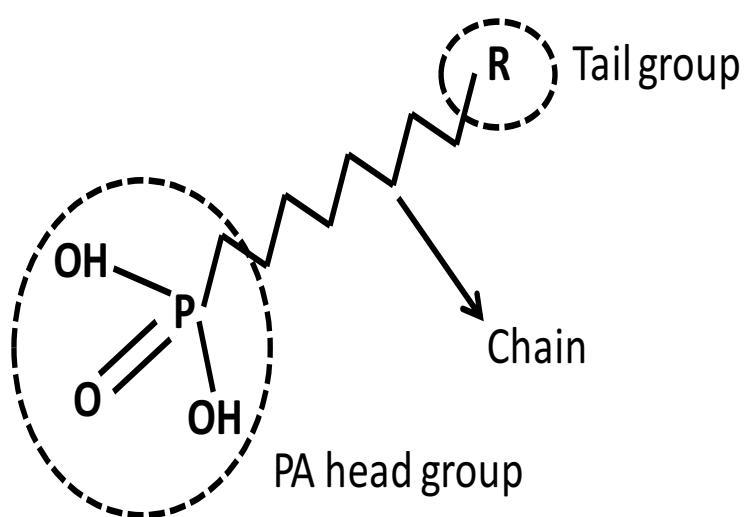


*Nat. Mat. 11, 502 (2012)*



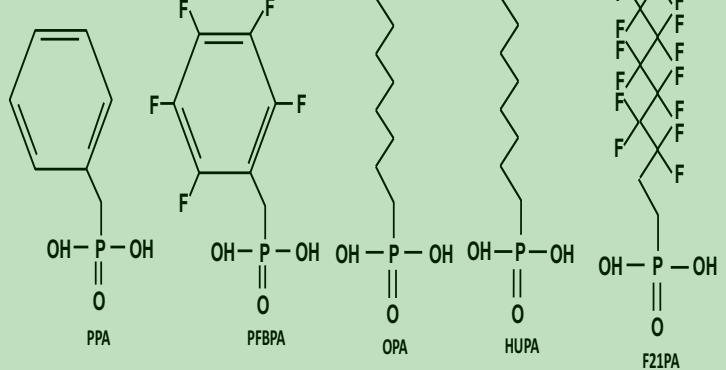
*Nano Lett. 12  
590 (2012)*

# Phosphonic acid interfaces: Size and mass control



Metals: Al, Au, Ni

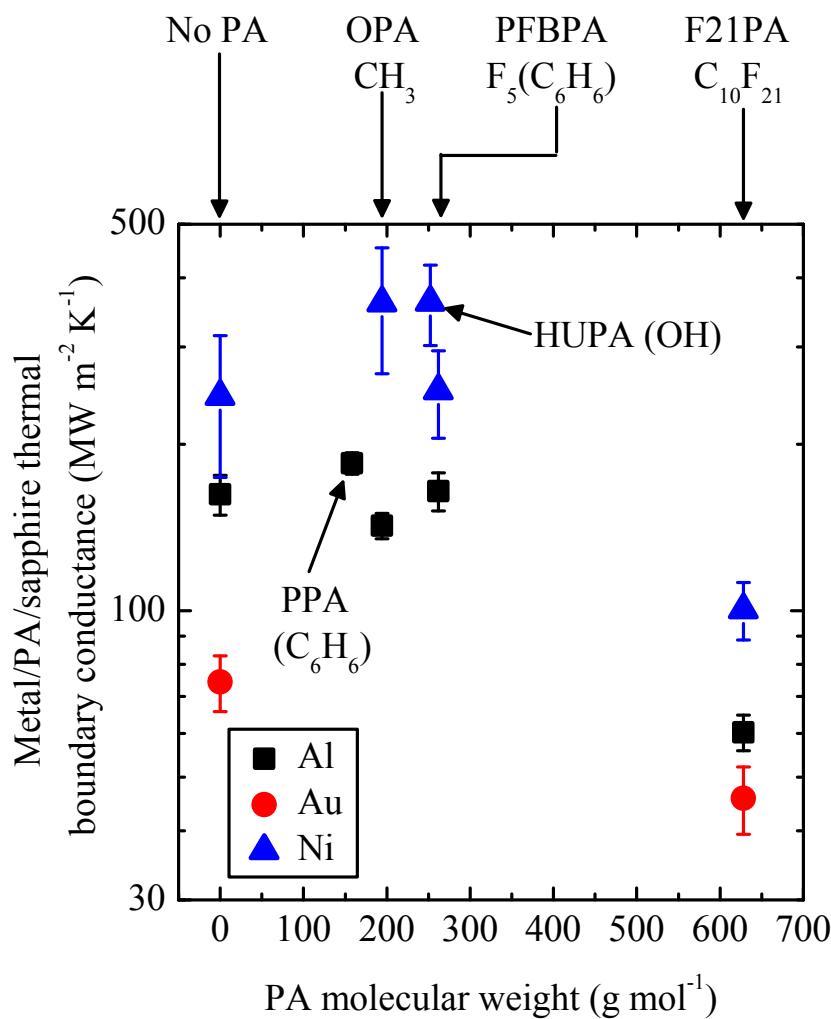
## Phosphonic Acids



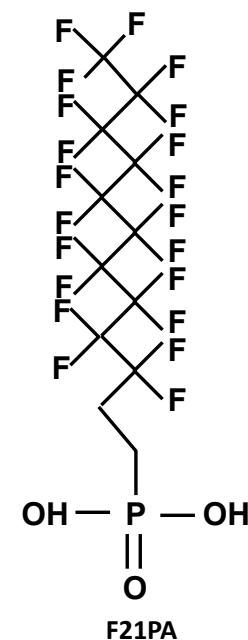
Sapphire substrate

Synthesis at Ga. Tech: Prof. Sam Graham

# Phosphonic acid interfaces: Size and mass control



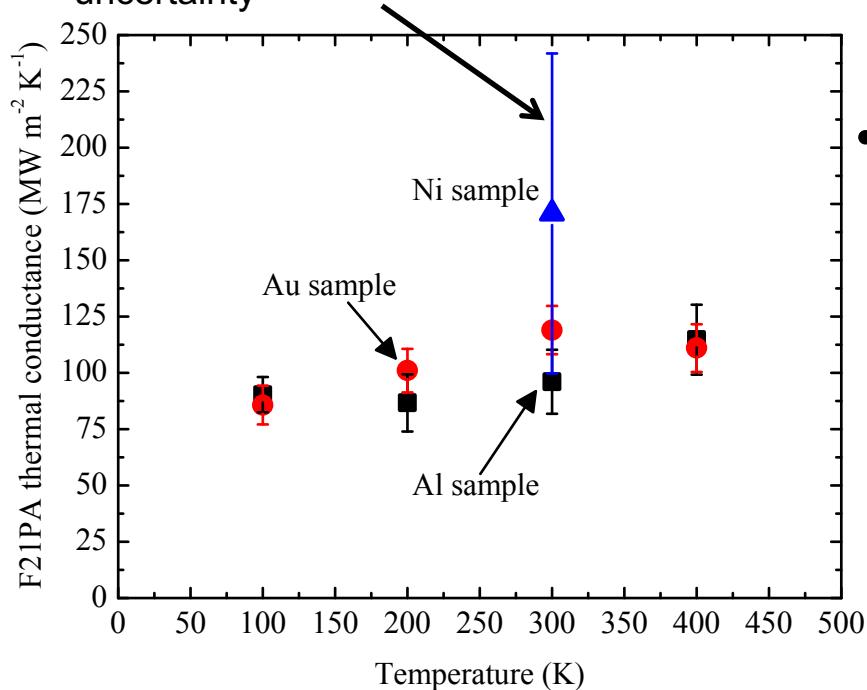
- No correlation with length
- Most noticeable change in TBC with large MW F21PA
- Probably not length effect (HUPA~ 1.7 nm, F21PA ~ 2 nm)



