



SCHOOL of ENGINEERING & APPLIED SCIENCE  
UNIVERSITY of VIRGINIA

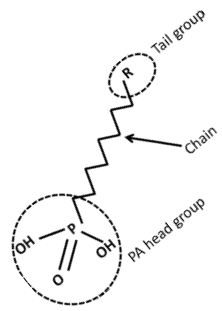
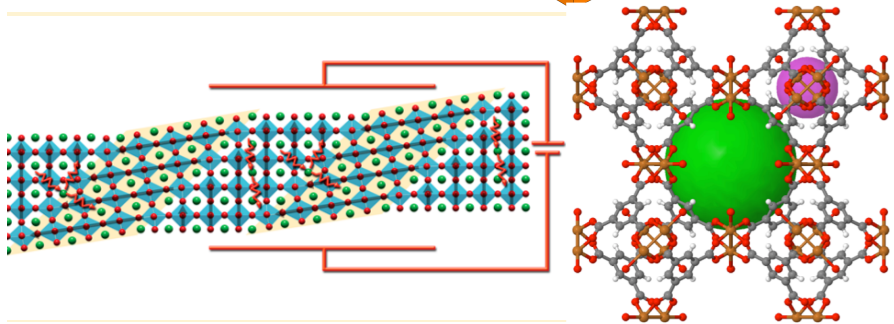
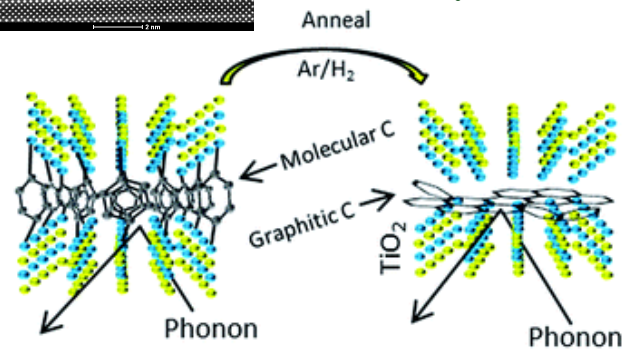
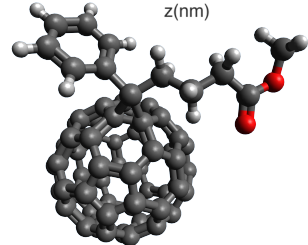
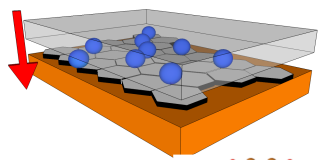
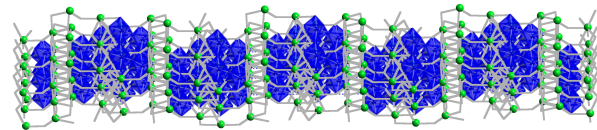
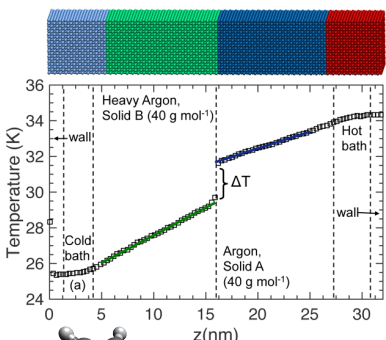
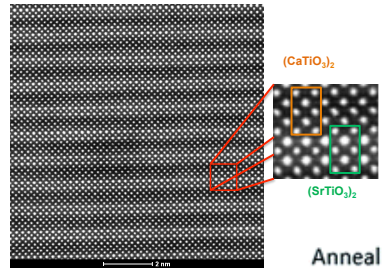
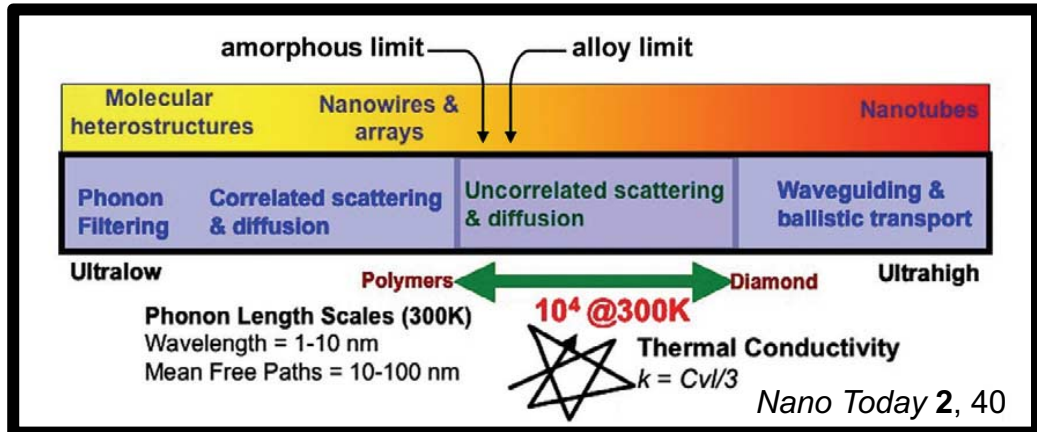
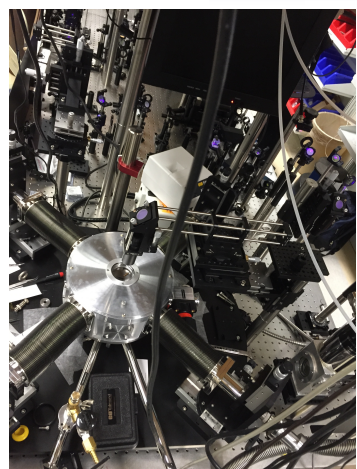
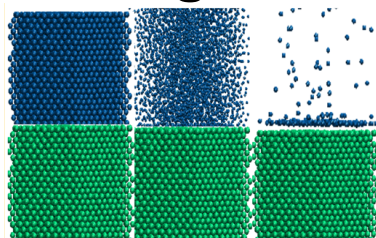
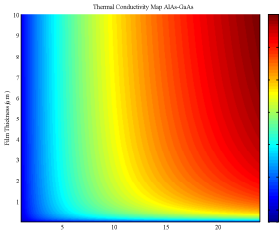
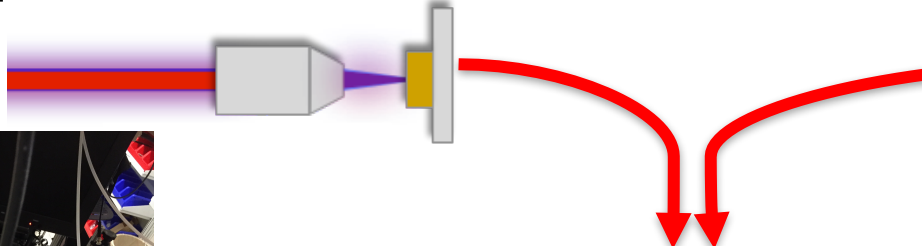
# Thermal transport processes in functional organic and organic/inorganic nanomaterials: The role of the chemical bond



**Patrick E. Hopkins**

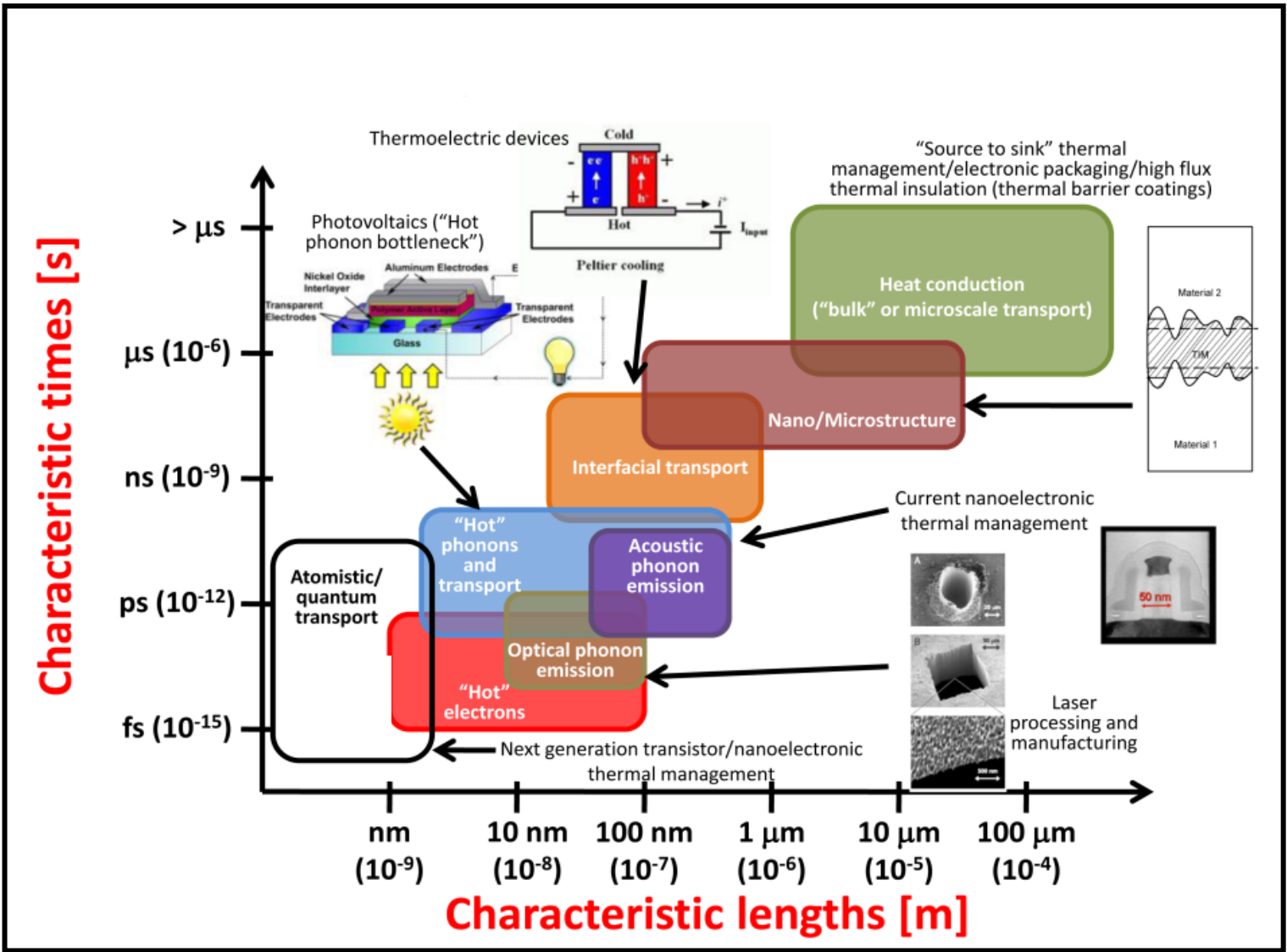
Associate Professor  
Dept. Mech. & Aero. Eng.  
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[phopkins@virginia.edu](mailto:phopkins@virginia.edu)  
[patrickehopkins.com](http://patrickehopkins.com)

Experiments and Simulations in Thermal Engineering

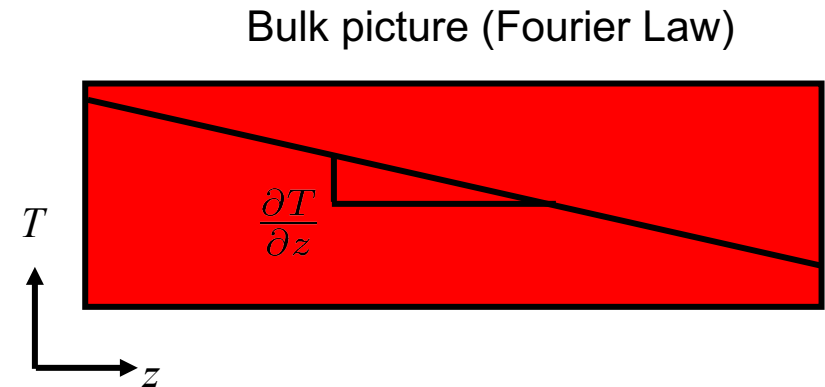
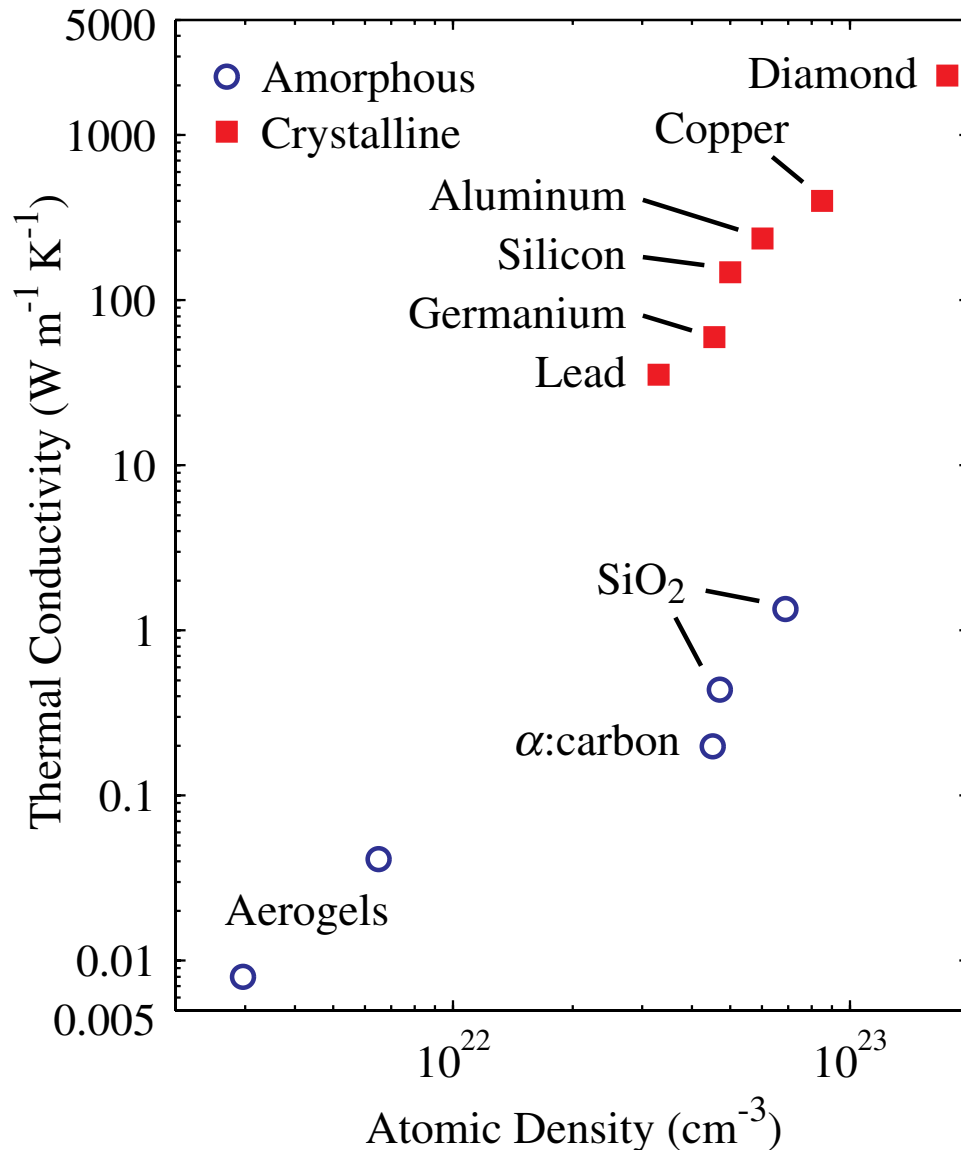




Engineering energy transport, conversion and storage in materials over multiple time and length scales



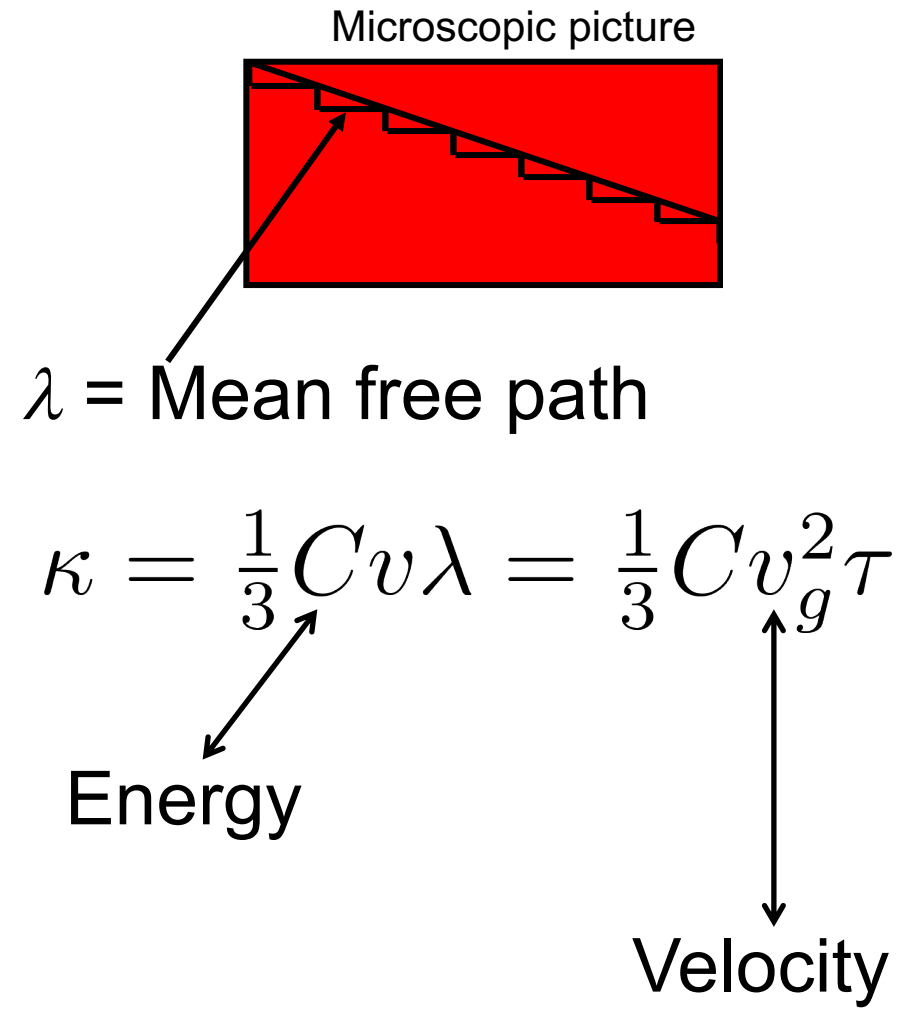
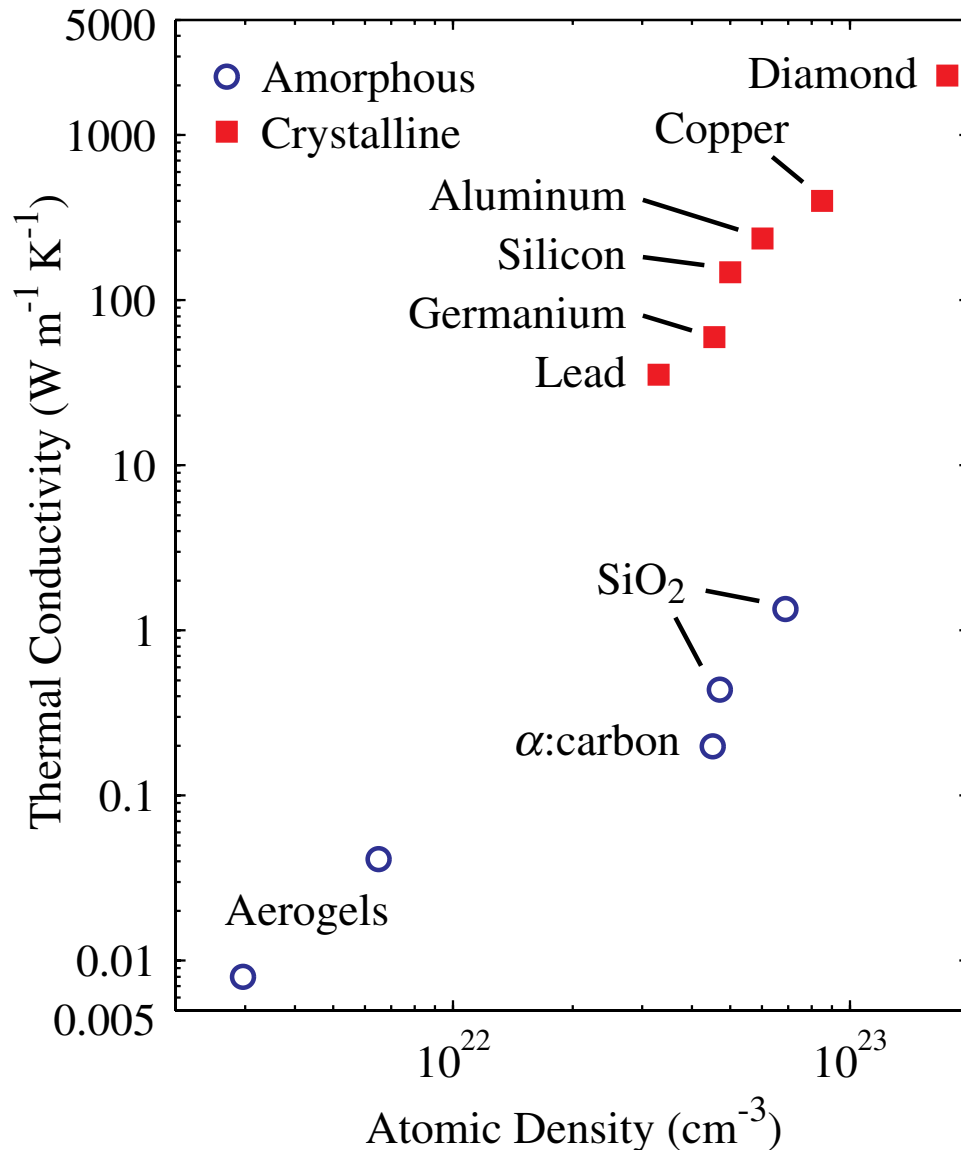
# Thermal conductivity of materials – Traditional picture



$$q_z = -\kappa \frac{\partial T}{\partial z}$$

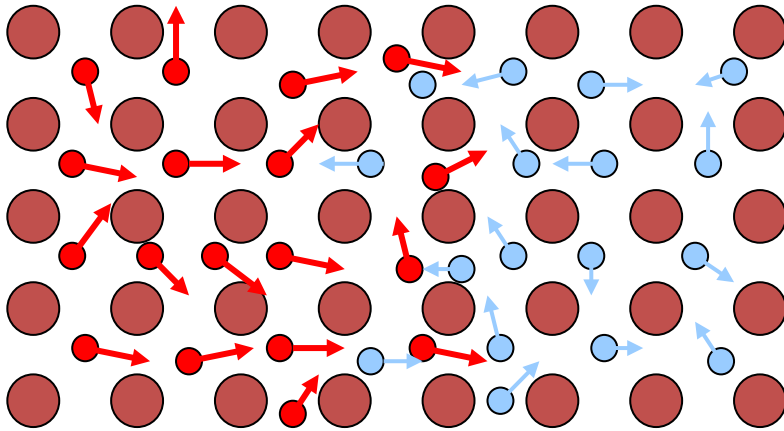


# Thermal conductivity of materials – Kinetic Theory picture



# Thermal conductivity of materials - nanoscopic

Diffusion of “hot” electrons



Metals:

Free electrons are the dominant energy carriers in metals, ballistic velocity  $\sim 10^6$  m/s

Electron carrier density:

in metals  $\sim 10^{23}$  cm $^{-3}$

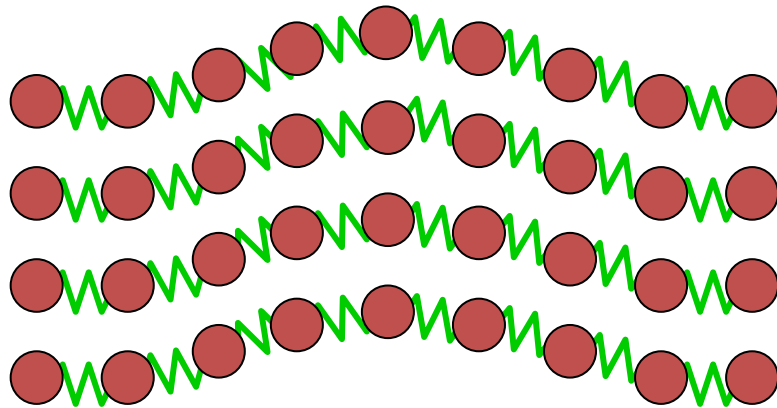
in semiconductors  $\sim 10^{18}$  cm $^{-3}$



atom

Red dot represents “hot” free electron

Blue dot represents “cold” free electron



Phonon propagation



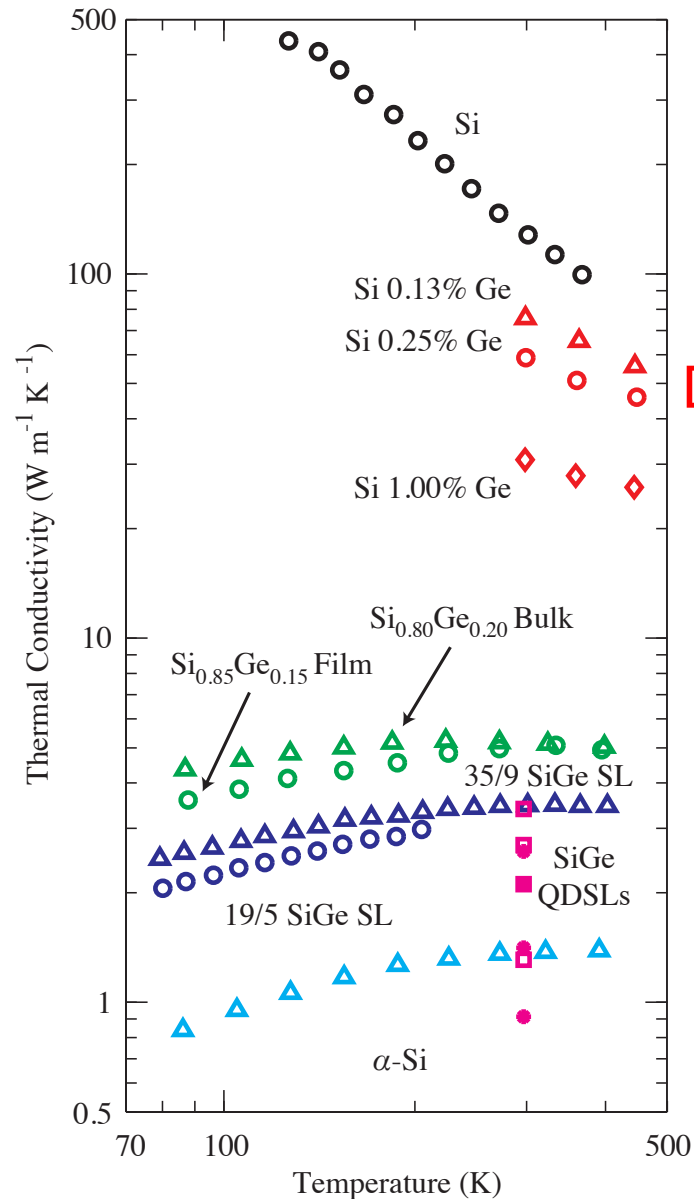
Semiconductors:

Phonons (lattice vibrations) are the dominant energy carriers in semiconductors, velocity  $\sim 10^3$  m/s

$$\kappa = \frac{1}{3} C v \lambda = \frac{1}{3} C v_g^2 \tau$$



# Thermal conductivity of materials - nanoscopic



Dilute alloy

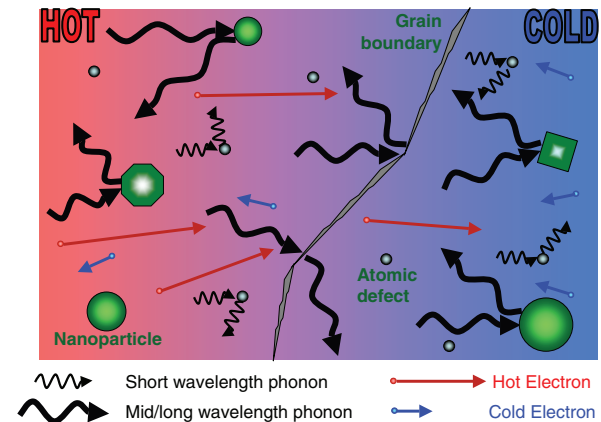
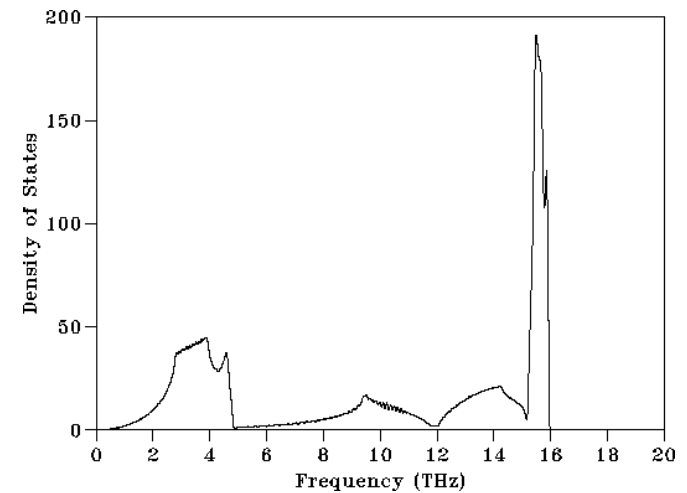
Alloy

Interface

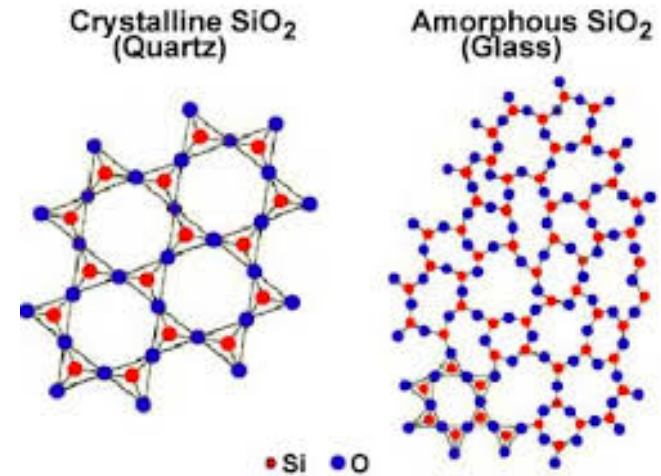
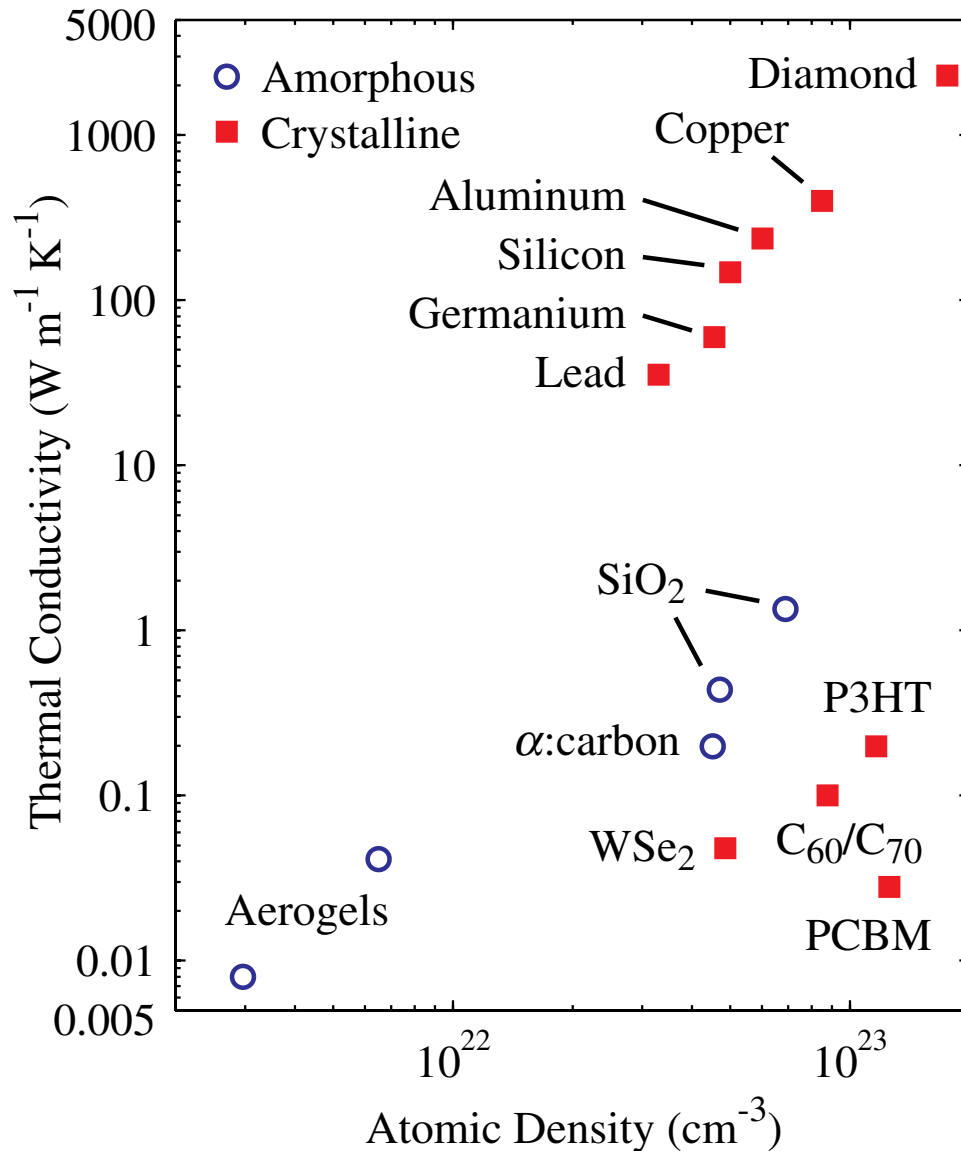
Disorder

$$\kappa = \frac{1}{3} C v \lambda = \frac{1}{3} C v_g^2 \tau$$

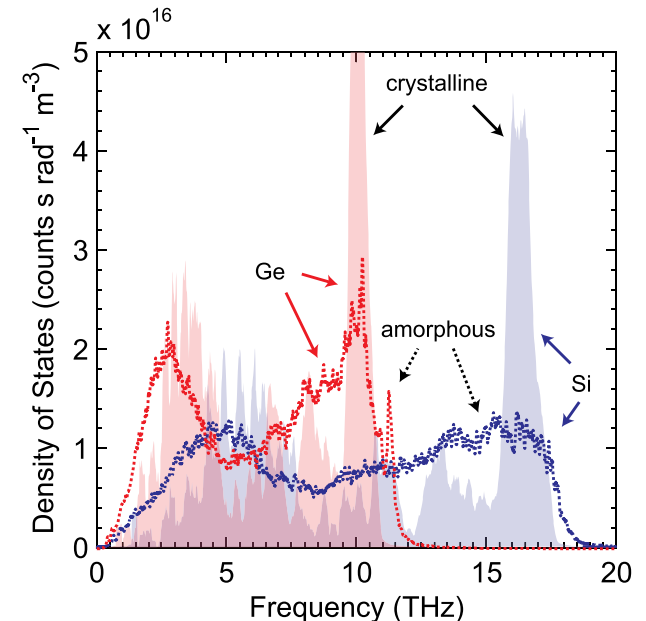
$$\tau = f(\tau_{\text{intrinsic}}, \tau_{\text{impurity}}, \tau_{\text{boundary}})$$



# Thermal conductivity of non-crystalline solids



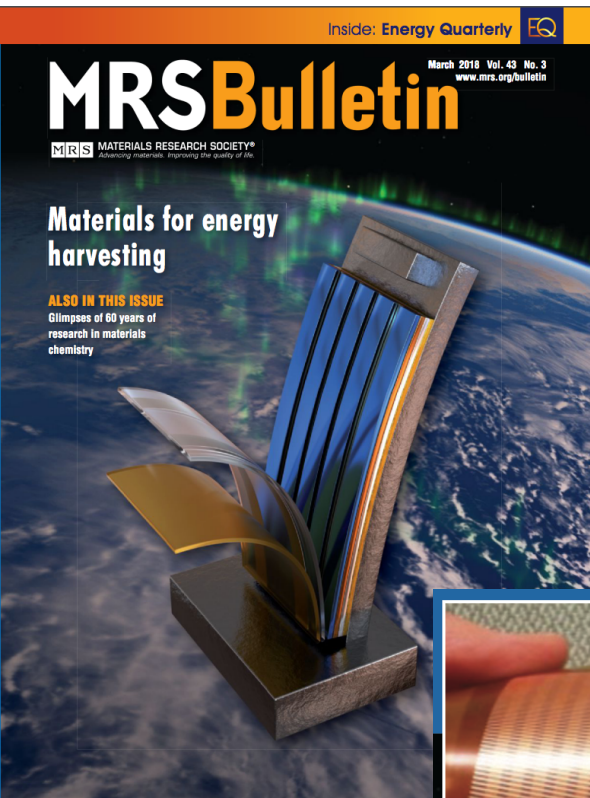
*NDT Resource Center*



*J. Appl. Phys.* **118**, 165303



# A future for energy recovery: wearable thermoelectrics



## March 2018 MRS Bulletin

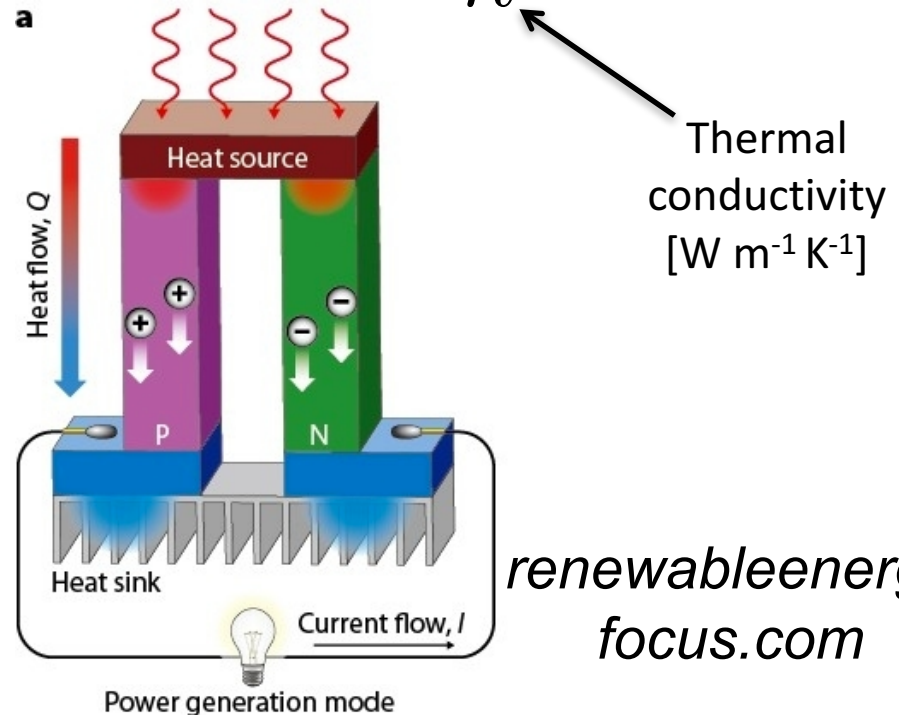
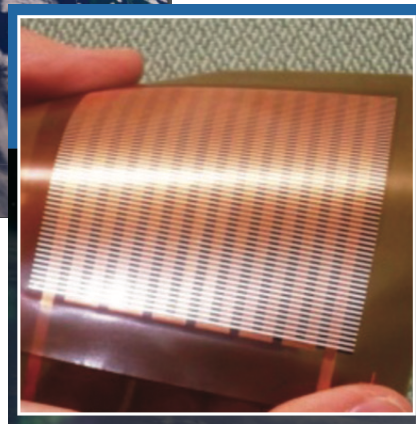
Seebeck coefficient [ $\text{V K}^{-1}$ ]

Electrical conductivity [ $\text{S m}^{-1}$ ]

Temperature [K]

Figure of merit [unitless]

$$ZT = \frac{S^2 \sigma T}{\kappa}$$



## Wearable and flexible thermoelectrics for energy harvesting

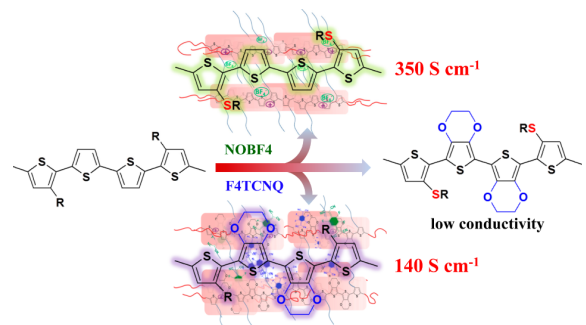
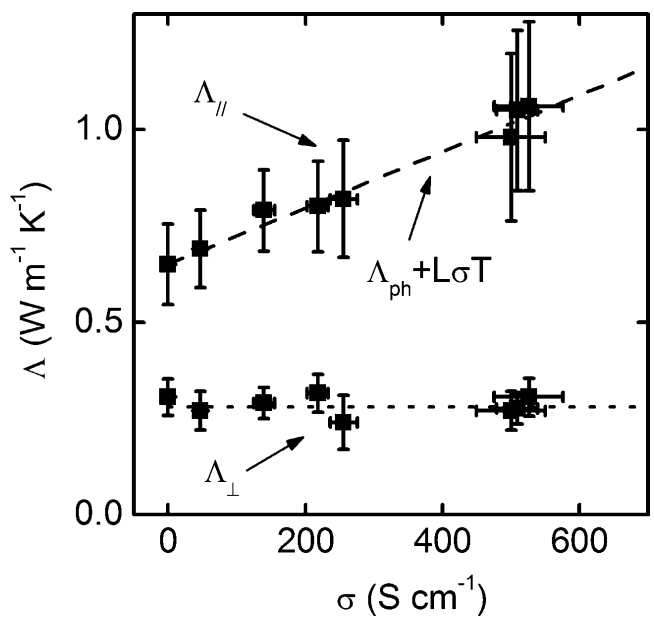
Ruoming Tian, Chunlei Wan, Naoyuki Hayashi, Toshiaki Aoi, and Kunihiro Koumoto

*renewableenergy  
focus.com*

# Engineering thermal conductivity in polymeric TE materials

$$ZT = \frac{S^2 \sigma T}{\kappa}$$

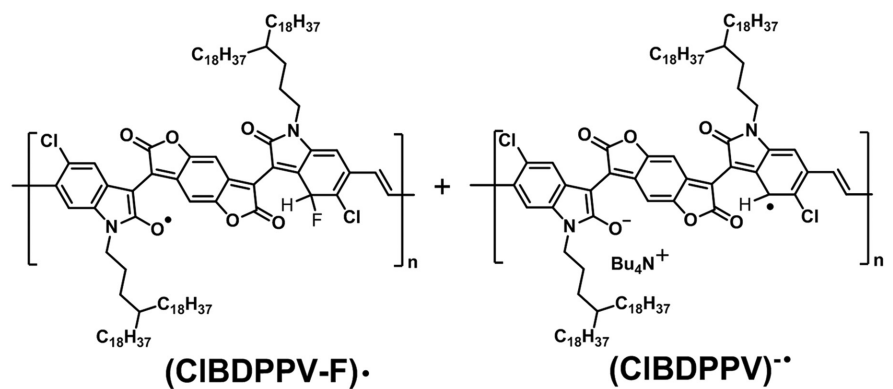
Macromolecules **48**, 585



J | A | C | S  
JOURNAL OF THE AMERICAN CHEMICAL SOCIETY  
Article  
pubs.acs.org/JACS

## Modification of the Poly(bisdodecylquaterthiophene) Structure for High and Predominantly Nonionic Conductivity with Matched Dopants