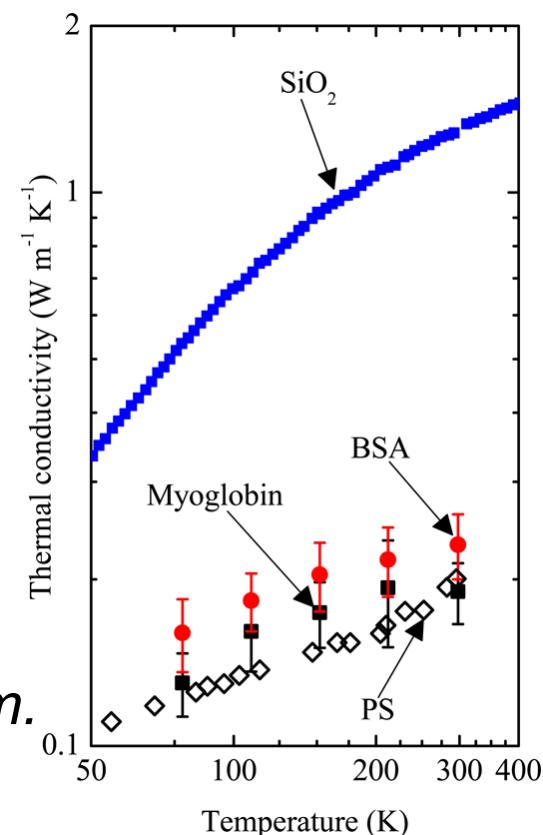
# Is the change in TBC the metal/PA interface or the F21PA itself?

# Heat transport in single F21PA molecules

Uncertainty in Ni thickness,  
but in process of post-  
TDTR analysis to reduce  
uncertainty



- No prior evidence of diffusive heat flow in a molecule (ballistic observed by Wang *et al.* *Science* **317**, 787 (2007))
- ~2 nm molecule, so effective thermal conductivity of F21PA~0.2 W/m/K at RT

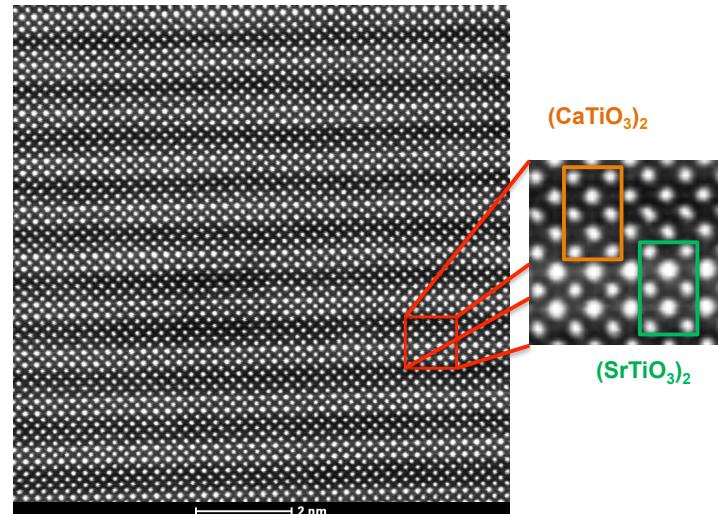


# Outline

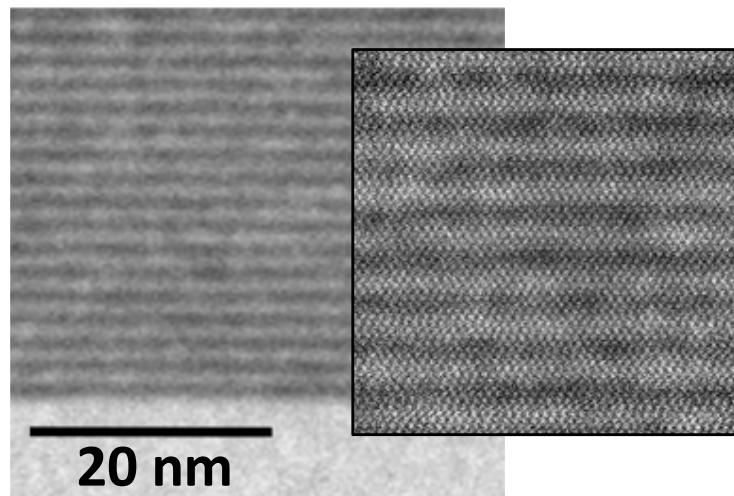
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- **Molecular interfaces in organic/inorganic composites: diffusive scattering via the vibron-phonon interaction**

# Heat transport mechanisms in superlattices

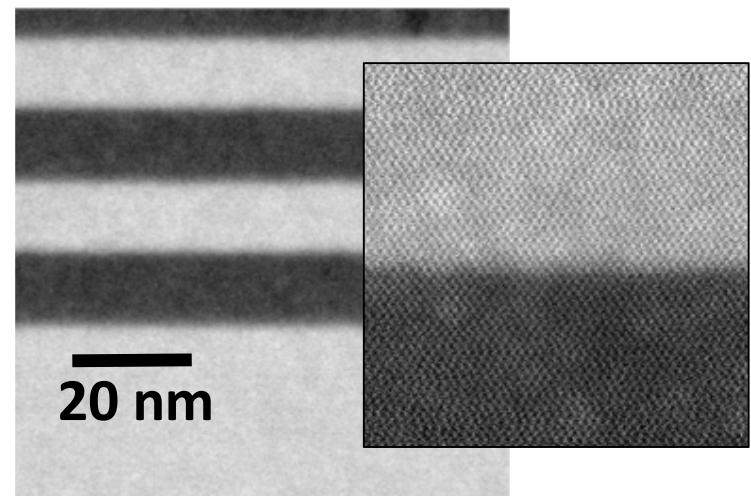
**CaTiO<sub>3</sub>/SrTiO<sub>3</sub>**



**(a)**

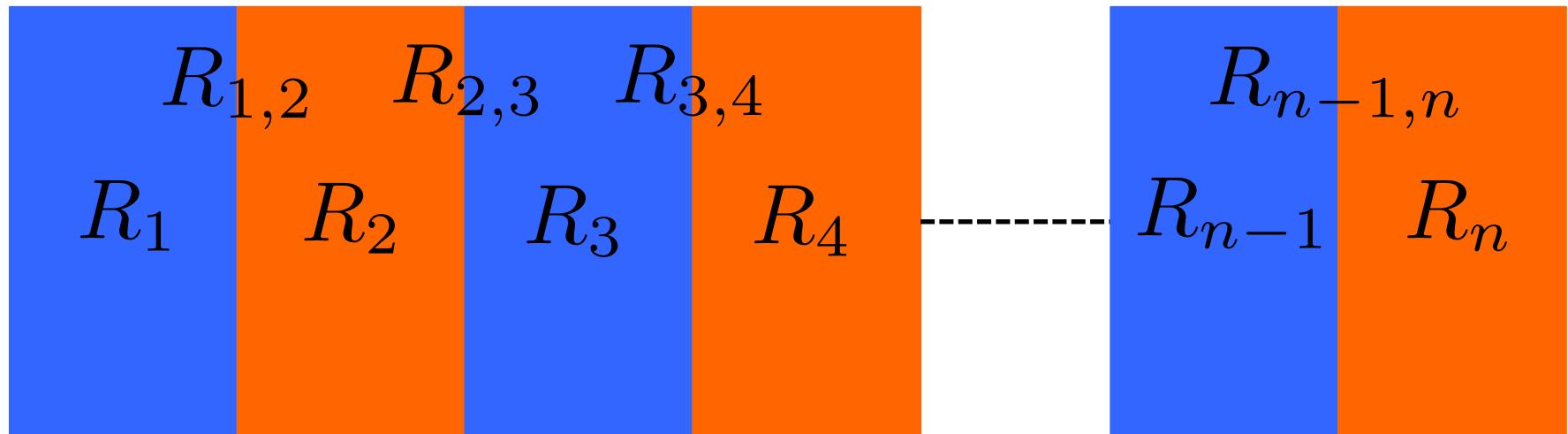


**(b)**



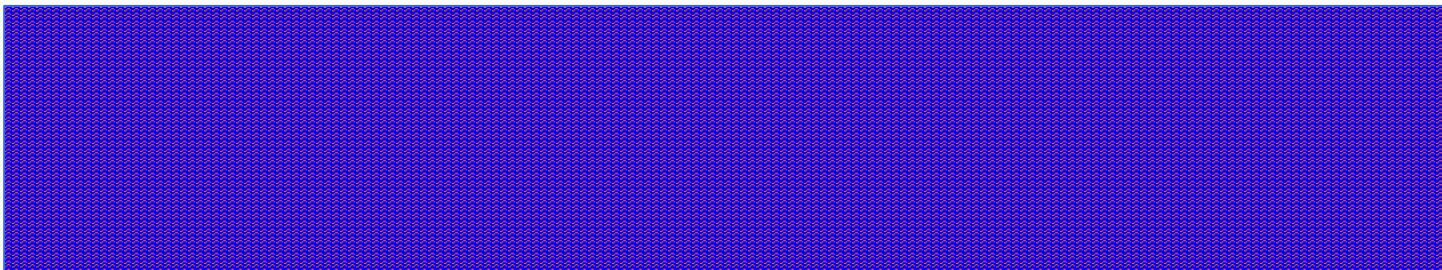
$d_{\text{SL}} = 2 \text{ nm}$  **GaAs/AlAs**  $d_{\text{SL}} = 24 \text{ nm}$

# Incoherent/particle picture of phonon transport in SLs



$$1/\kappa \propto 1/h = \sum_n R_j$$

What if layers are “linked”? – coherent transport



# Coherent transport in superlattices

PHYSICAL REVIEW B

VOLUME 25, NUMBER 6

15 MARCH 1982

VOLUME 84, NUMBER 5

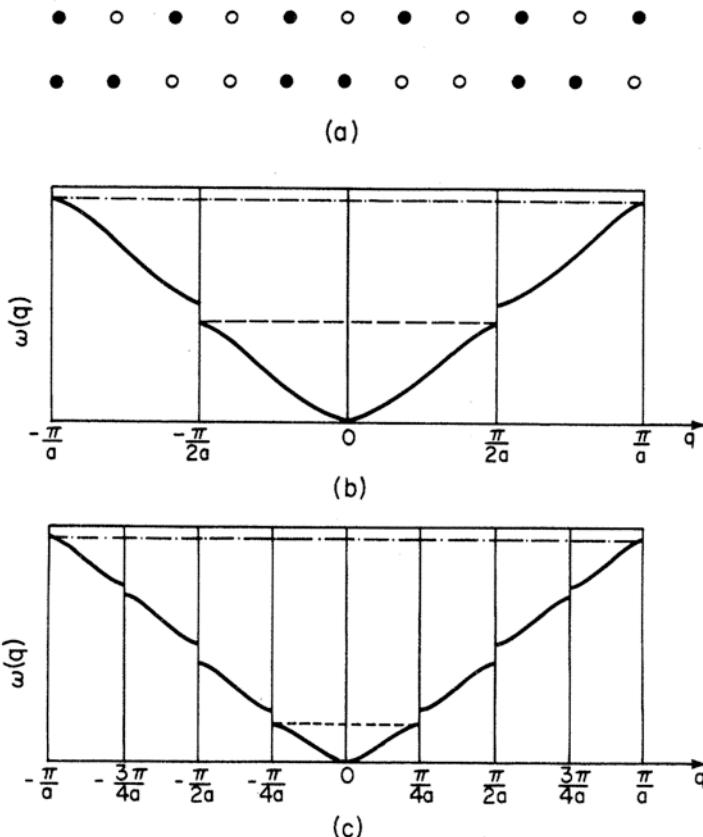
PHYSICAL REVIEW LETTERS

31 JANUARY 2000

## Thermal conductivity of superlattices

Shang Yuan Ren\* and John D. Dow

Department of Physics and Coordinated Science Laboratory, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801  
(Received 21 September 1981)

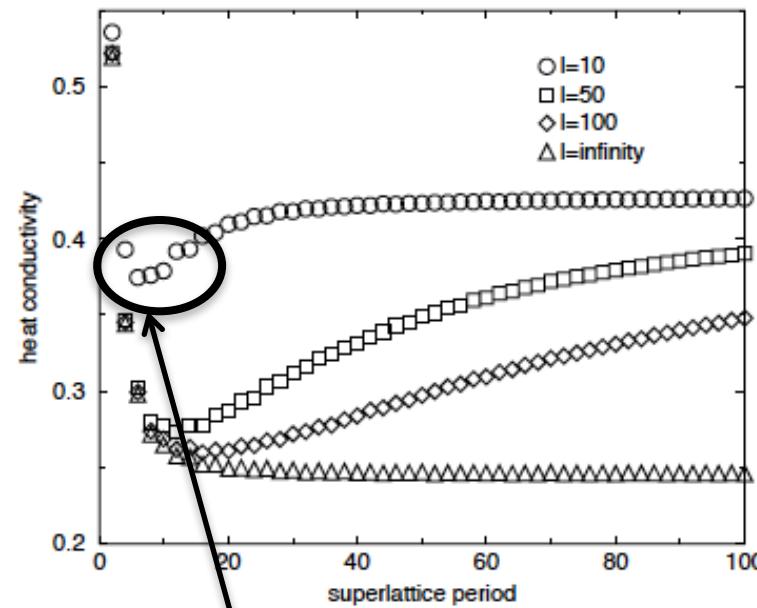


Interfacial periodicity can lead to “mini-band” formation

## Minimum Thermal Conductivity of Superlattices

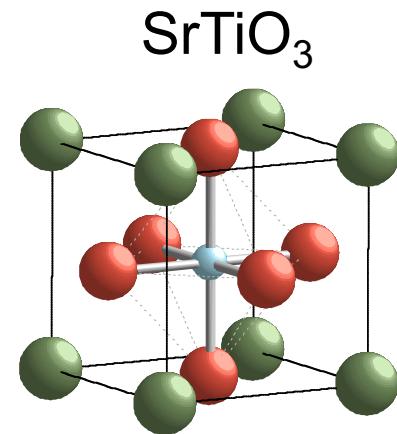
M. V. Simkin and G. D. Mahan

Department of Physics and Astronomy, University of Tennessee, Knoxville, Tennessee 37996-1200  
and Solid State Division, Oak Ridge National Laboratory, P.O. Box 2008, Oak Ridge, Tennessee 37831  
(Received 23 July 1999)