Hui Li,<sup>†</sup> Mallory E. DeCoster,<sup>\*</sup> Robert M. Ireland,<sup>†</sup> Jian Song,<sup>†</sup> Patrick E. Hopkins,<sup>‡</sup> and Howard E. Katz<sup>\*,†,§</sup>



COMMUNICATION  
Thermoelectrics  
ADVANCED MATERIALS  
www.advmat.de

## High Conductivity and Electron-Transfer Validation in an n-Type Fluoride-Anion-Doped Polymer for Thermoelectrics in Air

Xingang Zhao, Deepa Madan,<sup>\*</sup> Yan Cheng, Jiawang Zhou, Hui Li, Susanna M. Thon, Arthur E. Bragg, Mallory E. DeCoster, Patrick E. Hopkins, and Howard E. Katz<sup>\*</sup>



# Outline

- **TDTR: Measurement of thermal conductivity of thin films and thermal resistance across interfaces**
- Weakly bonded solids: new lower limits to thermal conductivity
- Functionalized interfaces at graphene contacts: tuning heat and electrical transport via the interfacial bond
- Heat transport across single molecule interfaces: when does a molecule become a defect?
- Molecular interfaces in organic/inorganic composites: diffusive scattering via the vibron-phonon interaction

# Measuring heat flow in thin films: TDTR

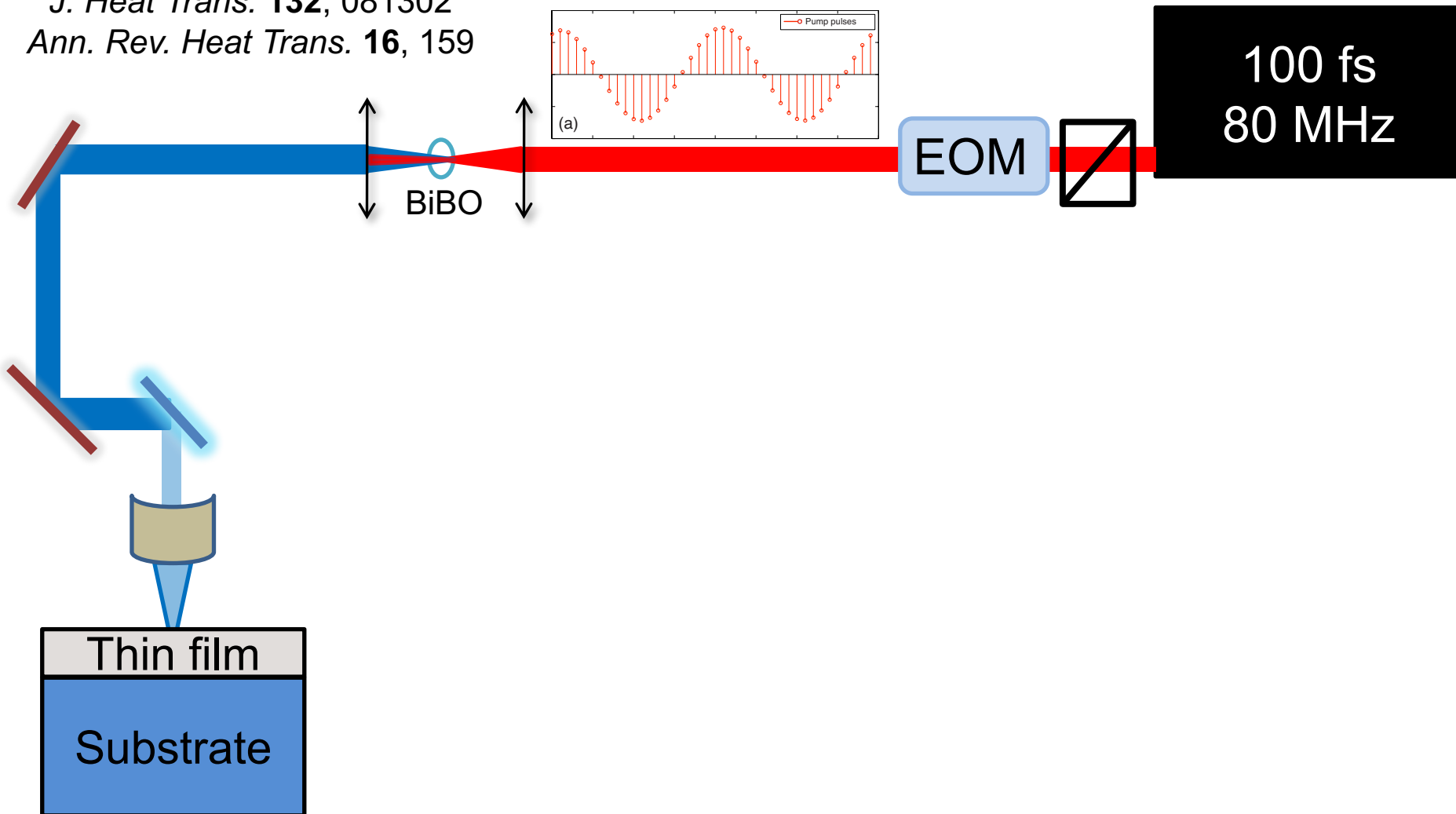
## TDTR Reviews and Analyses

*Rev. Sci. Instr.* **75**, 5119

*Rev. Sci. Instr.* **79**, 114902

*J. Heat Trans.* **132**, 081302

*Ann. Rev. Heat Trans.* **16**, 159



# Measuring heat flow in thin films: TDTR

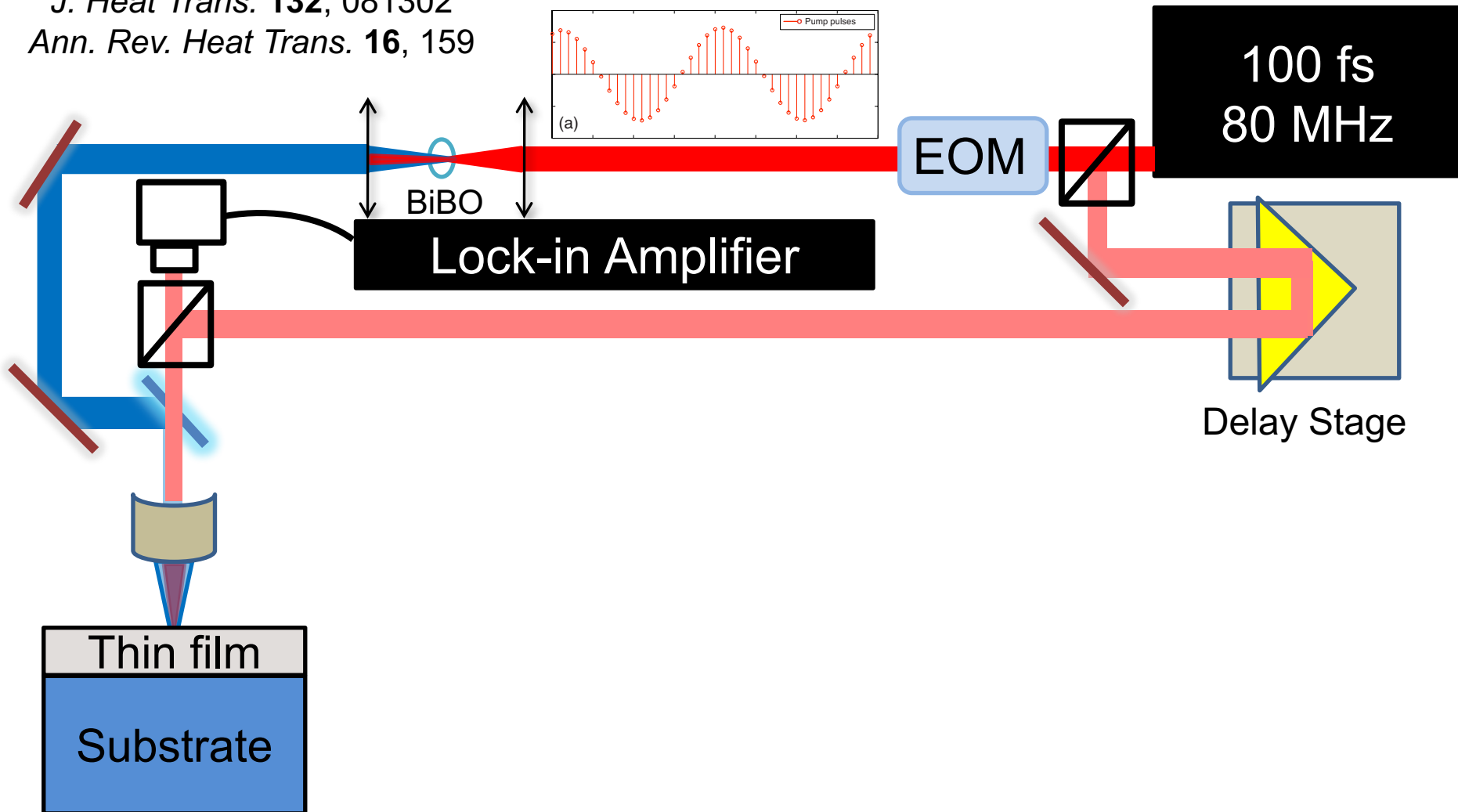
## TDTR Reviews and Analyses

*Rev. Sci. Instr.* **75**, 5119

*Rev. Sci. Instr.* **79**, 114902

*J. Heat Trans.* **132**, 081302

*Ann. Rev. Heat Trans.* **16**, 159



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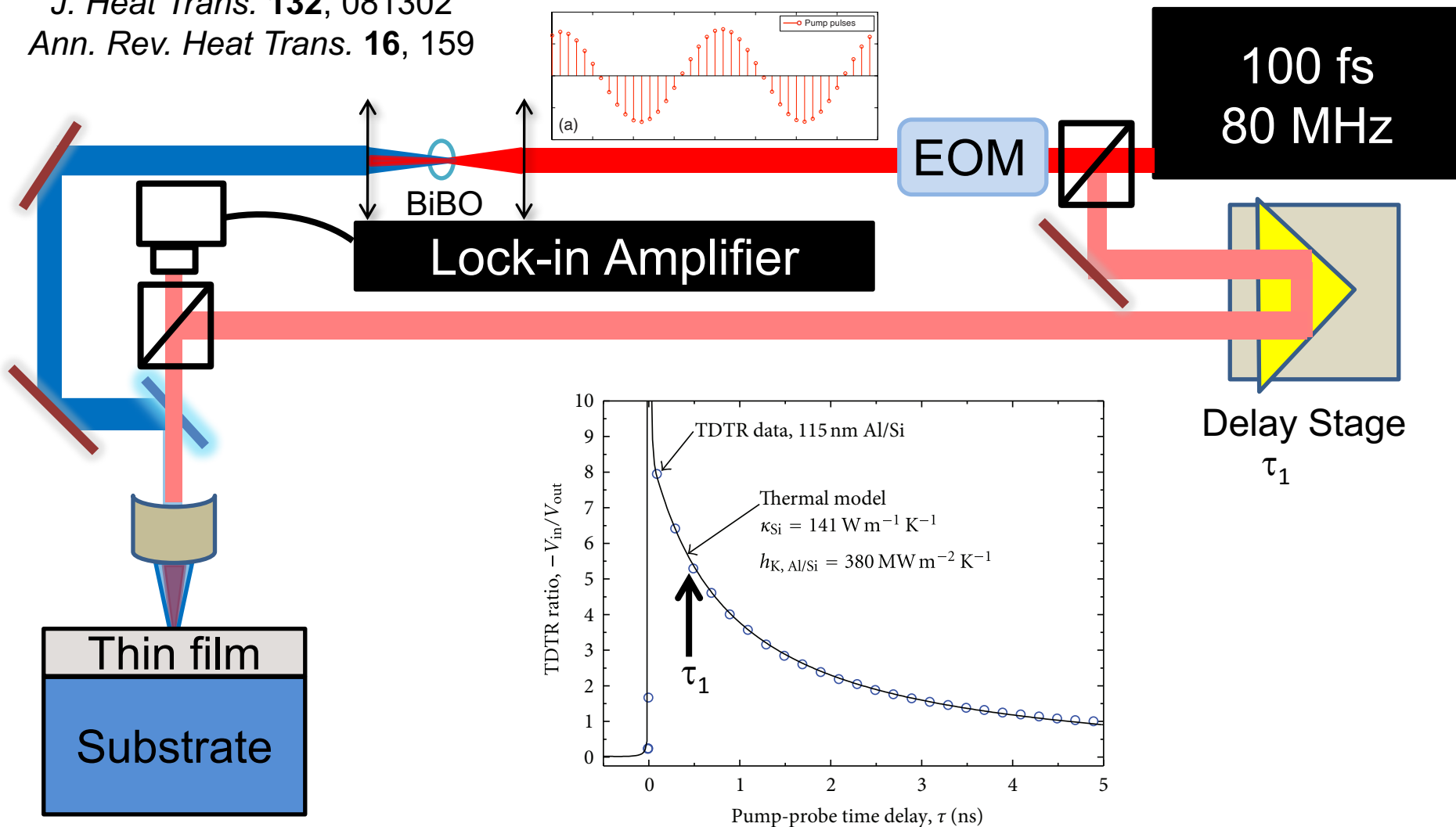
## TDTR Reviews and Analyses

*Rev. Sci. Instr.* **75**, 5119

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# Measuring heat flow in thin films: TDTR