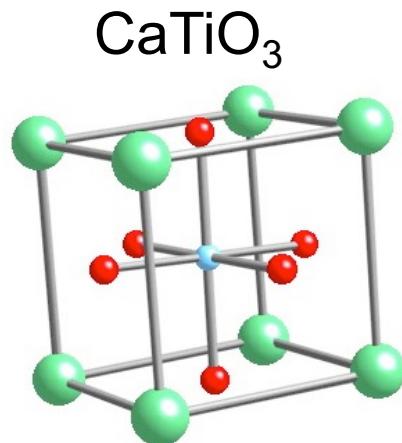


Mini-band formation leads to a minimum in the superlattice thermal conductivity

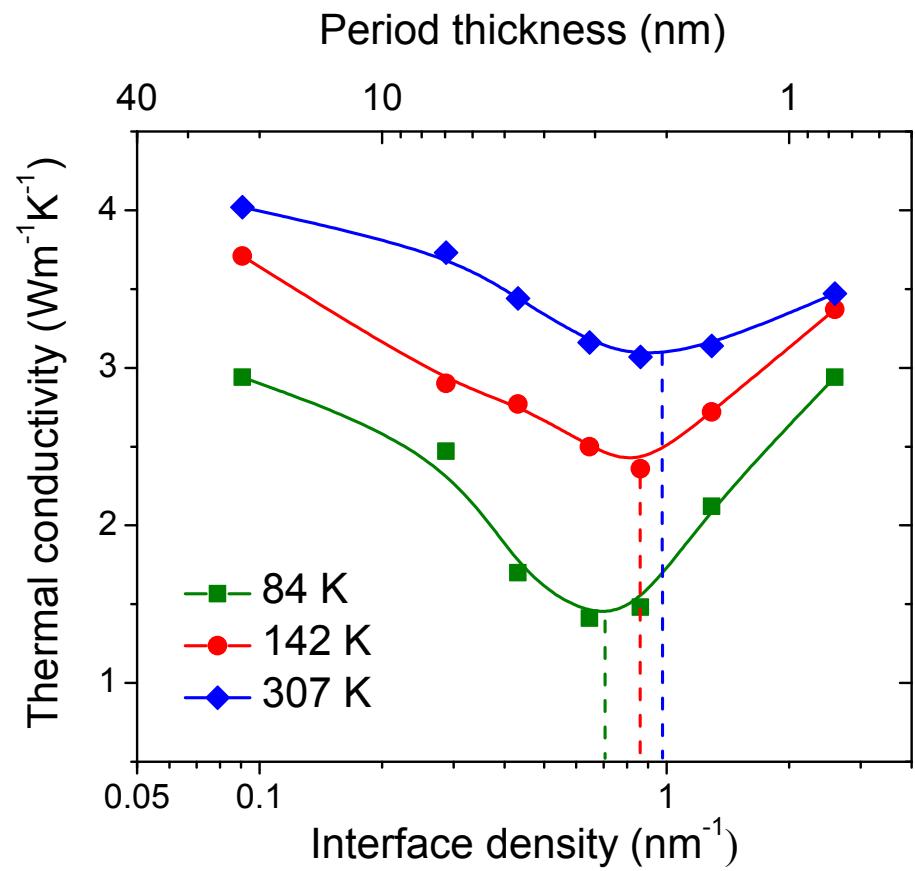
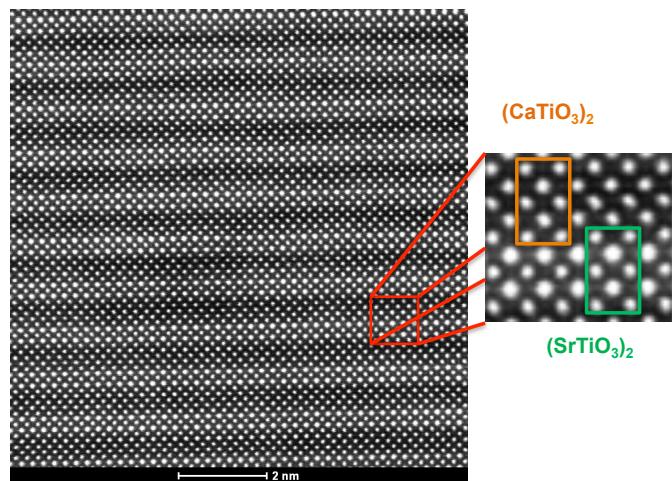
# Coherent transport in superlattices



$(a = 3.905 \text{ \AA})$   
 $\rho = 5.1 \text{ g/cc}$   
 $v_m = 5.41 \text{ km/s}$



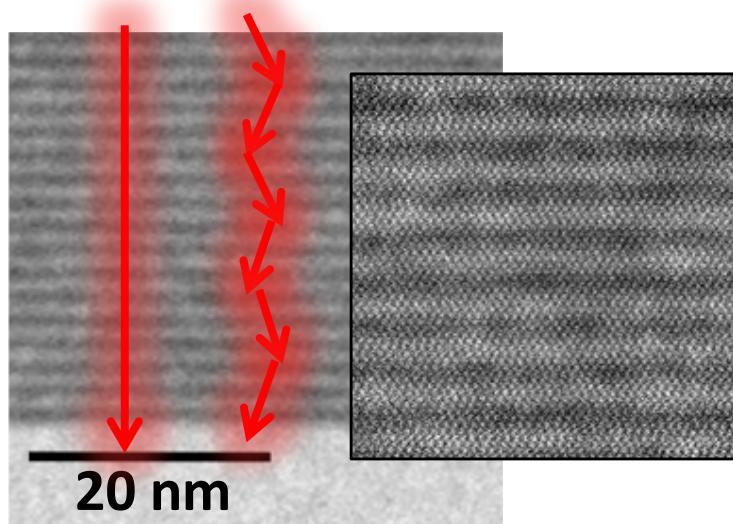
$(a_{pc} = 3.81 \text{ \AA})$   
 $\rho = 3.75 \text{ g/cc}$   
 $v_m = 5.71 \text{ km/s}$



**SL design to manipulate  
coherent phonon transport**  
*Nature Materials 13, 168 (2013)*

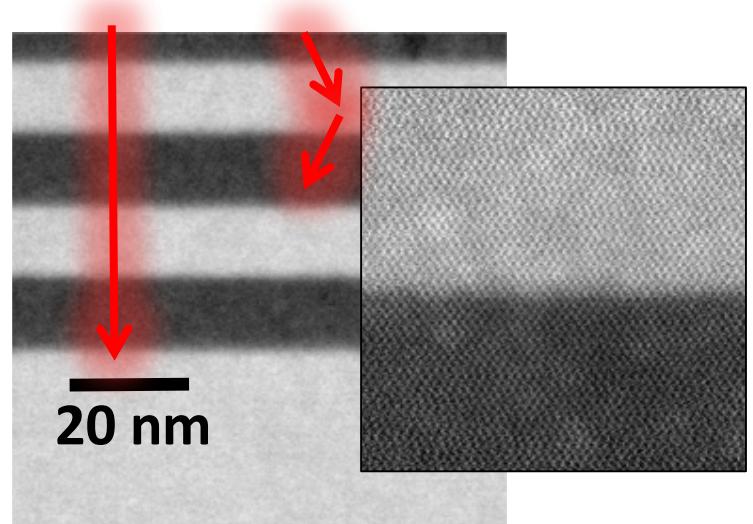
# Spectral phonon transport in SLs

(a)



$$d_{\text{SL}} = 2 \text{ nm}$$

(b)

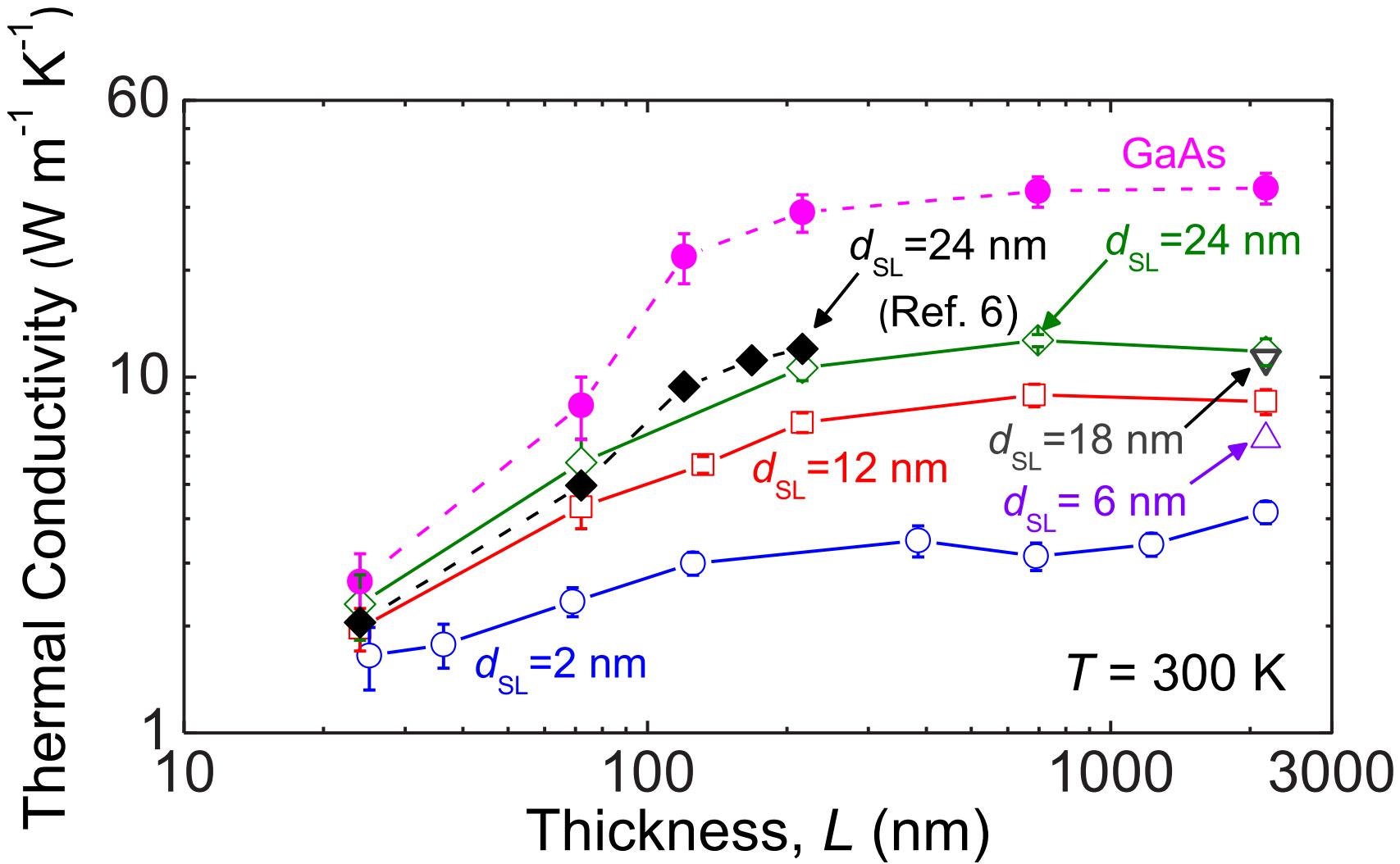


$$d_{\text{SL}} = 24 \text{ nm}$$

**Do all phonons really  
scatter at interfaces???**

**What if wavelength of  
phonon is  $\gg$  than  
periodicity?**

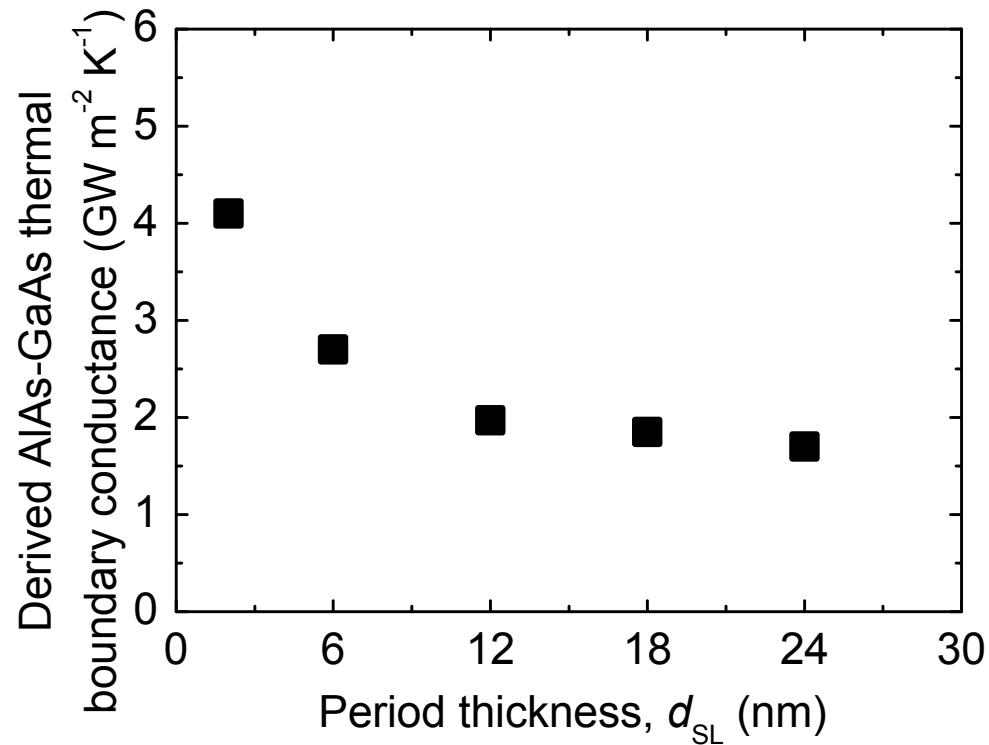
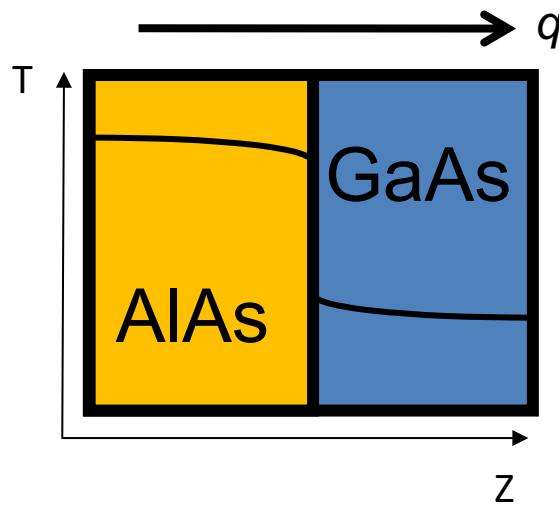
# Spectral phonon transport in SLs