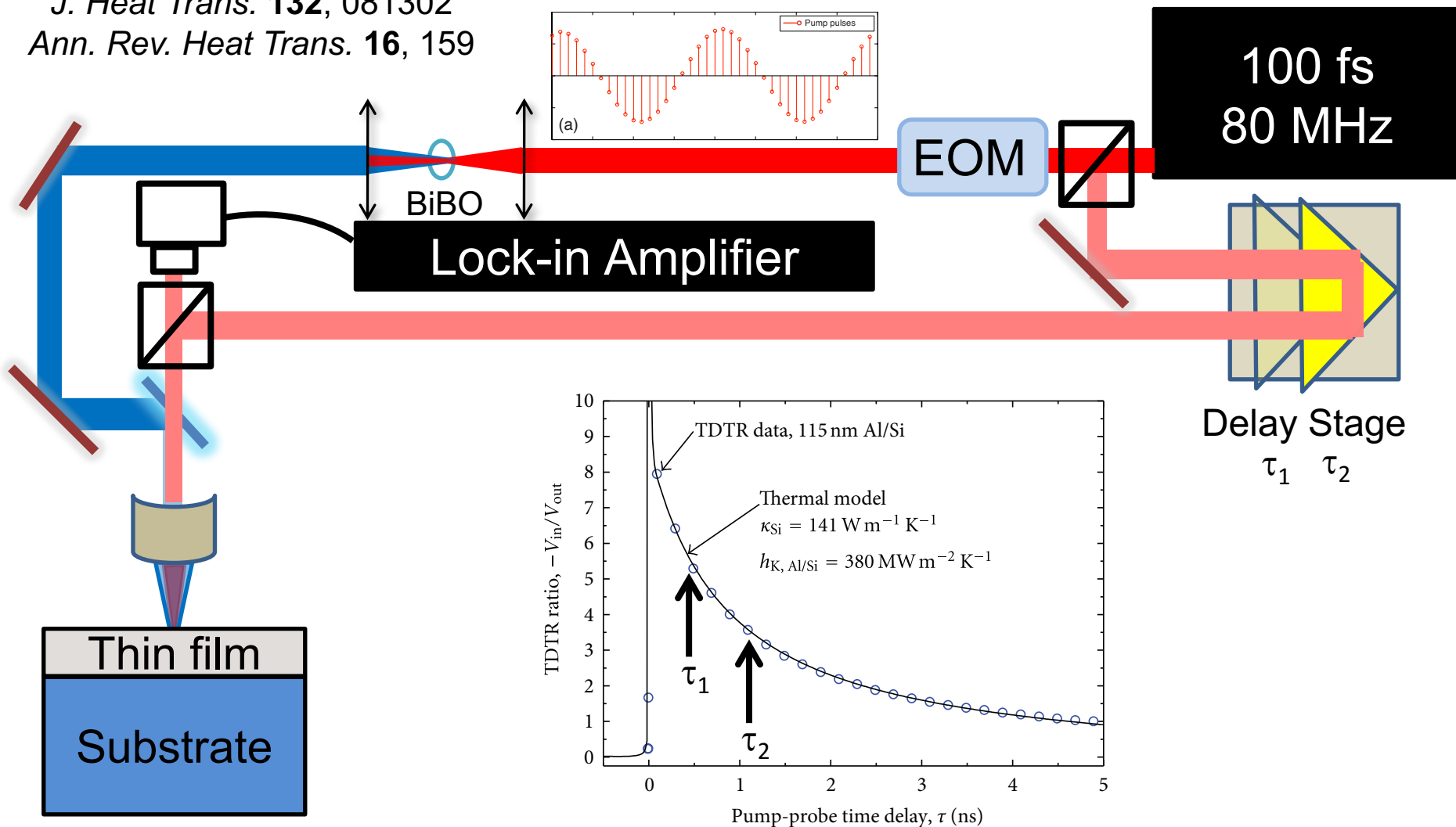
## TDTR Reviews and Analyses

*Rev. Sci. Instr.* **75**, 5119

*Rev. Sci. Instr.* **79**, 114902

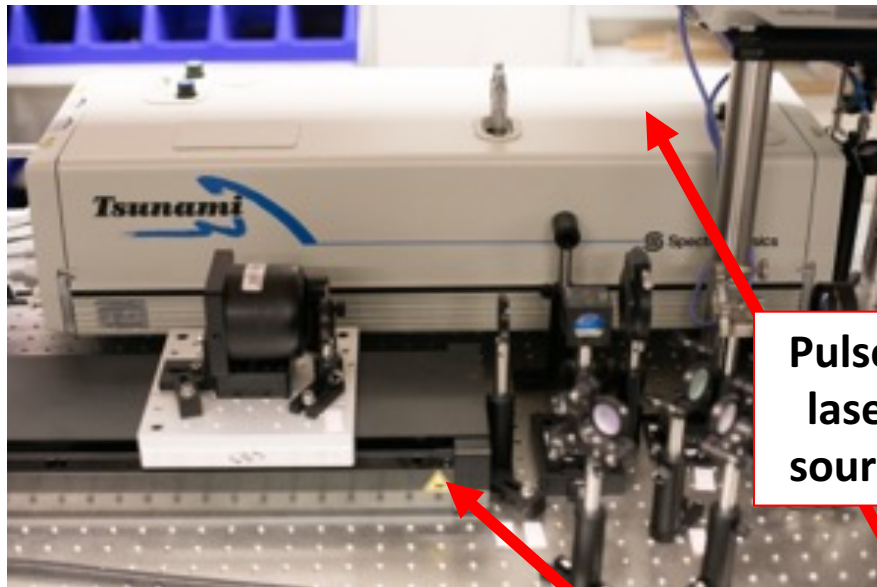
*J. Heat Trans.* **132**, 081302

*Ann. Rev. Heat Trans.* **16**, 159



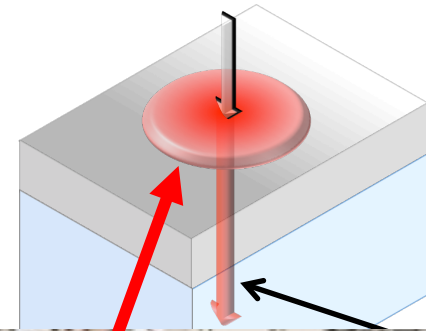


# Measuring heat flow in thin films: TDTR

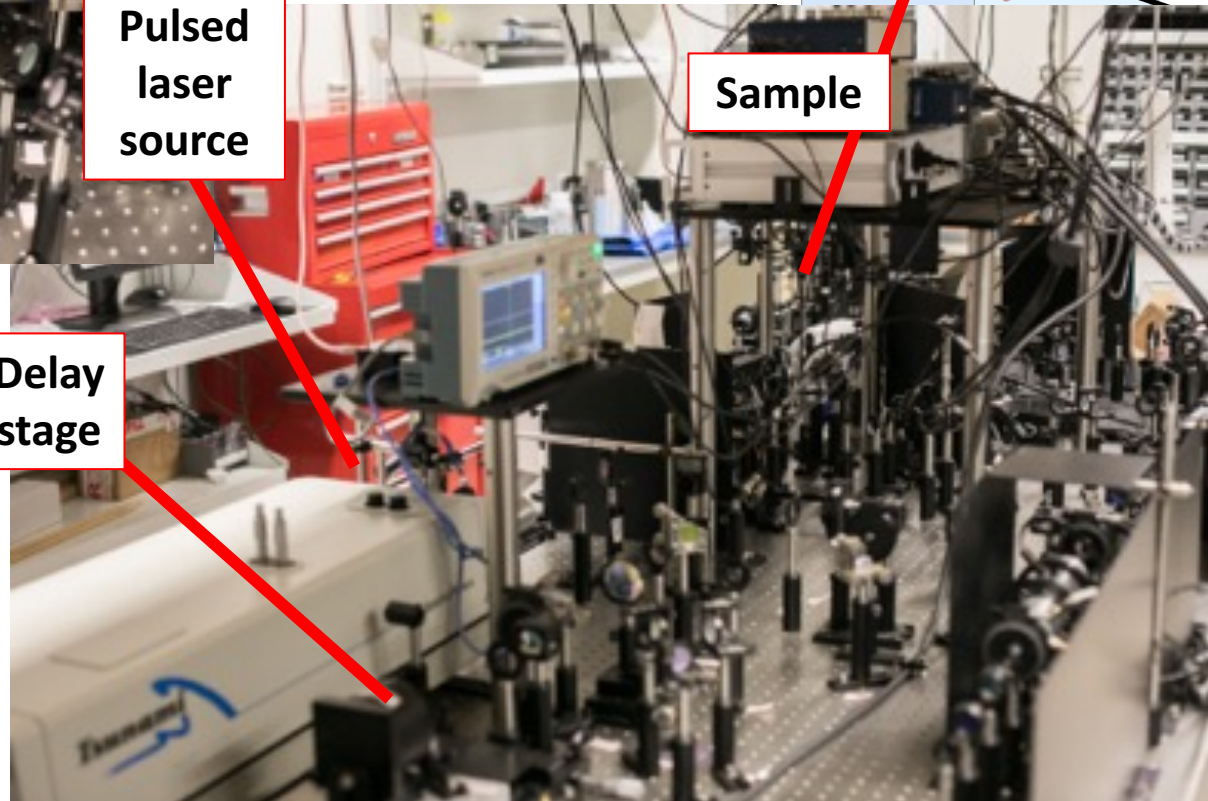


Pulsed  
laser  
source

Delay  
stage

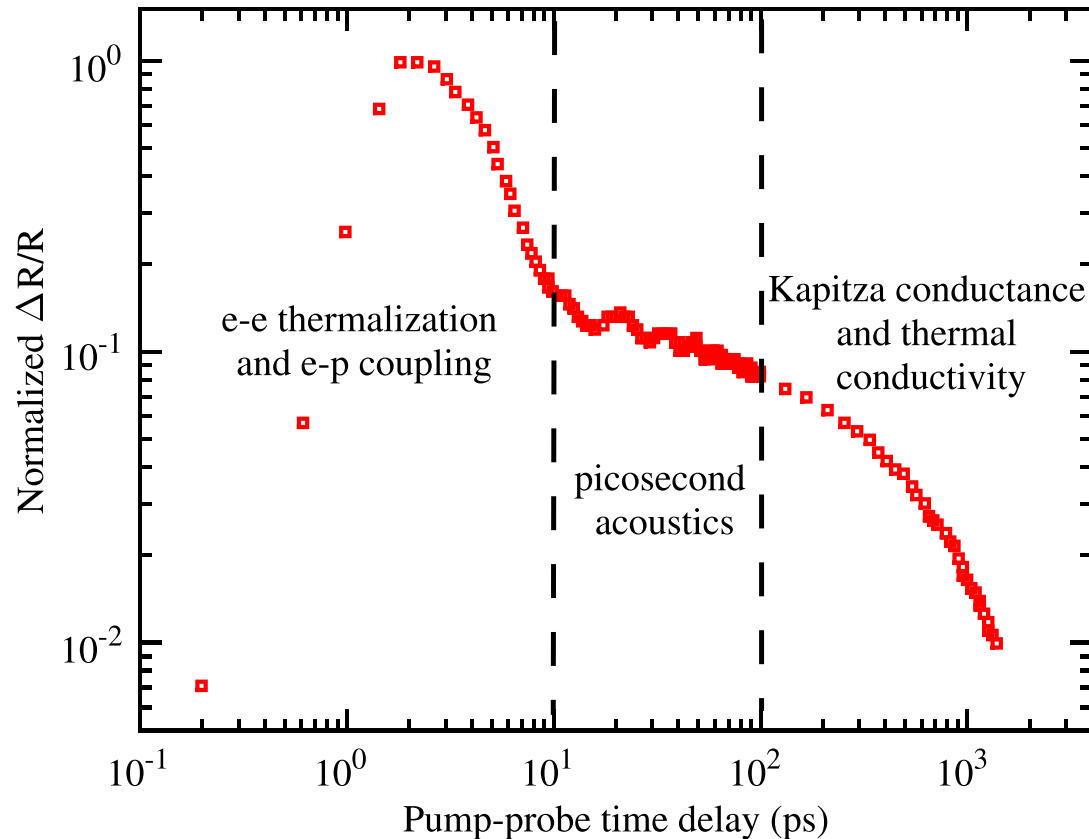


Sample



# How can we measure nanoscale heat transport processes?

Need time scale resolution < picoseconds



Pulse absorption ( $\sim 100$  fs)

↓  
Fermi relaxation and  
ballistic transport  
(few hundred fs)

↓  
Electron-phonon coupling  
(a few ps)

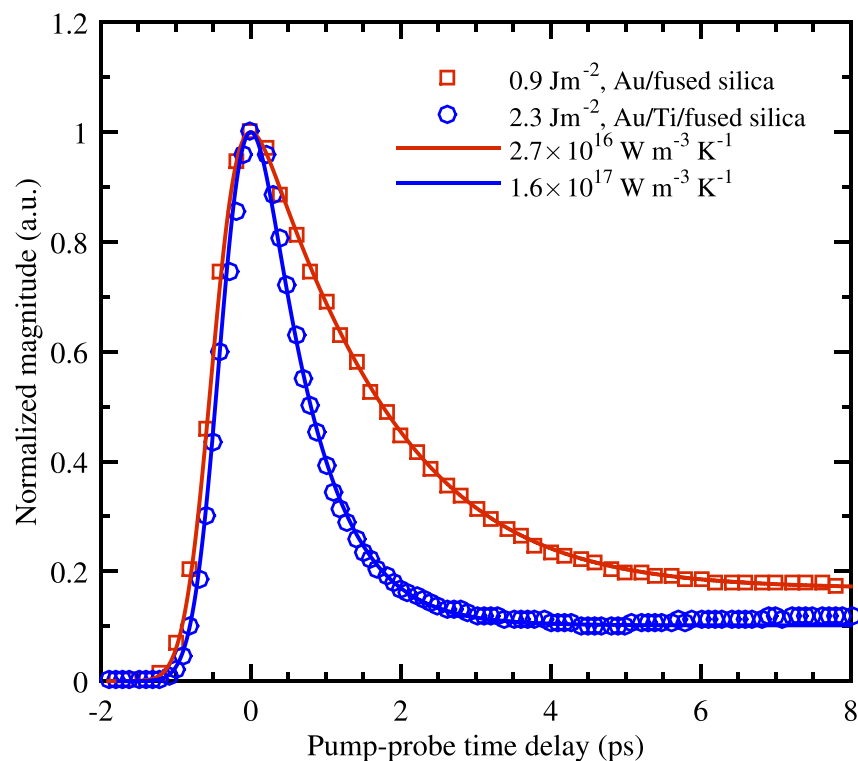
↓  
Strain propagation in film  
(10's of ps)

↓  
Thermal diffusion  
(hundreds of ps to ns)

# What can we measure with TDTR?

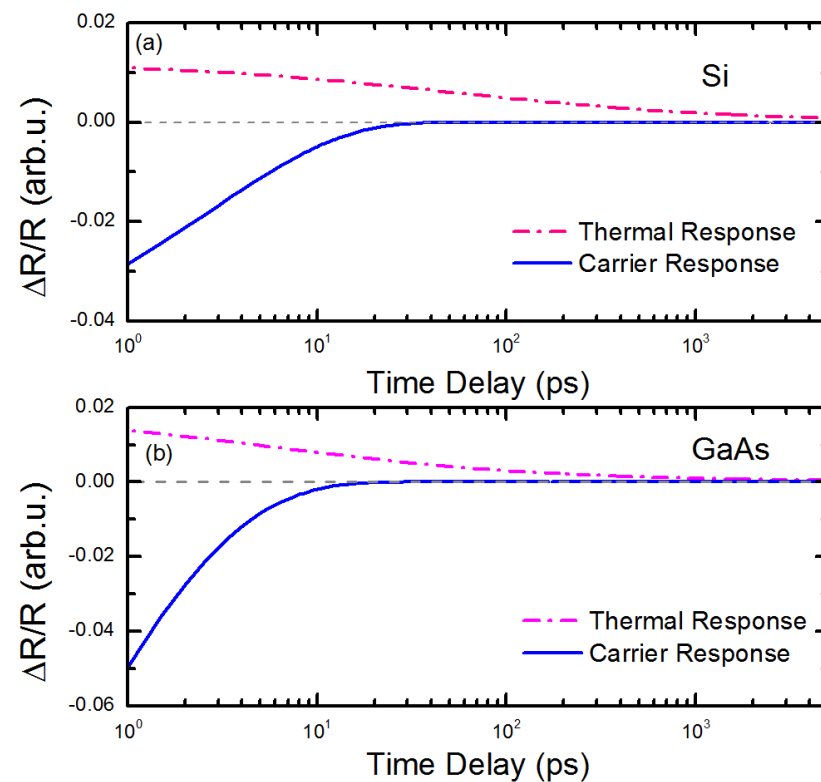
## Hot electron relaxation and recombination

### Electron-phonon coupling in metals



*JAP* **117**, 105105

### Excited state recombination in semiconductors

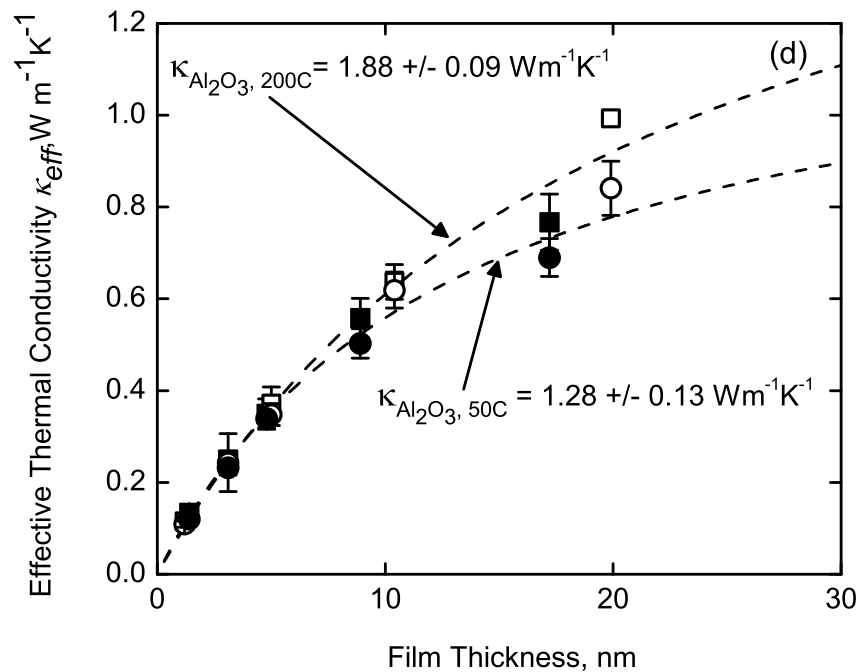


*Rev. Sci. Instrum.* **87**, 094902

# What can we measure with TDTR?

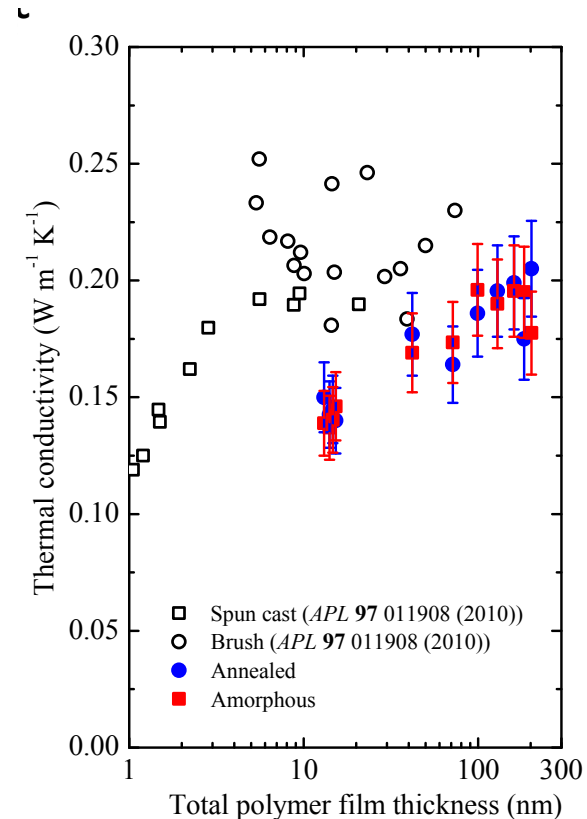
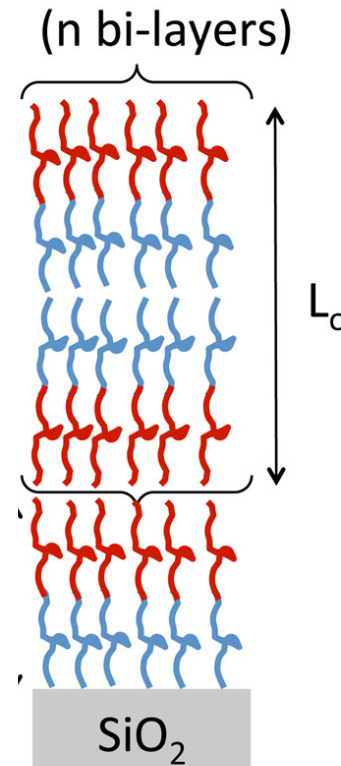
## Thermal conductivity of extremely thin films/interfaces

### ALD-grown thin films



*Thin Solid Films* **650**, 71

### Block co-polymer thin films

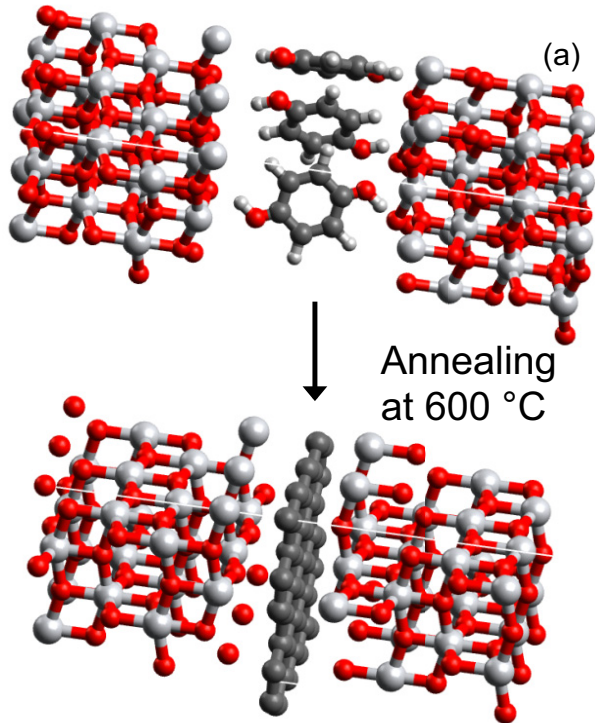


*J. Heat Trans.* **138** 024505

# What can we measure with TDTR?

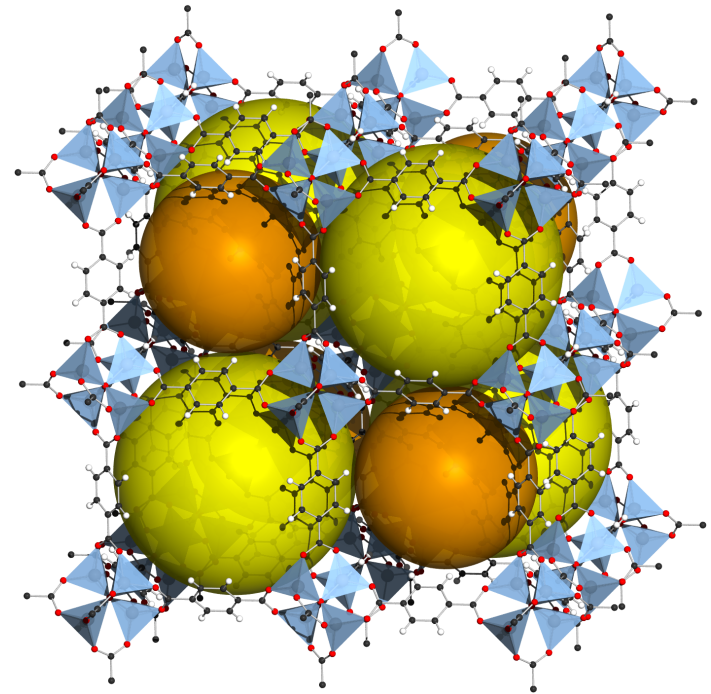
## Heat capacity of thin films and some bulk systems

### Organic/inorganic hybrid “superlattices”



*Phys. Rev. B* **93**, 024201

### Metal-organic frameworks “MOFs”

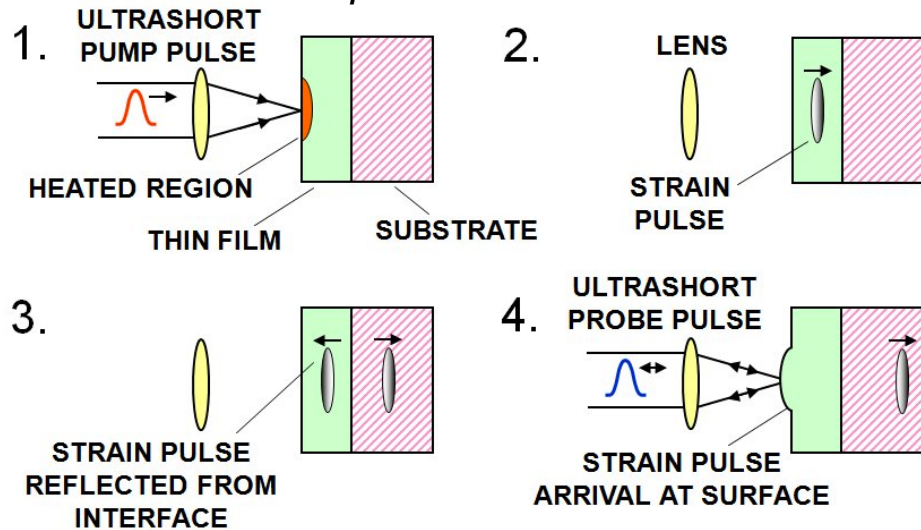


*Adv. Mater.* **27**, 3453  
*MRS Bulletin* **41**, 877

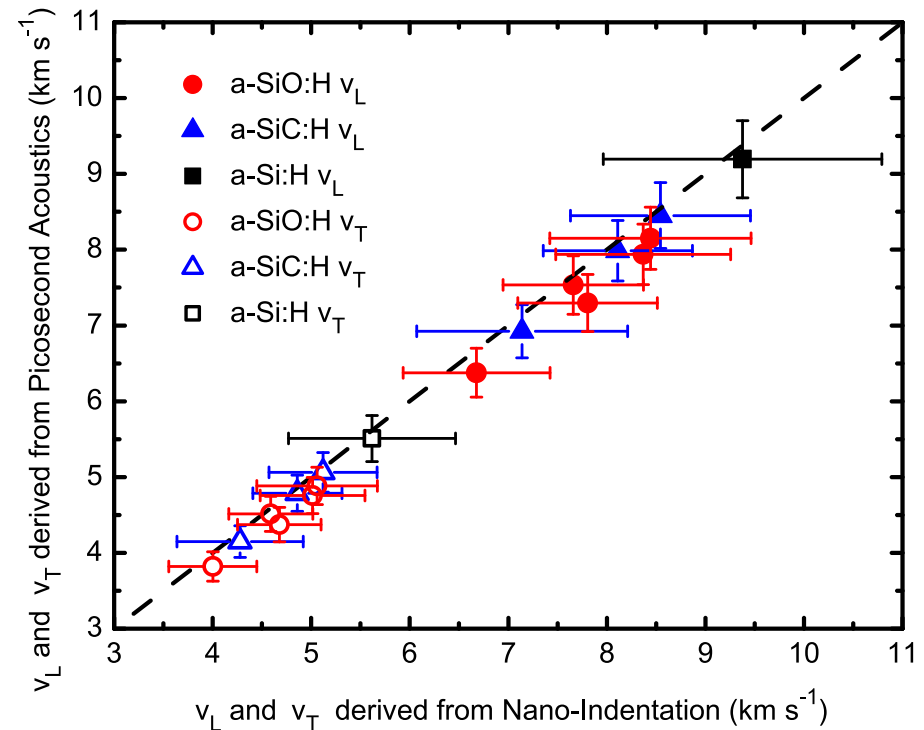
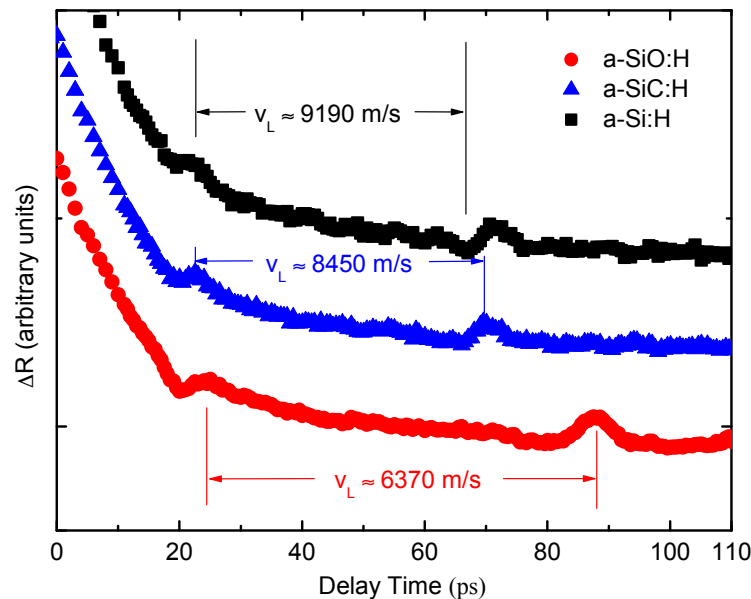


# What can we measure with TDTR?

Wikipedia



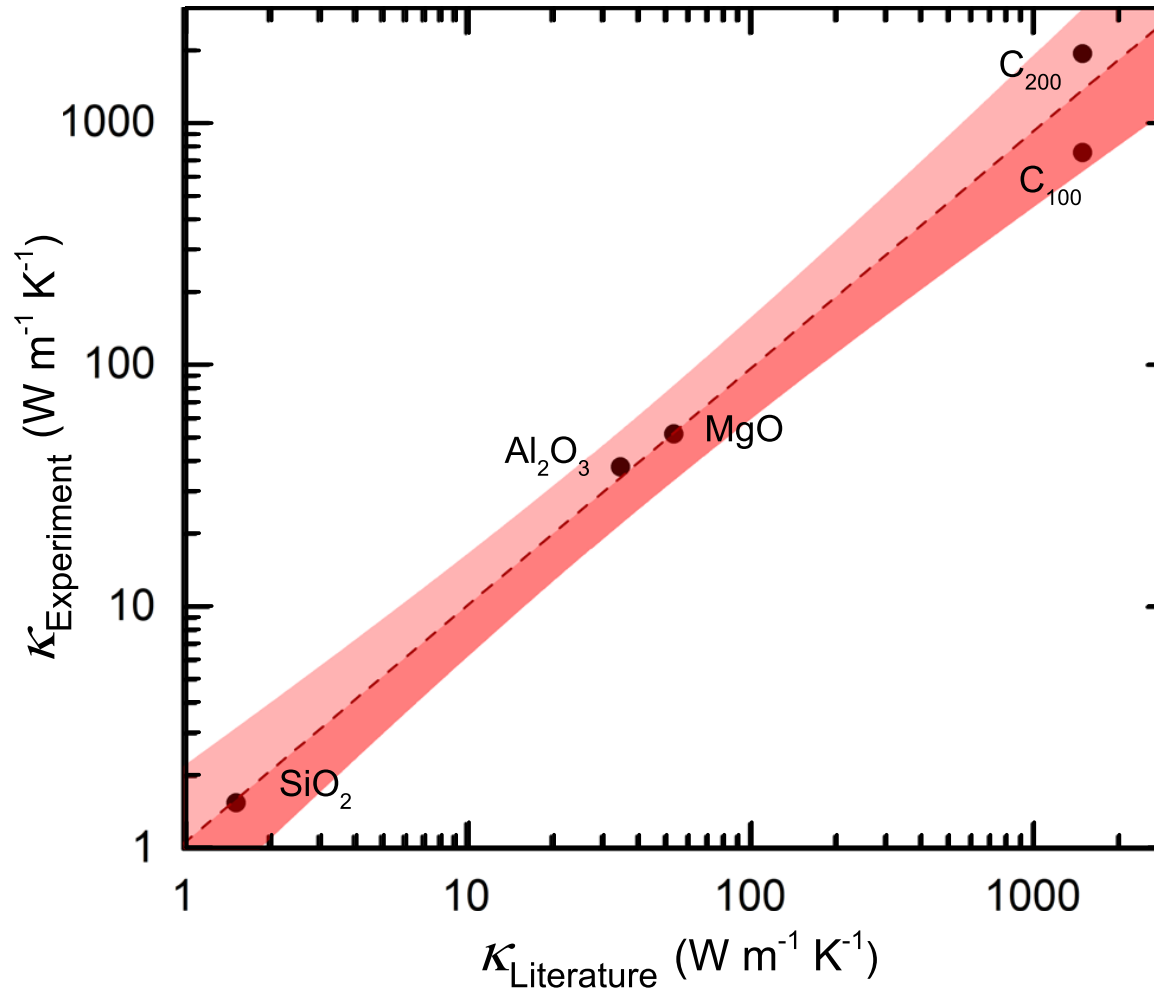
**Measurement of strain wave/acoustic wave propagation in thin films**



**APL 109, 191905**

# What can we measure with TDTR?

## Bulk materials



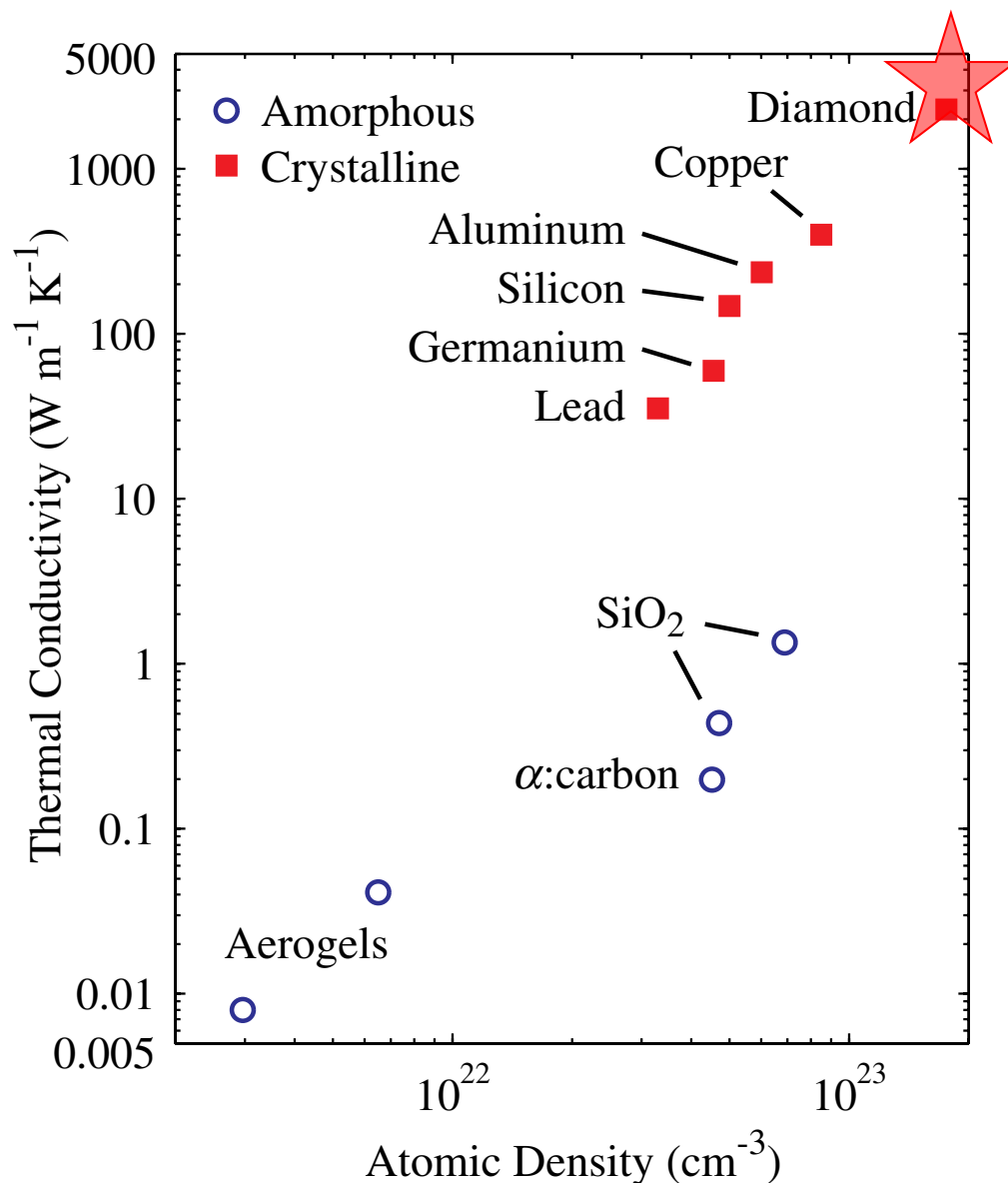
*Rev. Sci. Instrum.* **87**, 094902

*Appl. Phys. Lett.* **111**, 151902

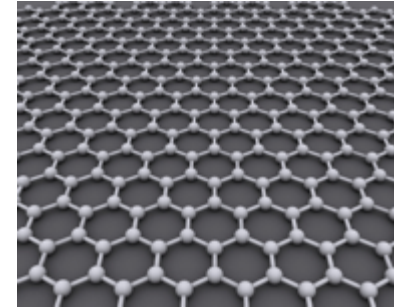
# Outline

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- **Weakly bonded solids: new lower limits to thermal conductivity**
- Functionalized interfaces at graphene contacts: tuning heat and electrical transport via the interfacial bond
- Heat transport across single molecule interfaces: when does a molecule become a defect?
- Molecular interfaces in organic/inorganic composites: diffusive scattering via the vibron-phonon interaction

# Thermal conductivity of materials – Role of the bond



Ultrahigh thermal conductivities (diamond, graphene, BAs, etc)



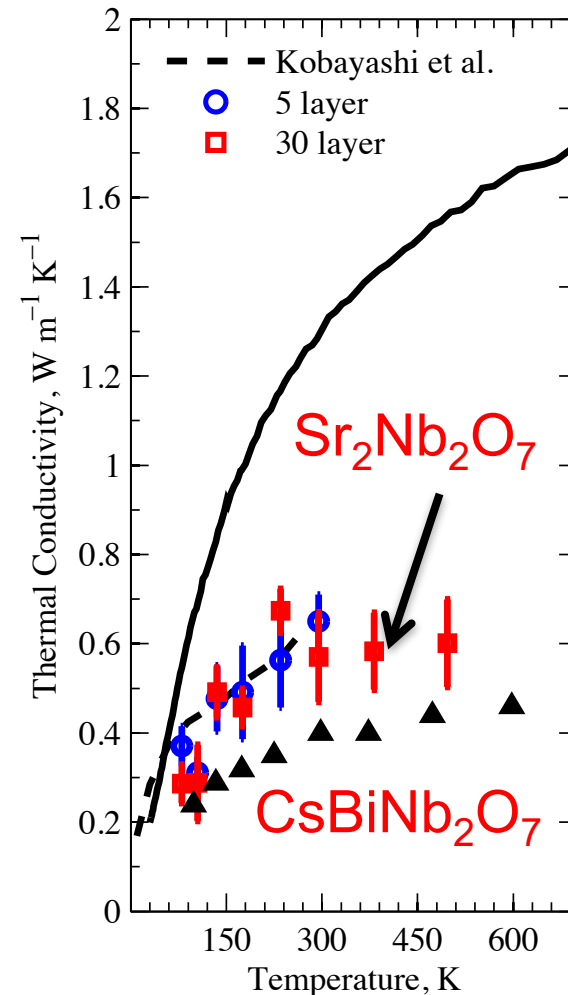
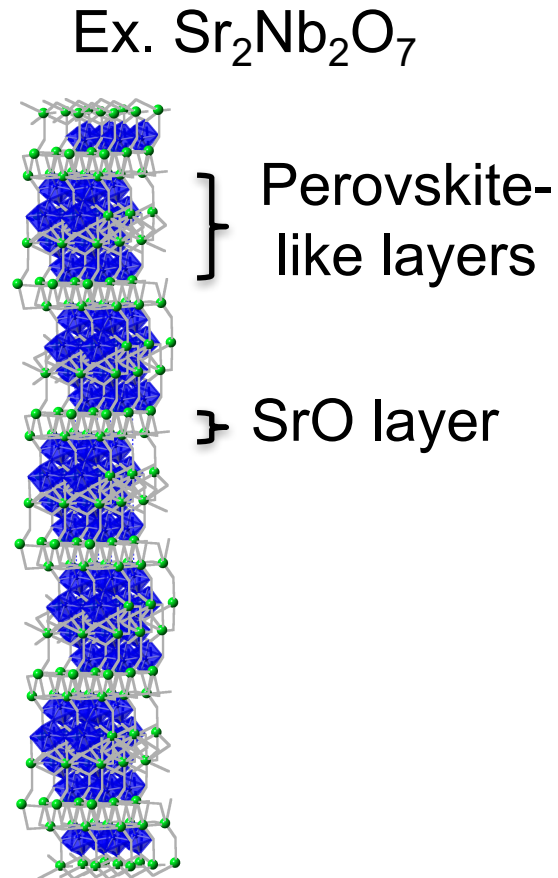
Light masses, stiff bonds

**How do we go the other way?**  
**Heavy masses, weak bonds**

*PRL* **110**, 015902 (2013)

# Thermal conductivity of materials – Role of the bond

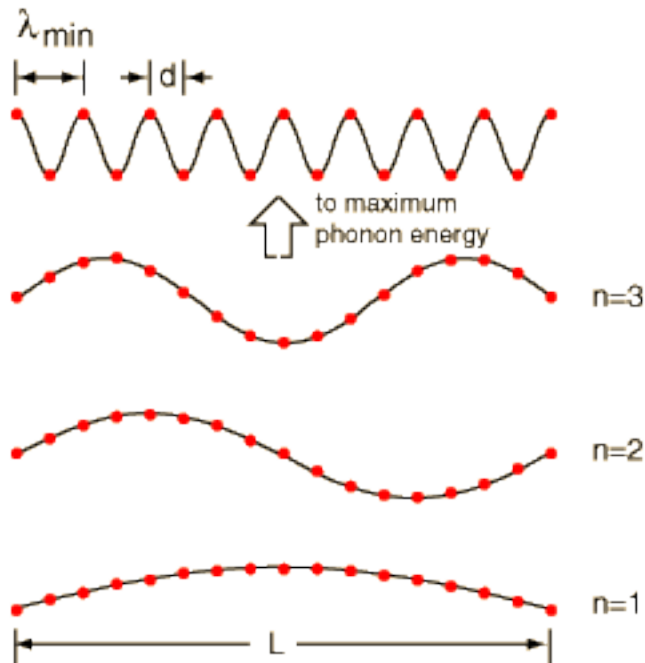
Layered structures can exhibit ultralow thermal conductivity





# The Einstein oscillator

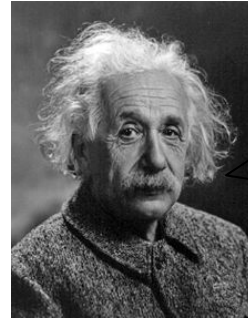
The phonon picture  
(coupled oscillators):  
several different wavelength in a  
lattice (many energies)



$$\lambda_{\min} = 2d$$

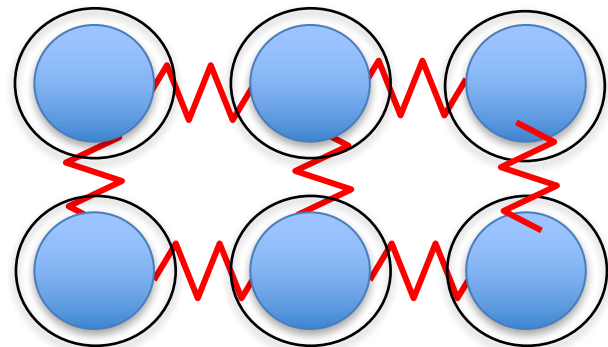
$$\lambda_n = \frac{2L}{n}$$

$$\lambda_{\max} = 2L$$



Vibrations of atoms  
are independent  
with random phases

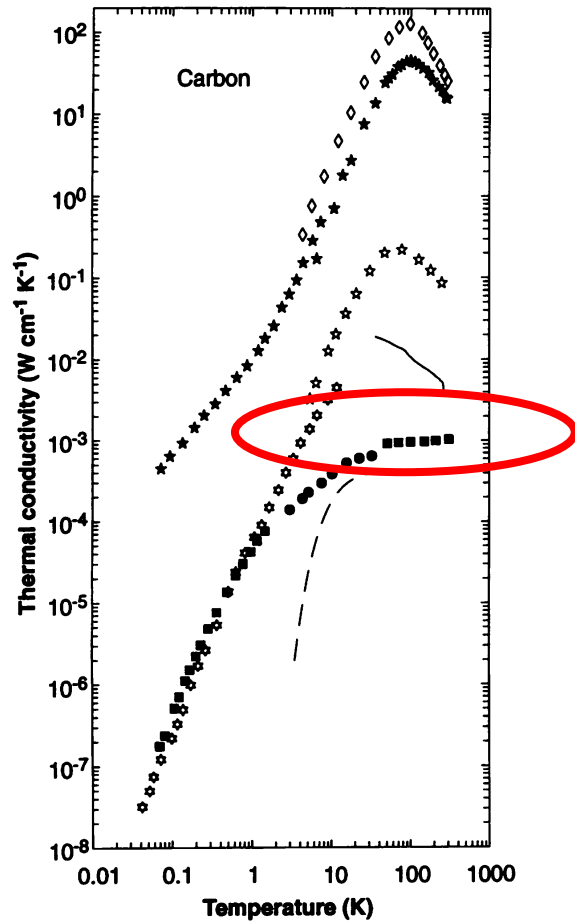
Single frequency of vibration of  
atom and energy “hops” from  
one site to another



**“Springs” are very very weak**

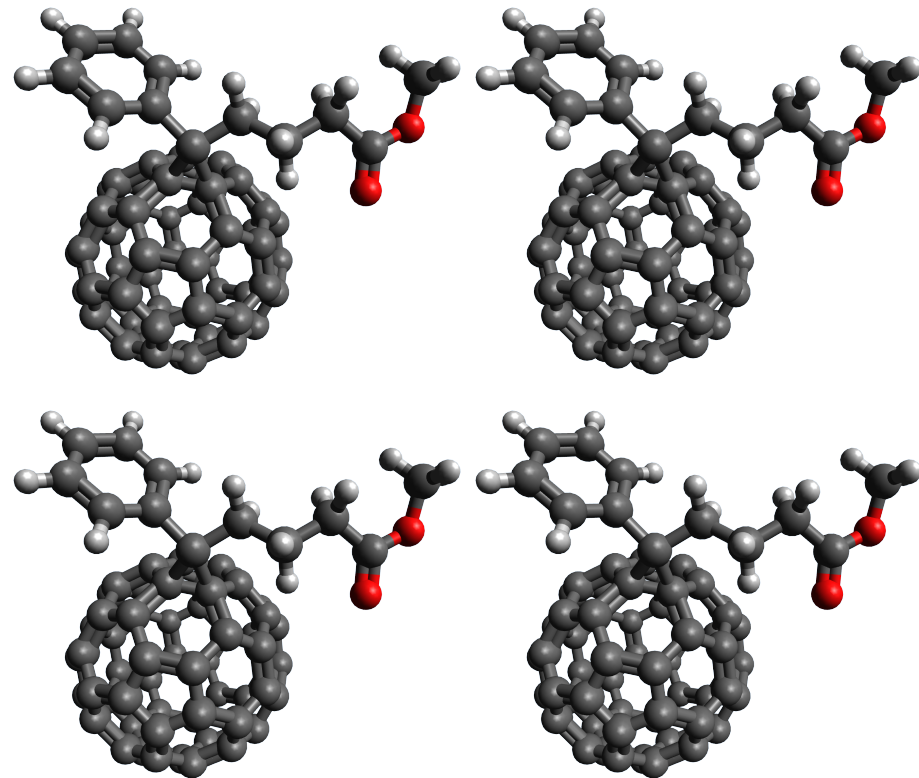
# The Einstein oscillator – weakly interacting buckyballs

## Fullerene films: Low thermal conductivities



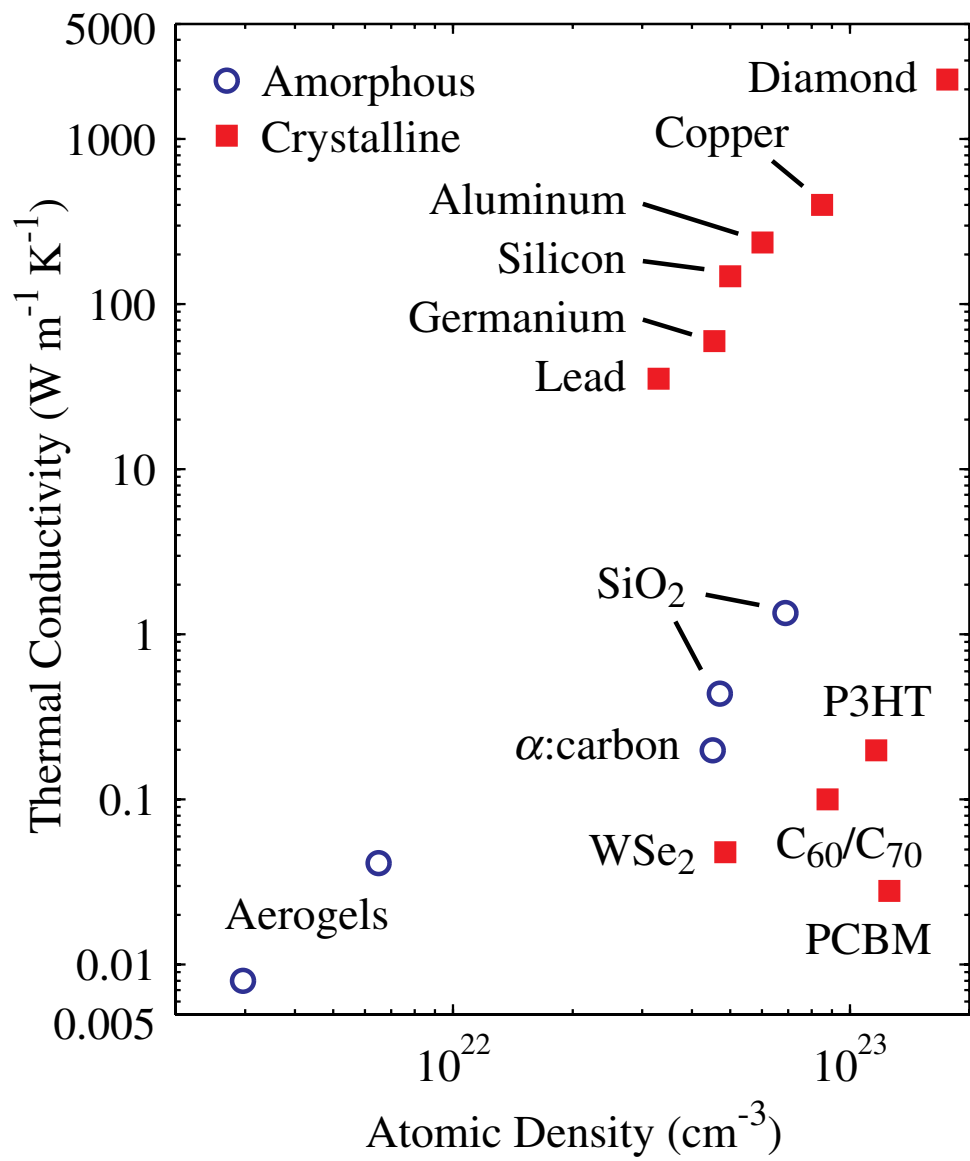
*Science* **259**, 1145

Add “disorder” with molecules:  
[6,6]-phenyl  $\text{C}_{61}$ -butyric acid  
methyl ester (PCBM)