Collaboration: G. Balakrishnan (UNM)  
*Phys. Rev. B* **97**, 085306

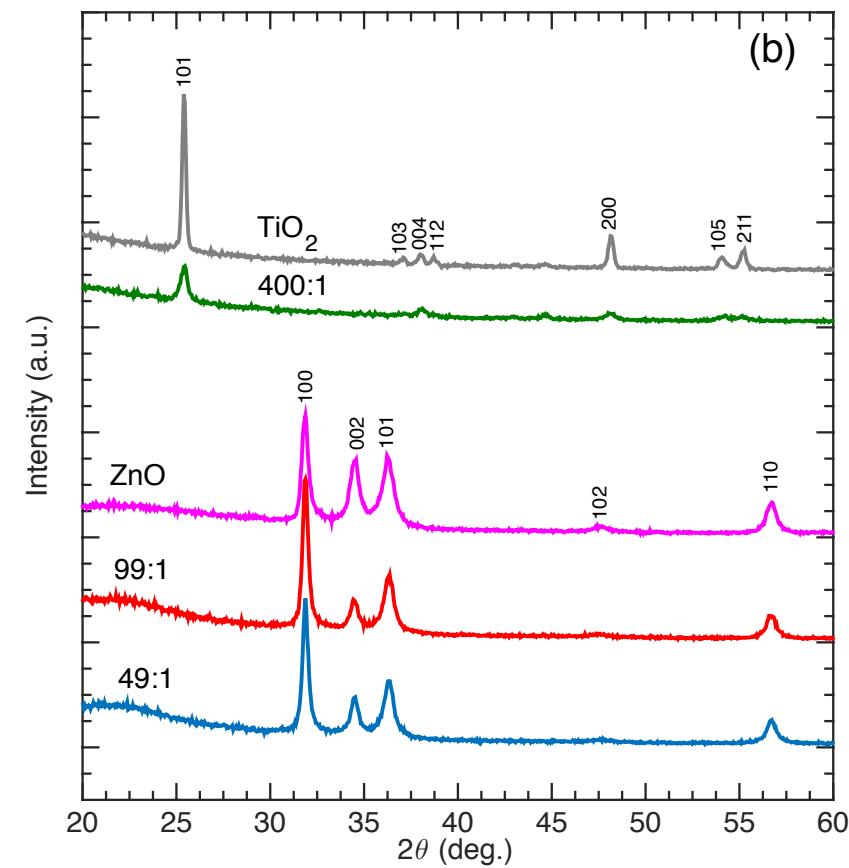
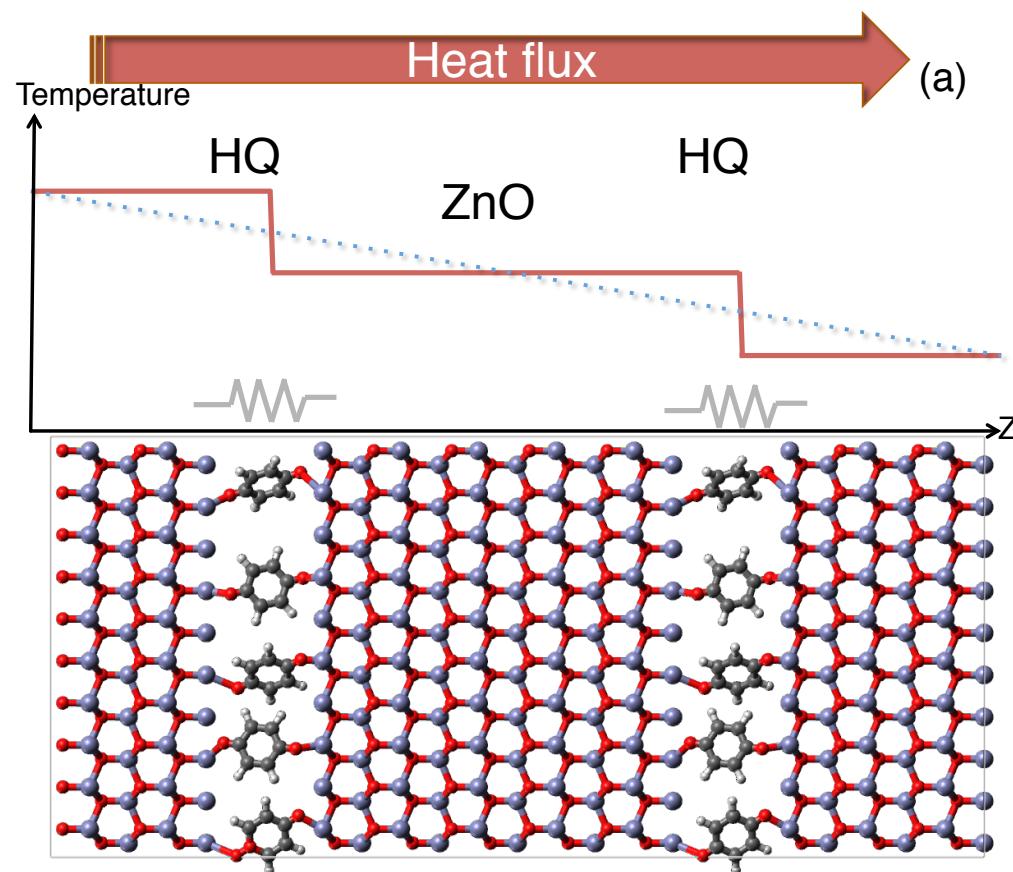
# TBC at AlAs/GaAs interface

$$q = h_K \Delta T = \frac{1}{R_K} \Delta T$$



Collaboration: G. Balakrishnan (UNM)  
*Phys. Rev. B* **97**, 085306

# Turn to molecular heterostructures

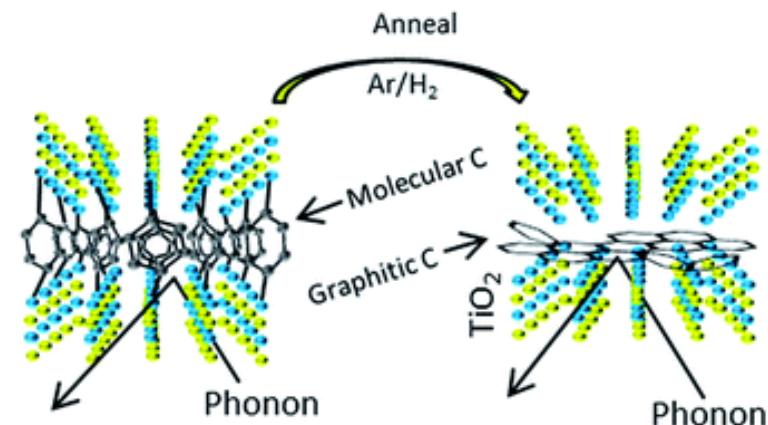
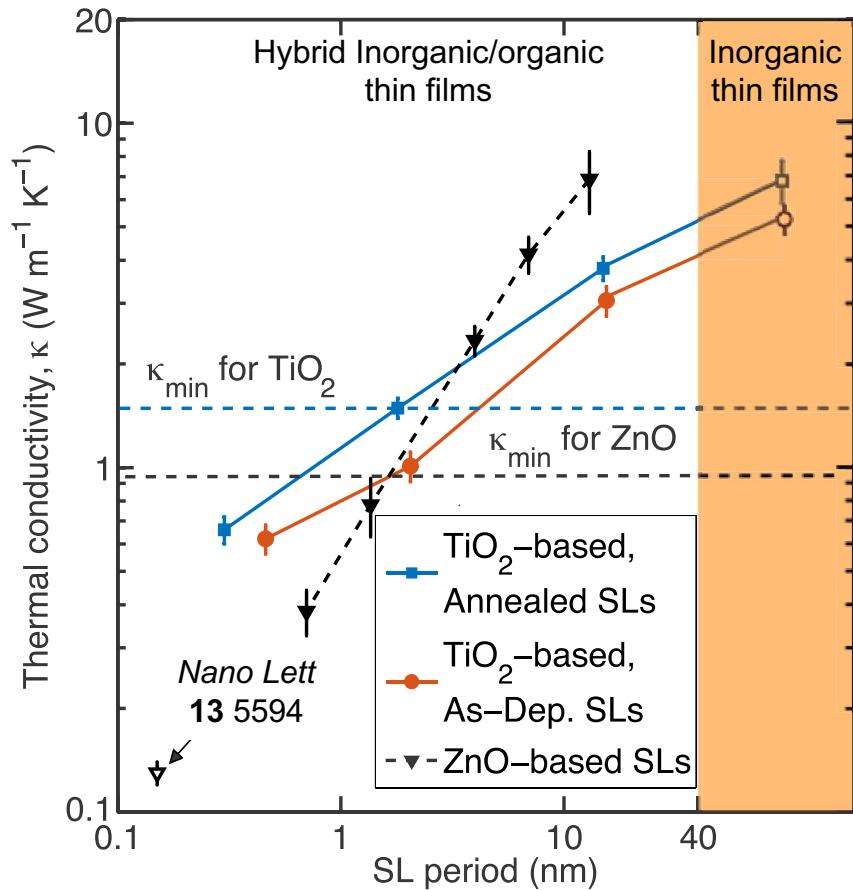


*PRB* 93, 115310; *PRB* 93, 024201

**Collaboration:** M. Karppinen (Aalto) – ALD/MLD growth

# Turn to molecular heterostructures

Phonon scattering at organic/inorganic interface can lead crystalline composites achieving  $\kappa$  less than amorphous phase



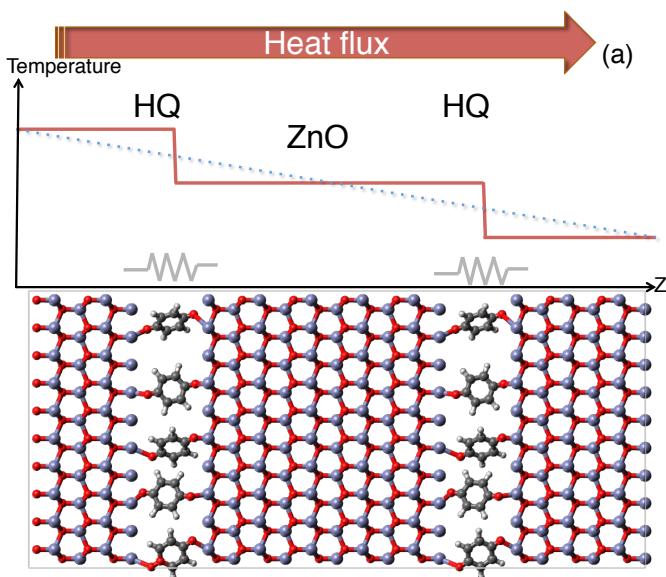
Reductions due to the thermal boundary conductance across organic/inorganic interfaces

PRB 93, 115310; PRB 93, 024201

Collaboration: M. Karppinen (Aalto) – ALD/MLD growth

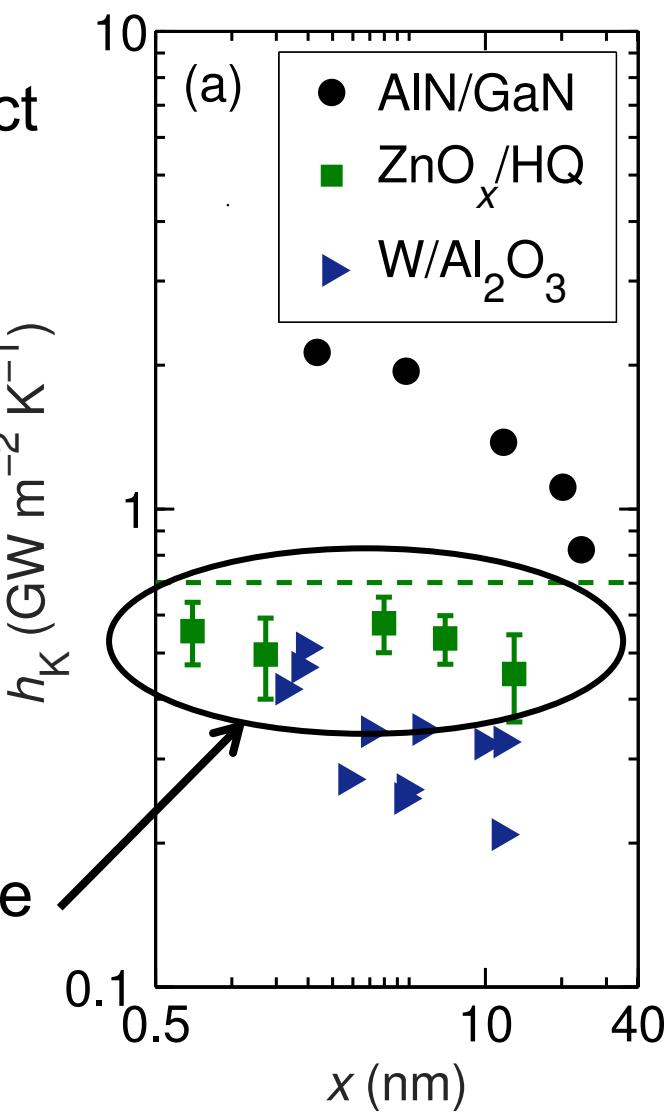
# Diffusive phonon scattering at organic/inorganic interface

Molecular interface causes all phonon modes to scatter at boundary (aka: no direct transmission of oxide modes across HQ)