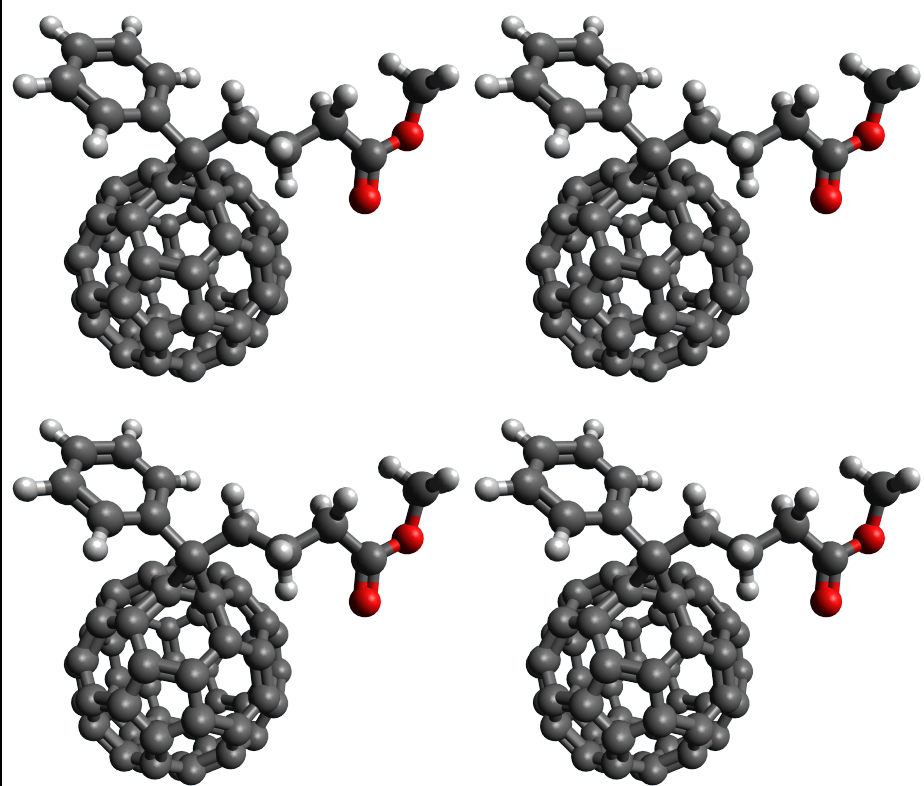


# New lower extreme of thermal conductivity of materials

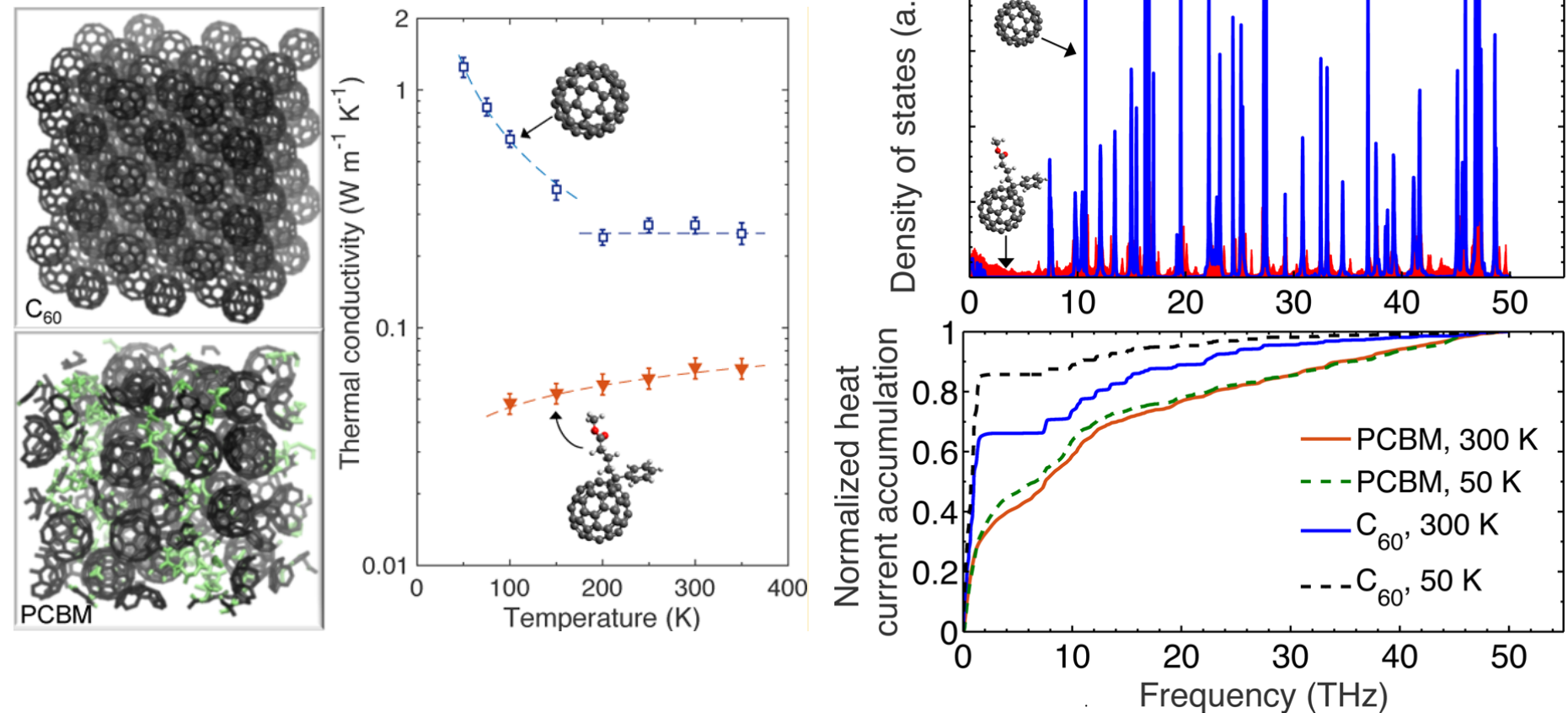
But why???



$$\kappa = \frac{1}{3} C v \lambda = \frac{1}{3} C v_g^2 \tau$$



# Turn to molecular dynamics simulations

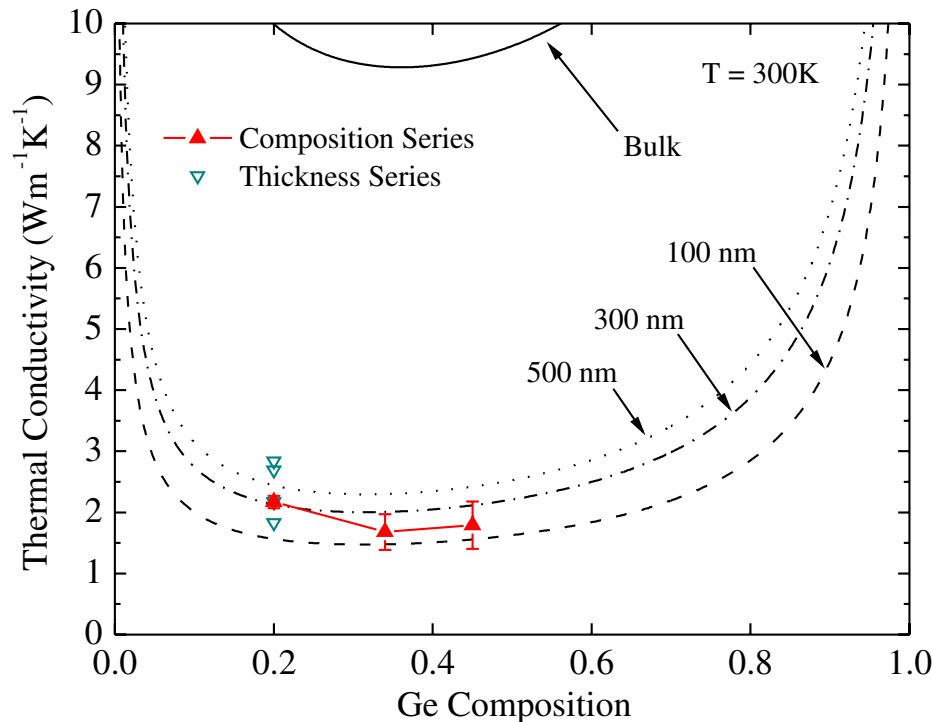


*J. Phys. Chem. Lett.* **8**, 2153 (2017)

*Phys. Rev. B* **96**, 220303 (2017)

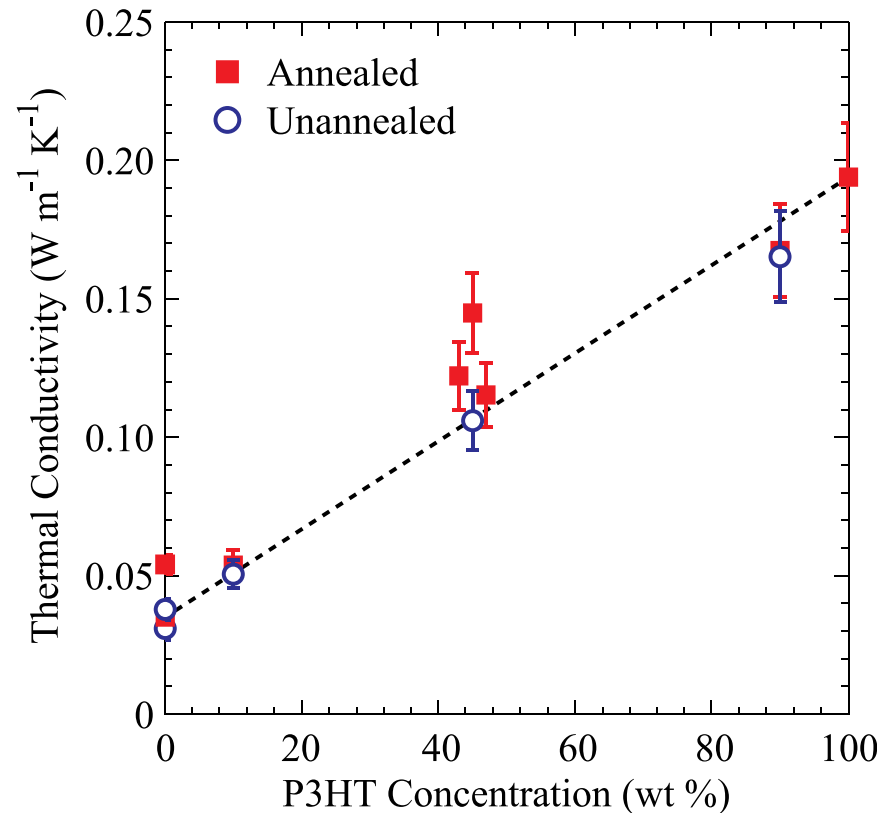
# Tuning the thermal conductivity of molecular films

**Covalent bonds**  
**Crystals/“phonons”**  
 **$\text{Si}_x\text{Ge}_{1-x}$  alloy**



*Phys. Rev. Lett.* **109**, 195901

**Weak bonds**  
**Vibrons**  
**PCBM-P3HT blend**



*Appl. Phys. Lett.* **102**, 251912

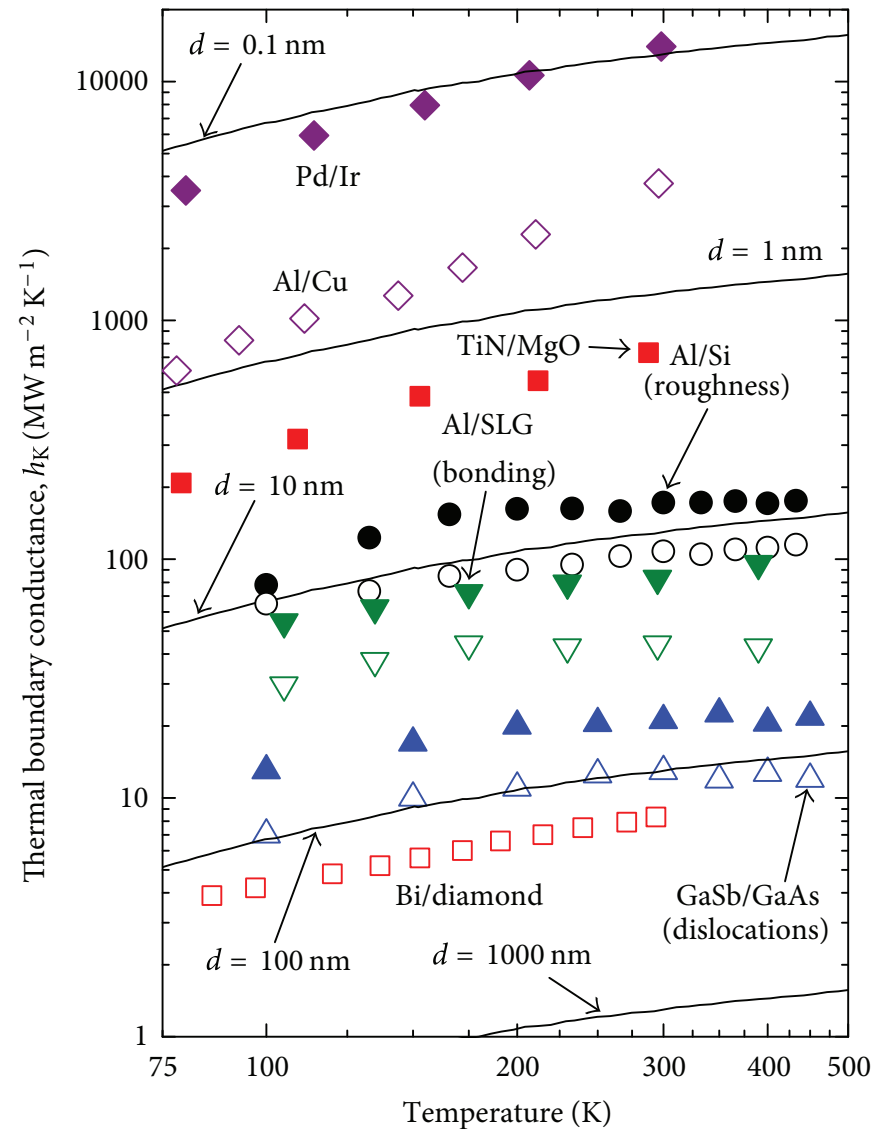
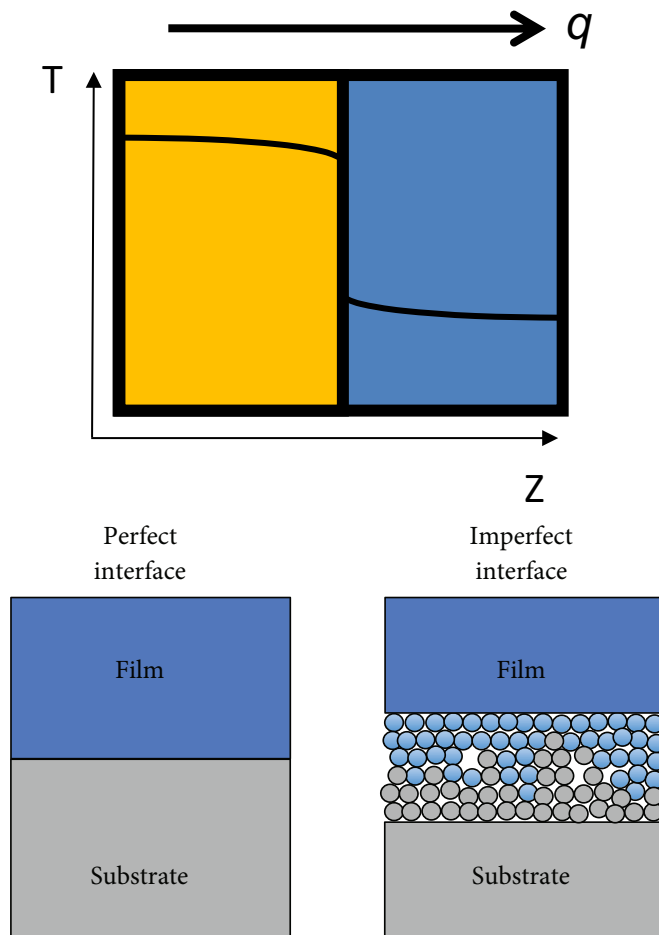


# Outline

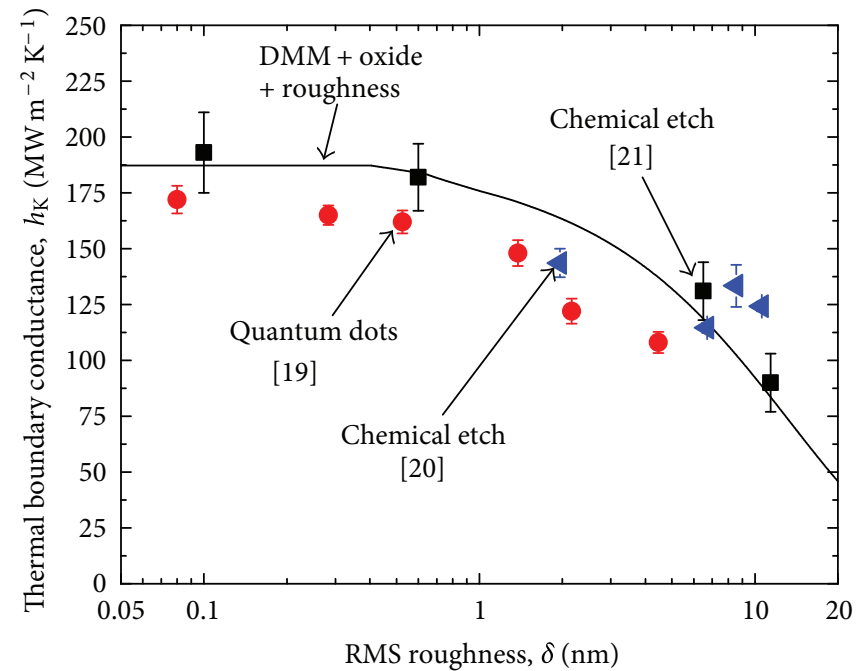
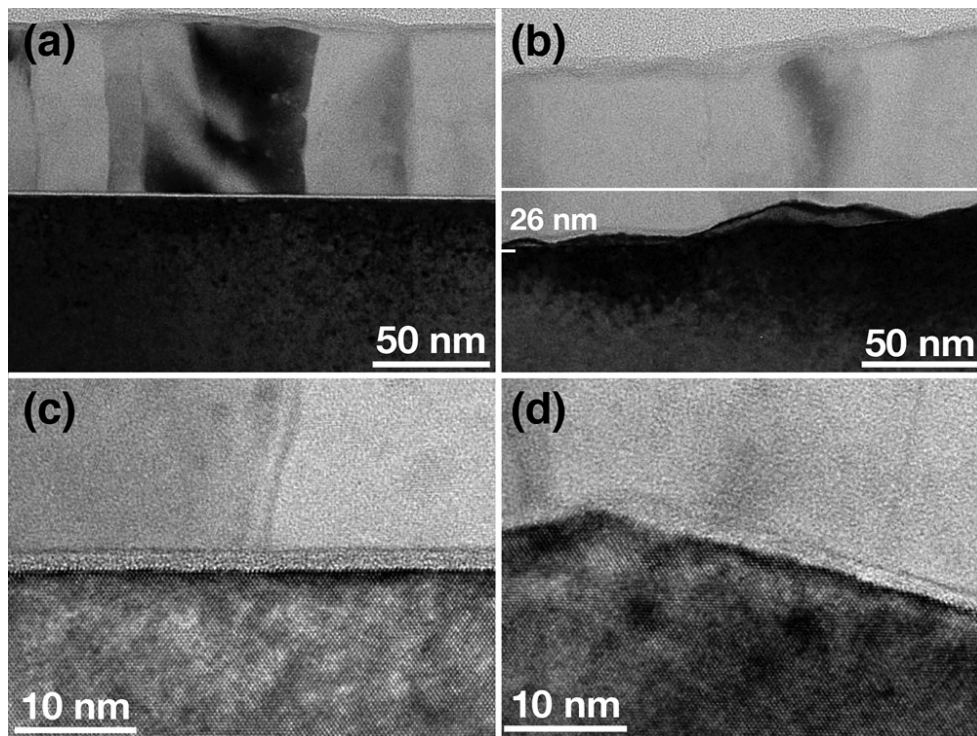
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# Thermal boundary conductance – nanoscale resistances

$$q = h_K \Delta T = \frac{1}{R_K} \Delta T$$

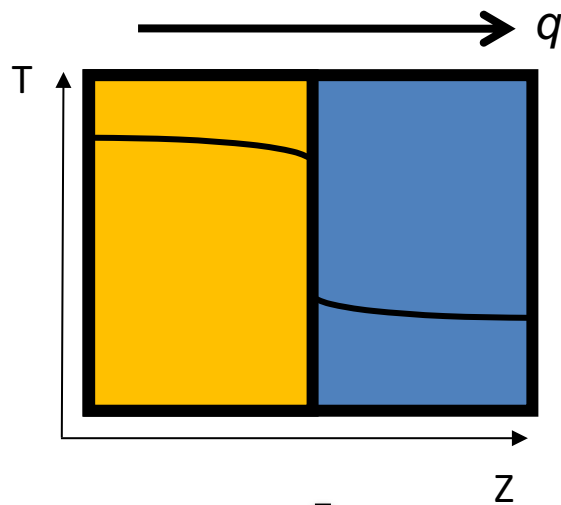


# Thermal boundary conductance – nanoscale resistances



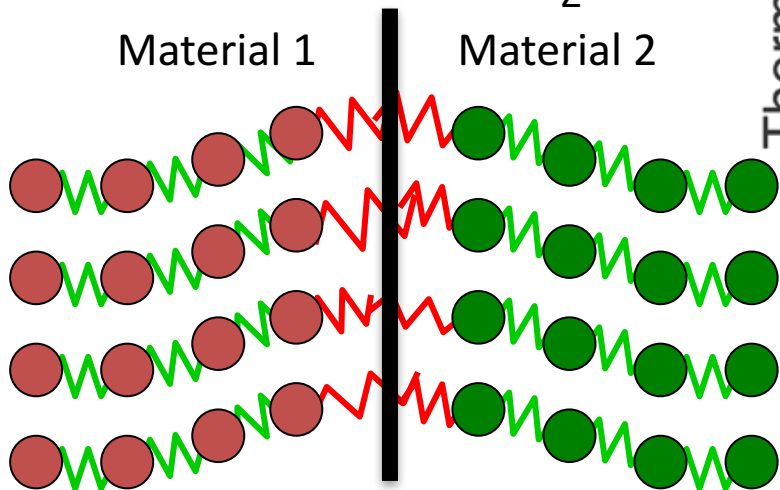
# Increase in bonding increases solid/solid TBC

$$q = h_K \Delta T = \frac{1}{R_K} \Delta T$$

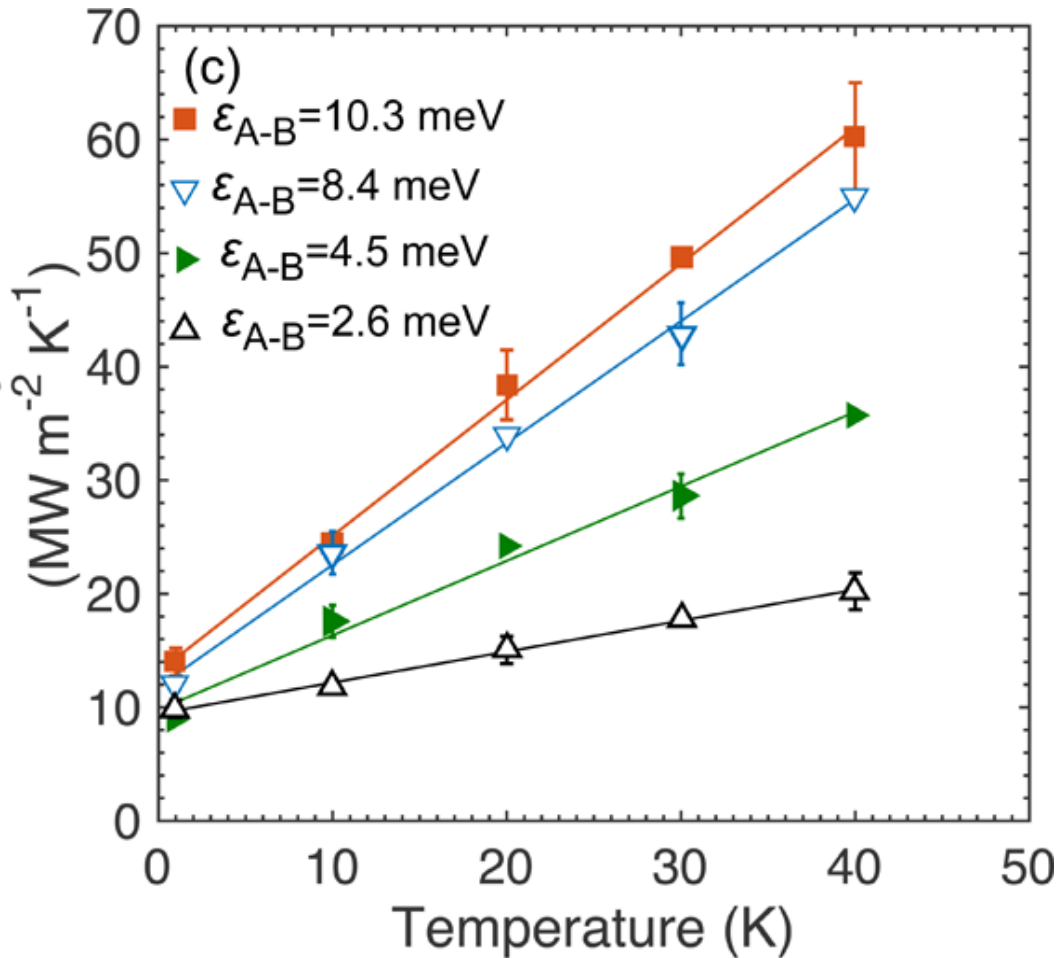


Material 1

Material 2

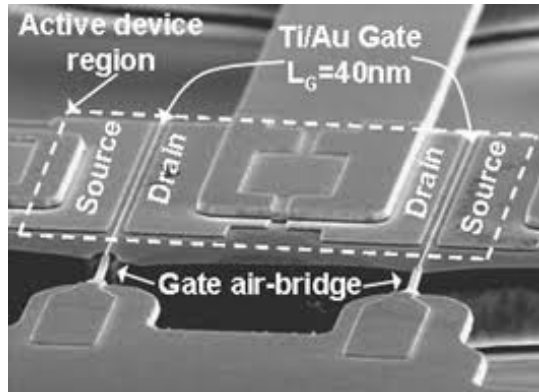


Thermal boundary conductance

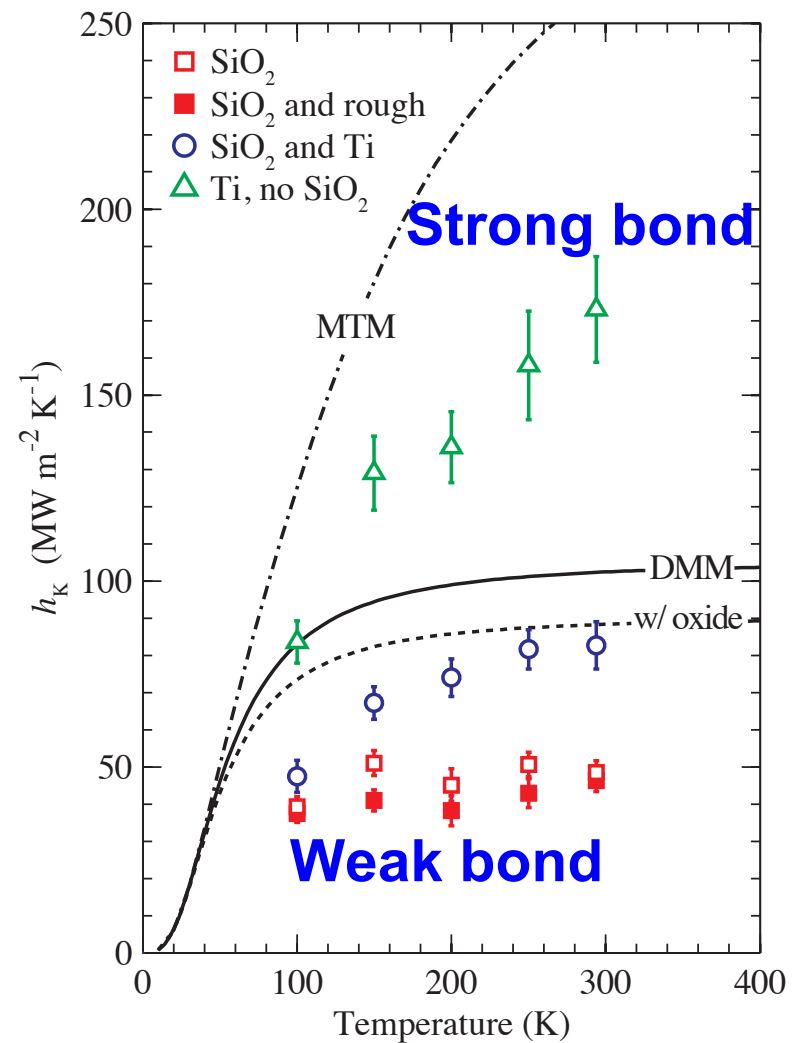
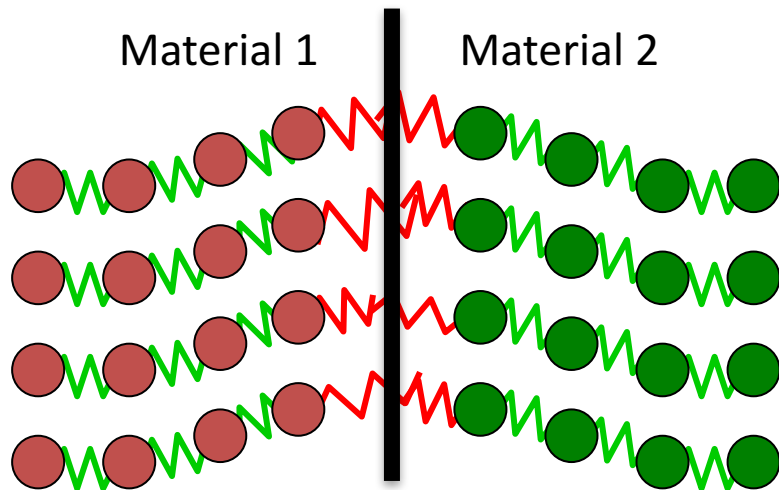


*J. Phys. Chem. C* **120**, 24847

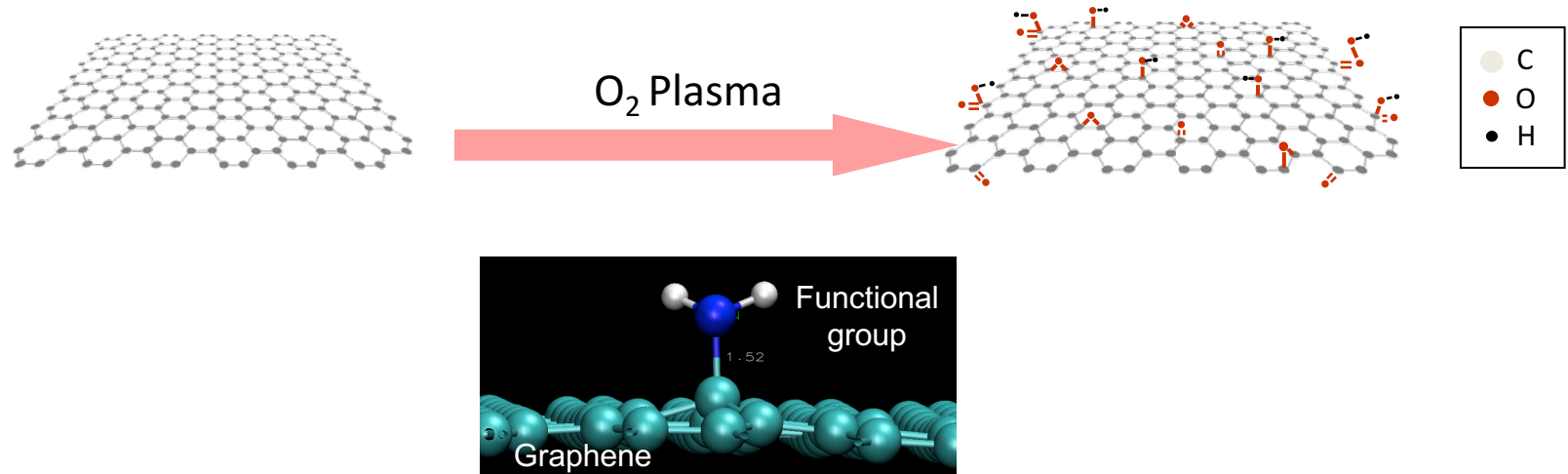
# Thermal boundary conductance – bonding effects



However, bonding plays a **HUGE ROLE** at the Au/Si interface



# Atmospheric plasma functionalization of graphene surfaces



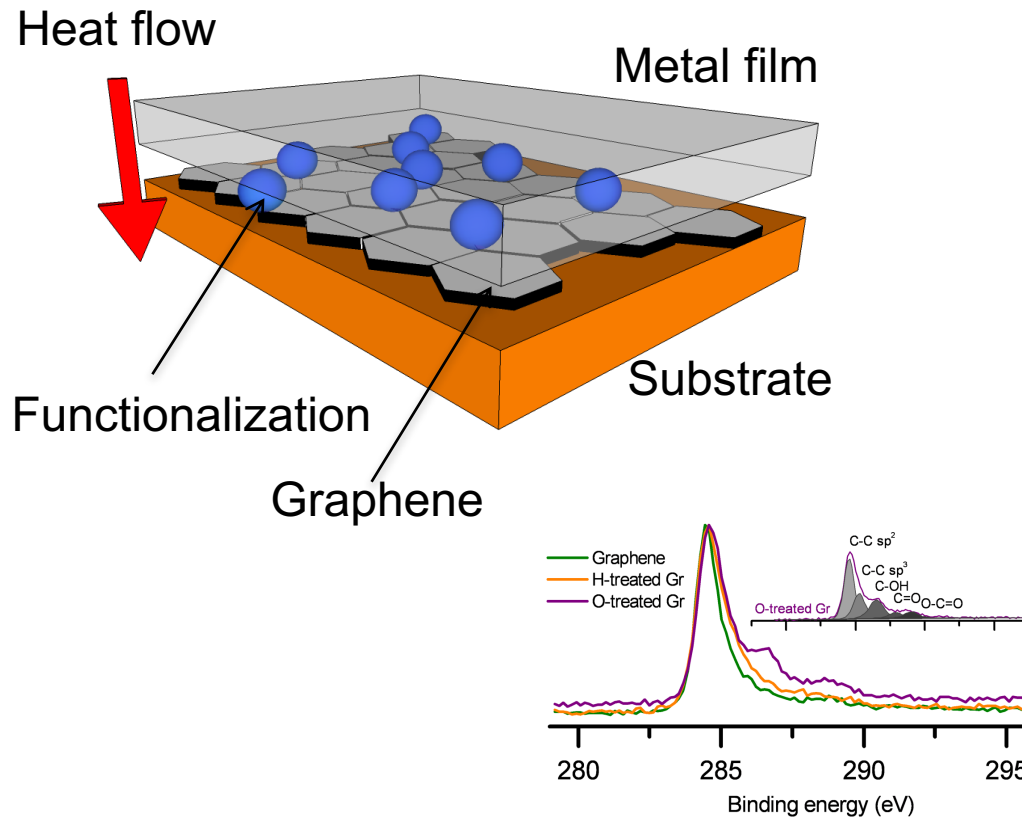
**Functional groups covalently bound to graphene**  
**Reversible after anneal**

*Appl. Phys. Lett.* **96**, 231501

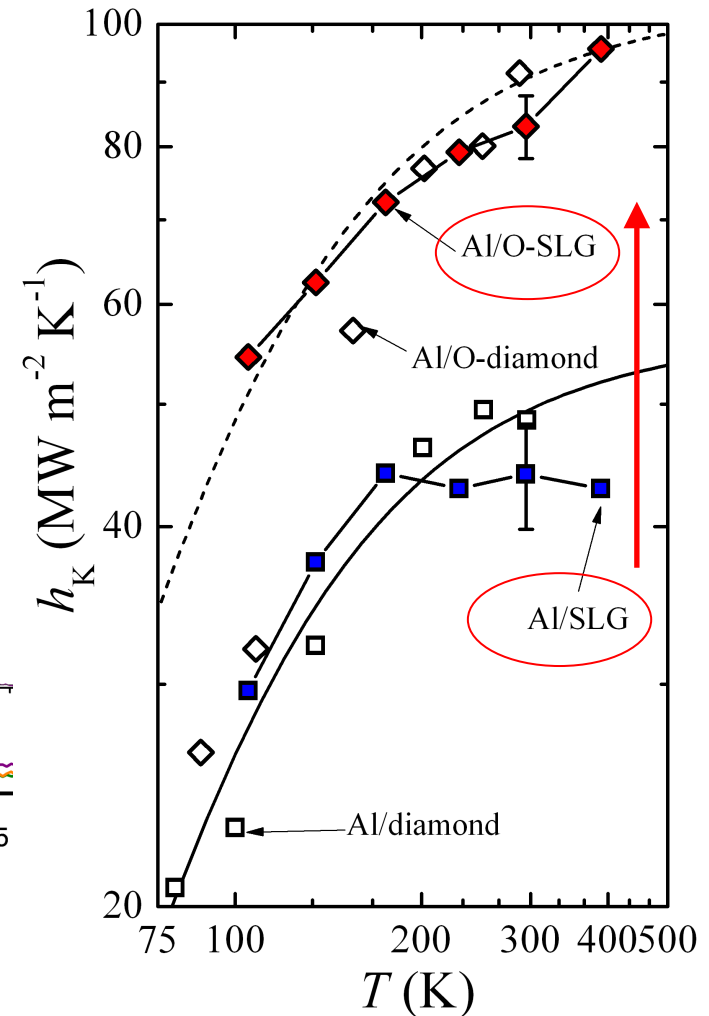
**Collaboration:** Scott Walton (NRL)



# Atmospheric plasma functionalization of graphene surfaces



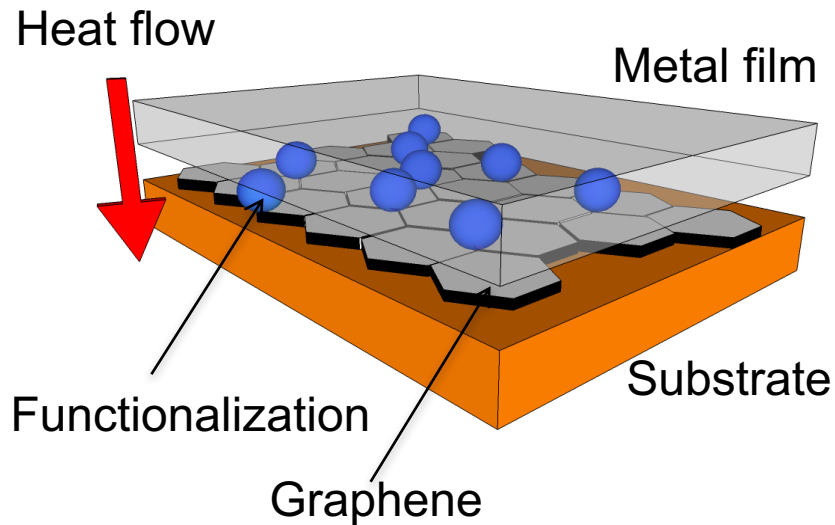
- **Al/graphene interaction increased with oxygen bond (Al-O bond)**
- **But what implications does this have for SLG devices?**



*Nano Lett.* **12**, 590 (2012)

**Collaboration: Scott Walton (NRL)**

# Au/graphene electronic contacts



**C-F bond inert  
does not want to  
interact!**

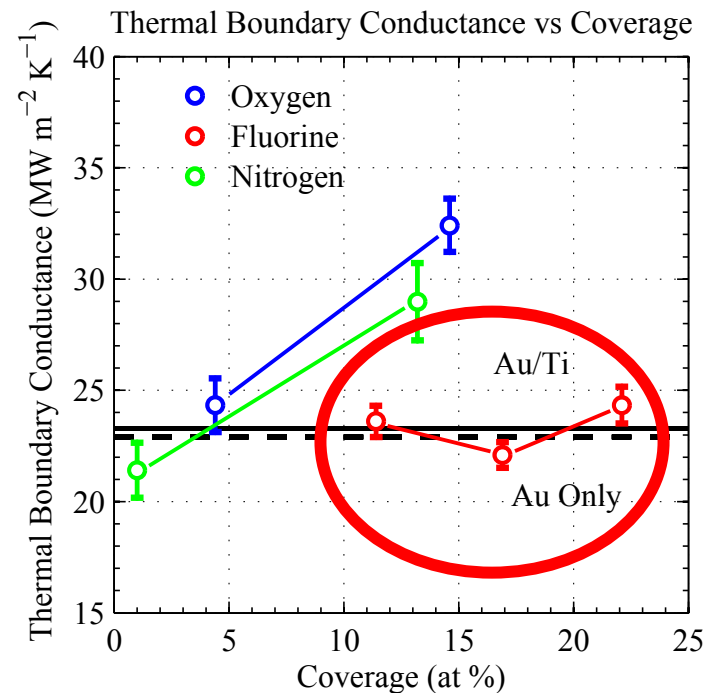
*Nano Lett.* **15**, 4876 (2015)

## Thermal boundary conductance results

Similar trends for oxygen as Al

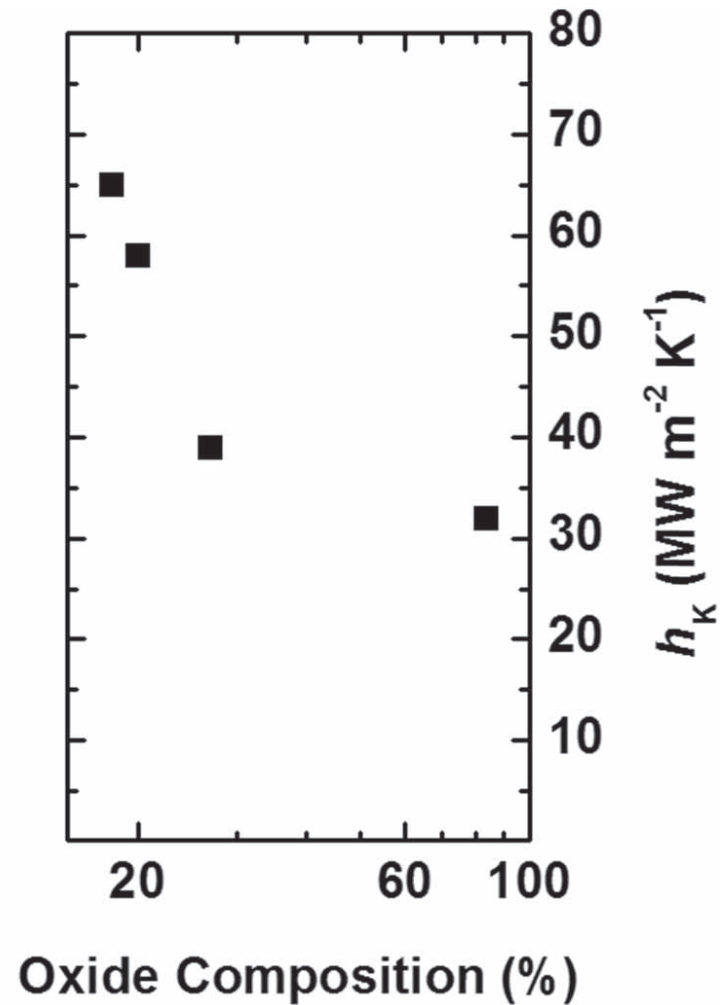
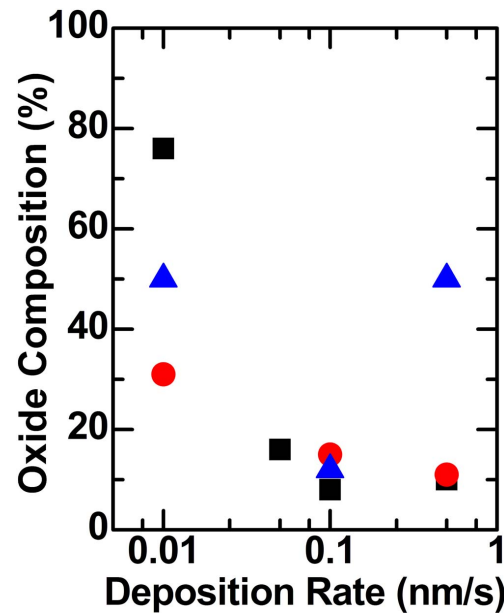
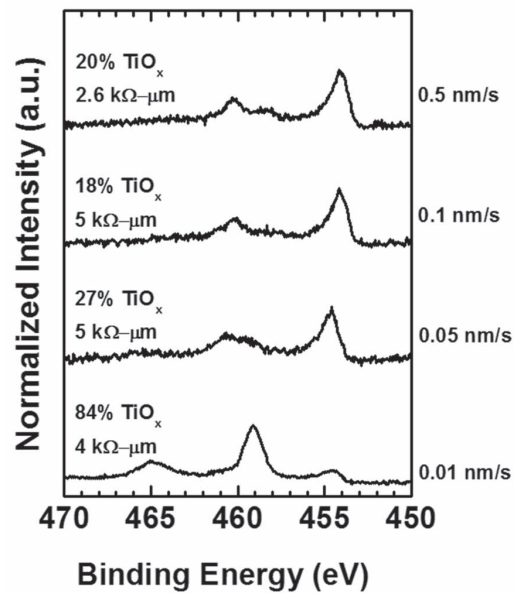
**Ti adhesion layer does nothing**