$$q = h_K \Delta T = \frac{1}{R_K} \Delta T$$

Molecular interface constant with  $x$

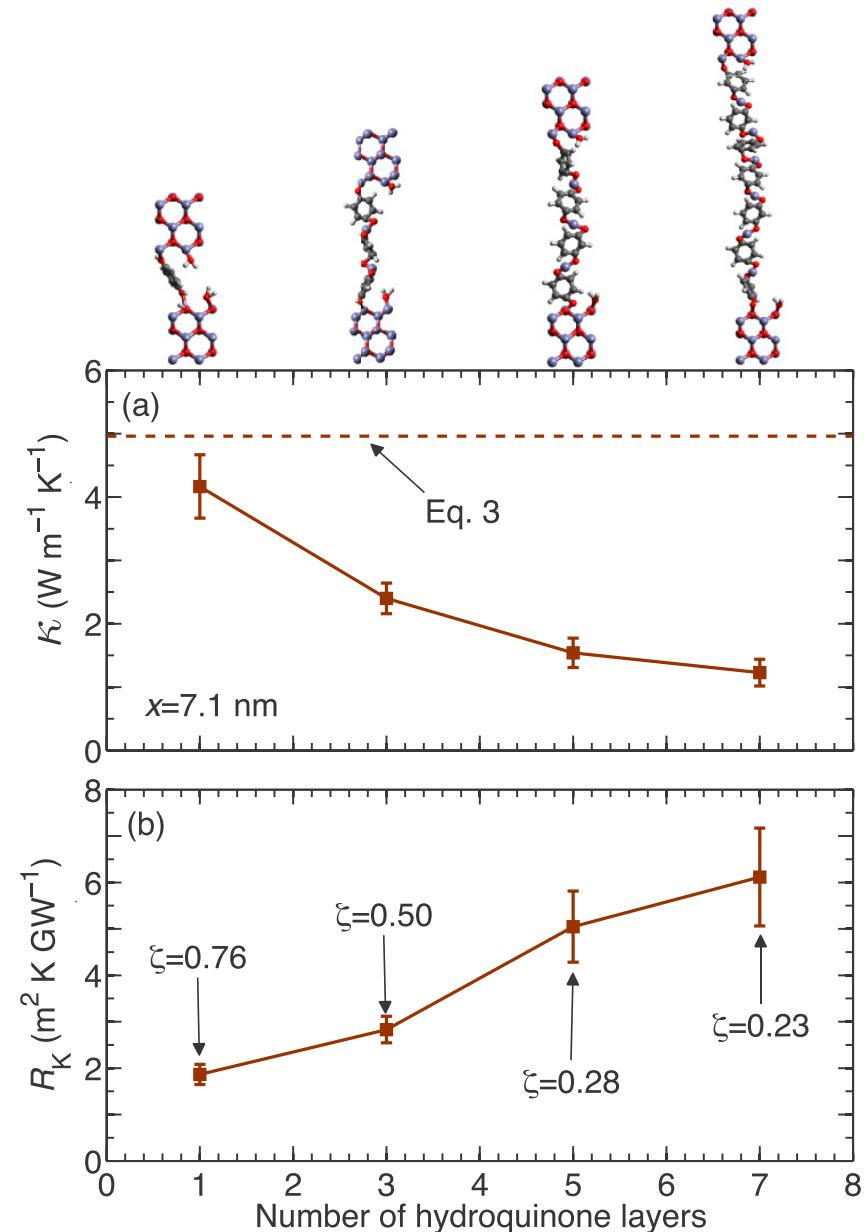


PRB 93, 115310; PRB 93, 024201

Collaboration: M. Karppinen (Aalto) – ALD/MLD growth

# Diffusive phonon scattering at organic/inorganic interface

Can use hybrid materials to study diffusive scattering in few molecule thick films

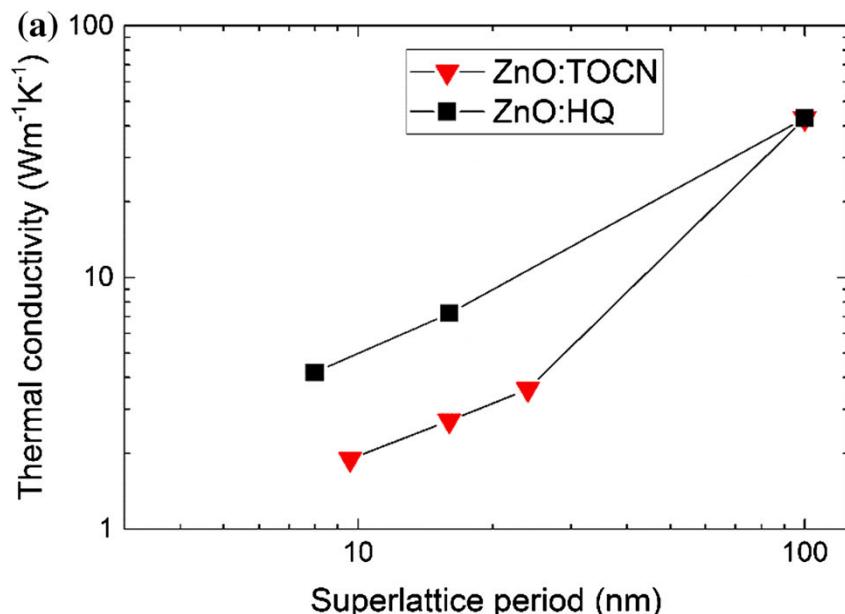
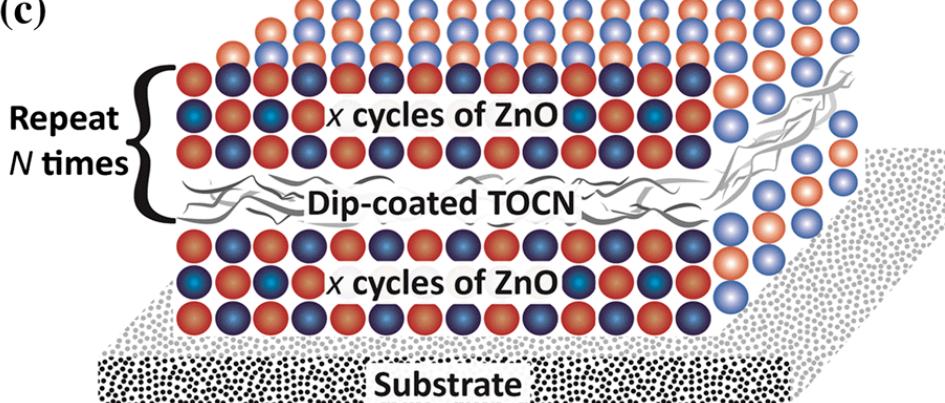


PRB 93, 115310; Collaboration: M. Karppinen (Aalto) – ALD/MLD growth

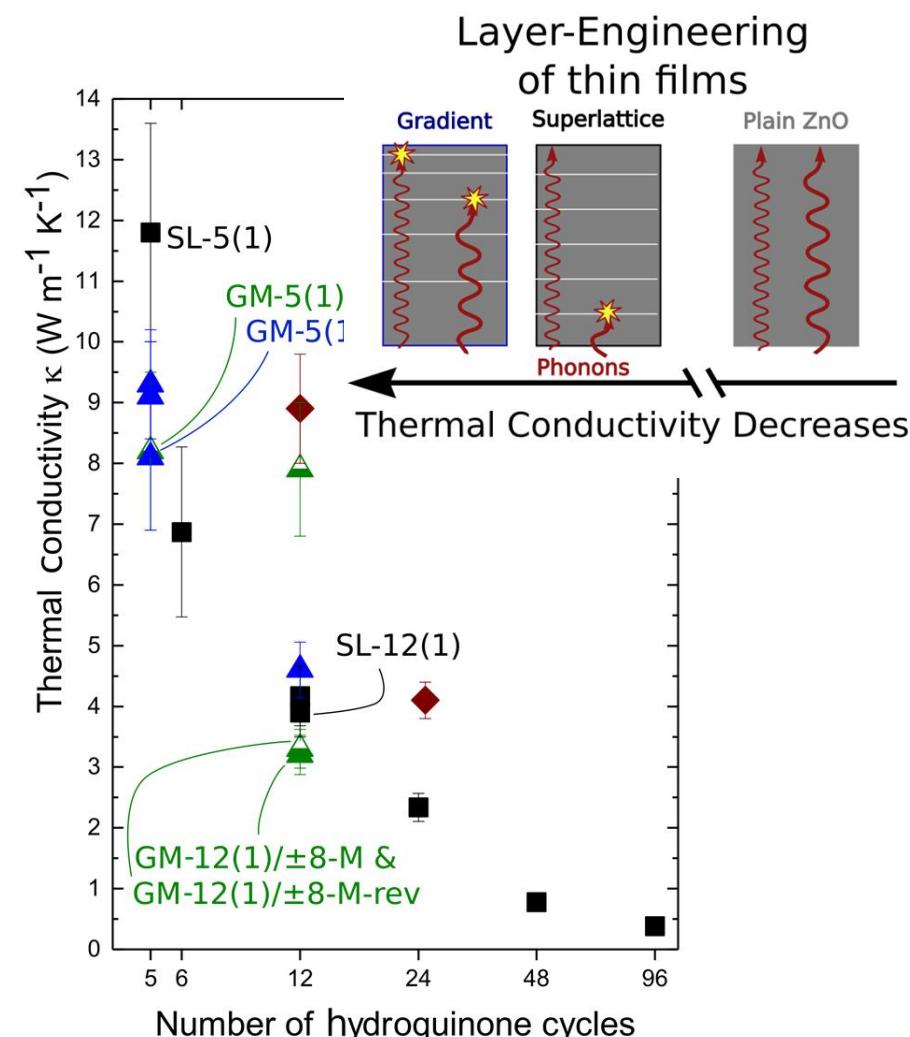
# Heat transport mechanisms in novel hybrid materials

## ZnO/nanocellulose

(c)



## A periodic hybrid “SLs”



# Conclusions/outlook – Heterogeneous material interfaces and engineering the chemical bond can lead to novel regimes of vibrational heat transport in materials