**Fluorine does nothing**



**Collaboration: Scott Walton (NRL)**

# Chemistry effects on the TBC across Au/Ti/graphene

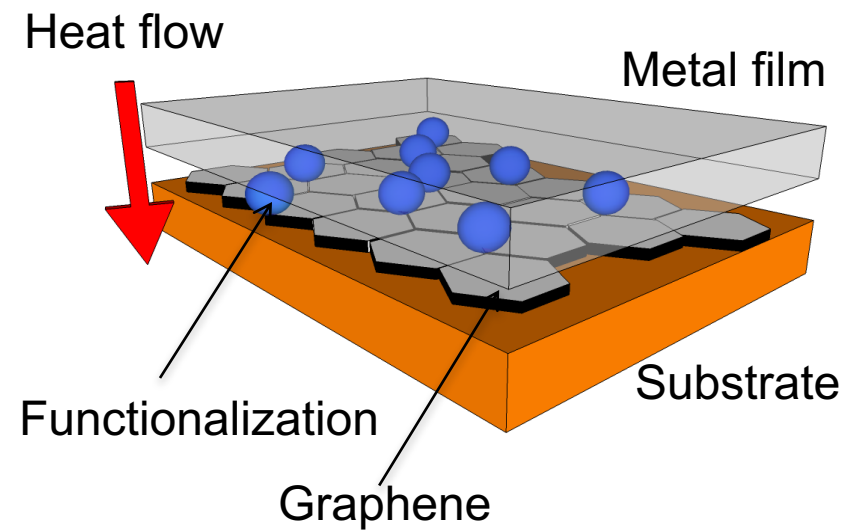


*Nanotechnology* **29**, 145201

**Collaboration:** Stephen McDonnell (UVA)

# Chemistry effects on the TBC across Au/Ti/graphene

## Bonding engineering at interface



## Defect engineering at interface

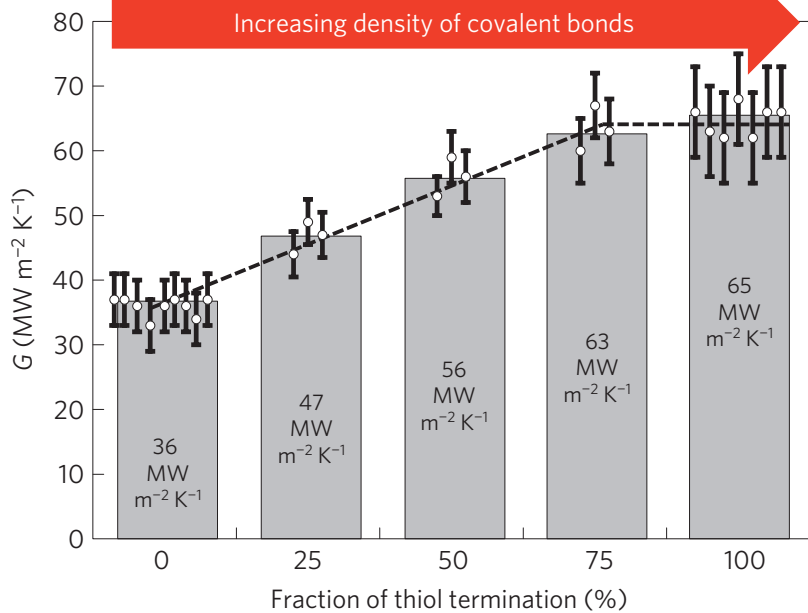
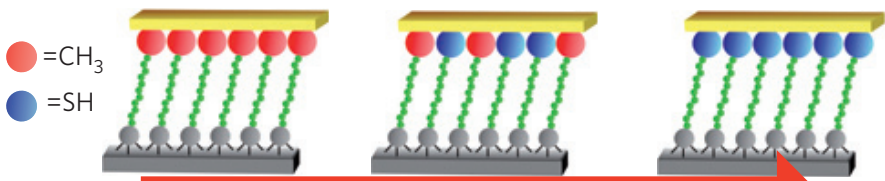


**When does a molecular bond become a defect?**

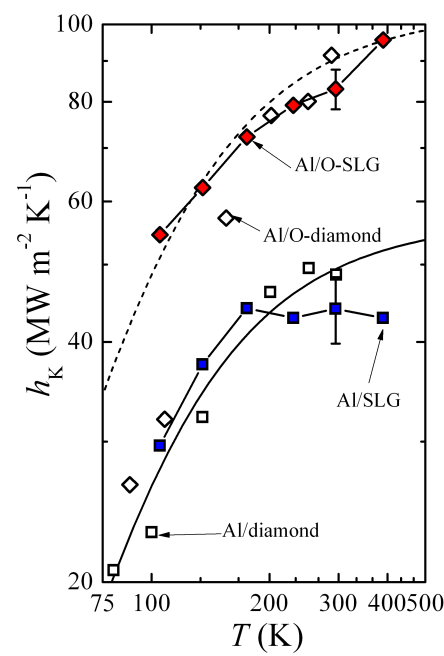
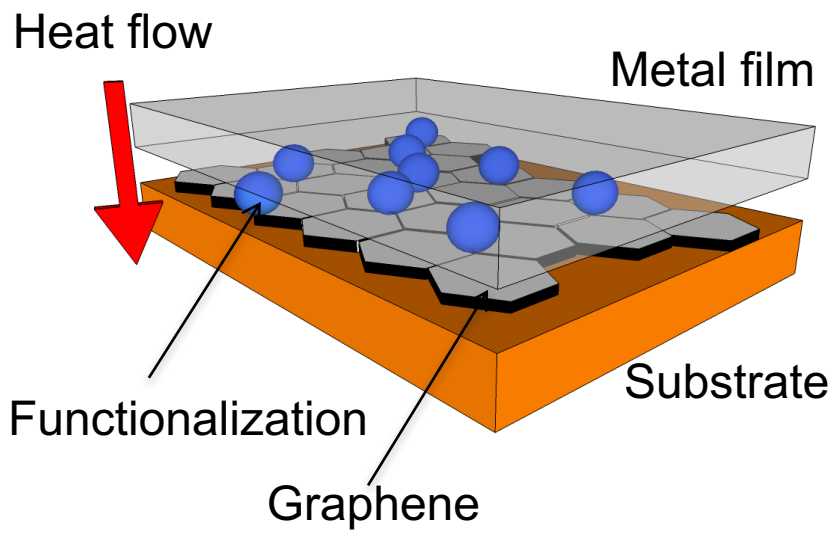
# Outline

- TDTR: Measurement of thermal conductivity of thin films and thermal resistance across interfaces
- Weakly bonded solids: new lower limits to thermal conductivity
- Functionalized interfaces at graphene contacts: tuning heat and electrical transport via the interfacial bond
- **Heat transport across single molecule interfaces: when does a molecule become a defect?**
- Molecular interfaces in organic/inorganic composites: diffusive scattering via the vibron-phonon interaction

# Contact chemistry to manipulate TBC

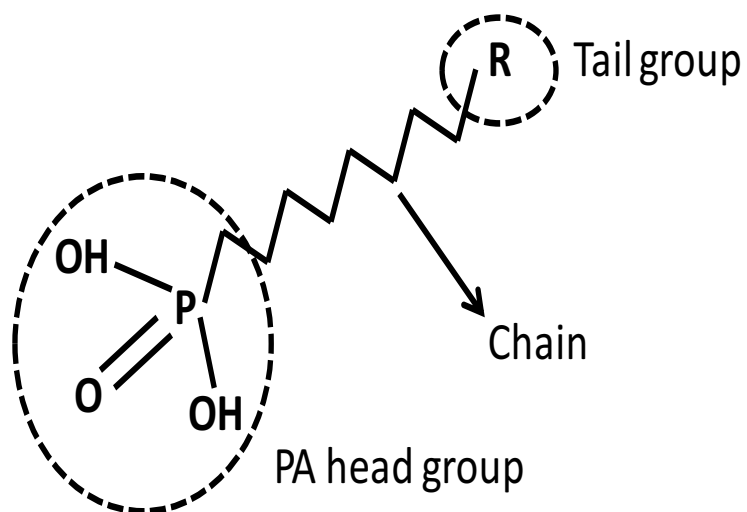


Nat. Mat. 11, 502 (2012)



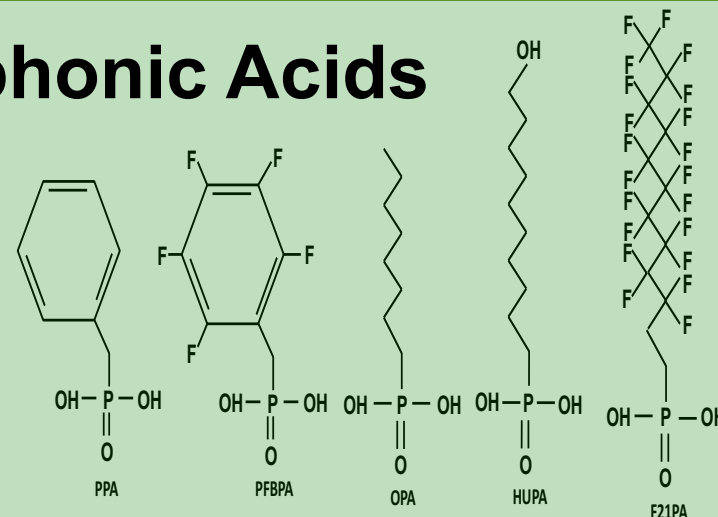
Nano Lett. 12  
590 (2012)

# Phosphonic acid interfaces: Size and mass control



**Metals: Al, Au, Ni**

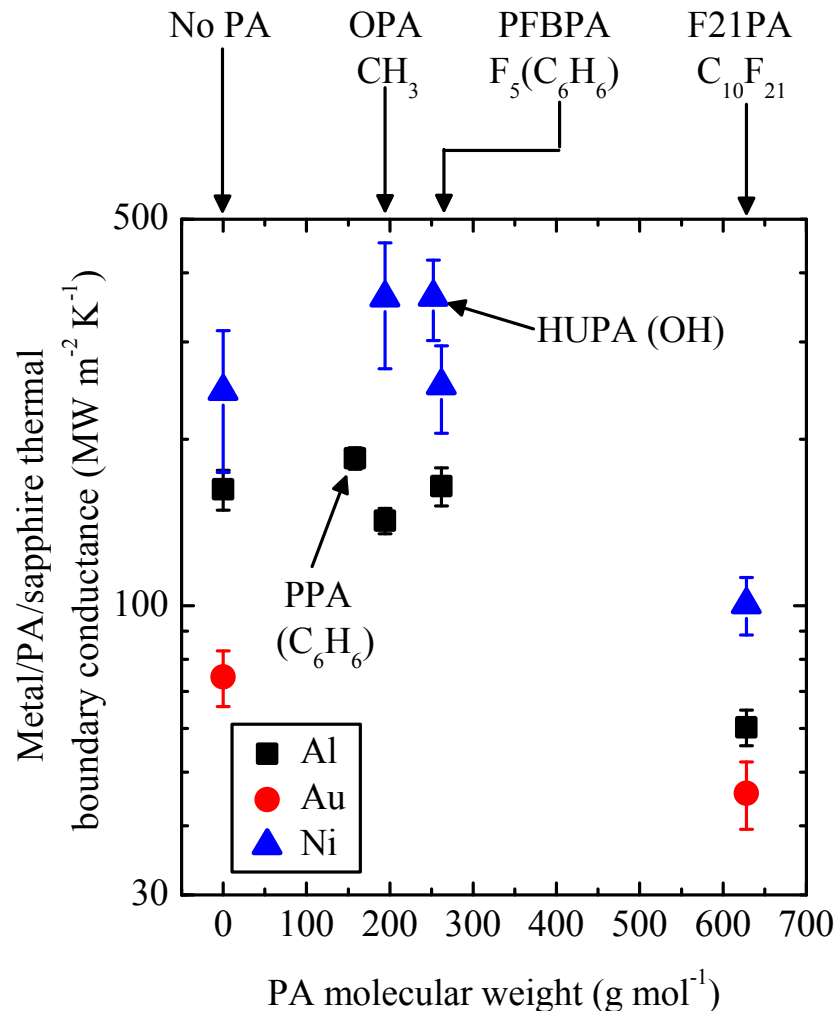
## Phosphonic Acids



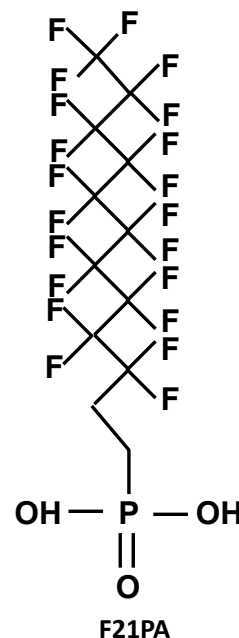
**Sapphire substrate**



# Phosphonic acid interfaces: Size and mass control



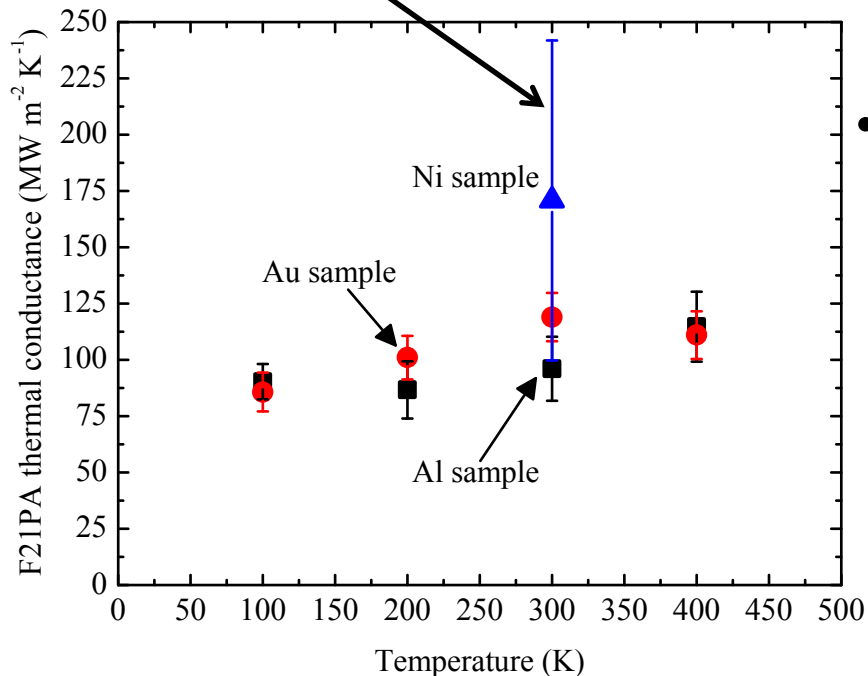
- No correlation with length
- Most noticeable change in TBC with large MW F21PA
- Probably not length effect (HUPA ~ 1.7 nm, F21PA ~ 2 nm)



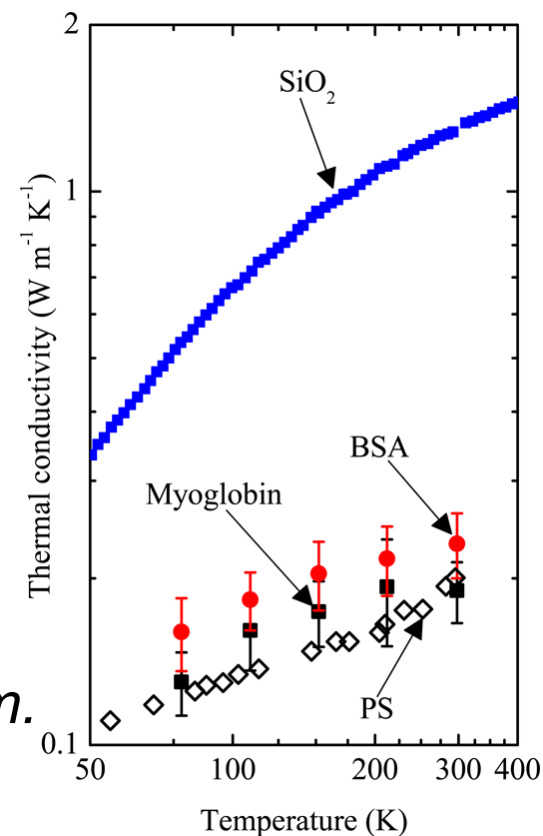
**Is the change in TBC the metal/PA interface or the F21PA itself?**

# Heat transport in single F21PA molecules

Uncertainty in Ni thickness,  
but in process of post-  
TDTR analysis to reduce  
uncertainty



- No prior evidence of diffusive heat flow in a molecule (ballistic observed by Wang *et al. Science* **317**, 787 (2007))
- ~2 nm molecule, so effective thermal conductivity of F21PA ~0.2 W/m/K at RT



*J. Phys. Chem. C* **119**, 20931

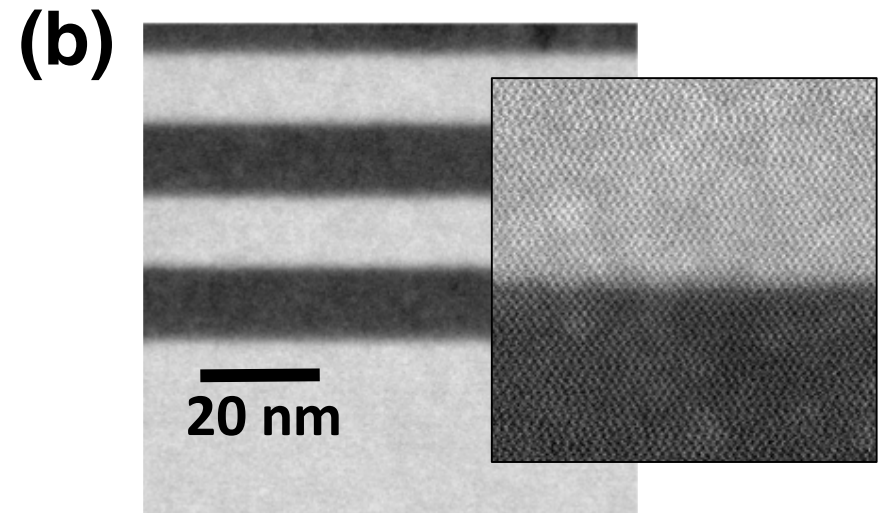
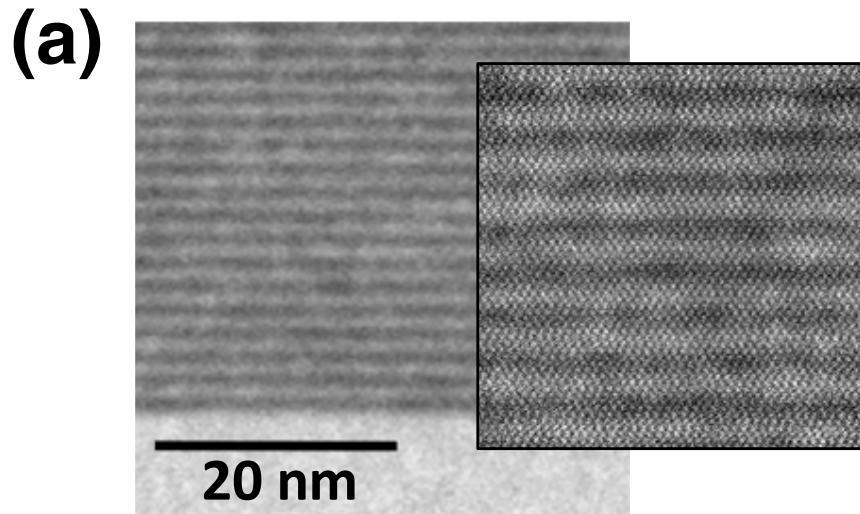
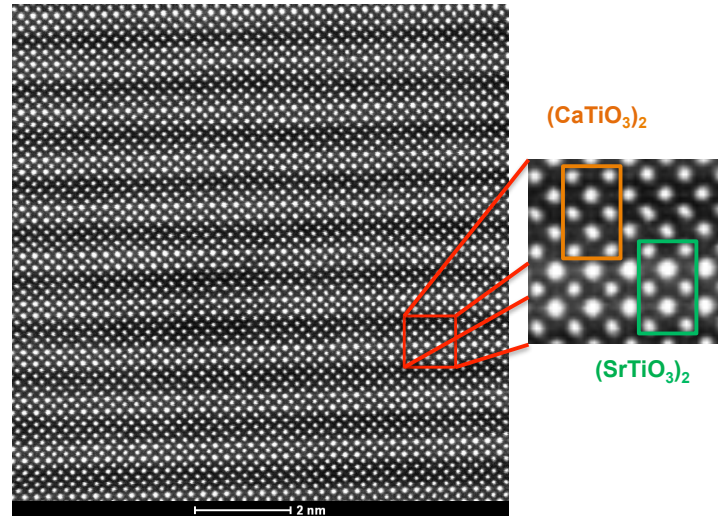
*J. Phys. Chem. Lett.* **5**, 1077

# Outline

- TDTR: Measurement of thermal conductivity of thin films and thermal resistance across interfaces
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- **Molecular interfaces in organic/inorganic composites: diffusive scattering via the vibron-phonon interaction**

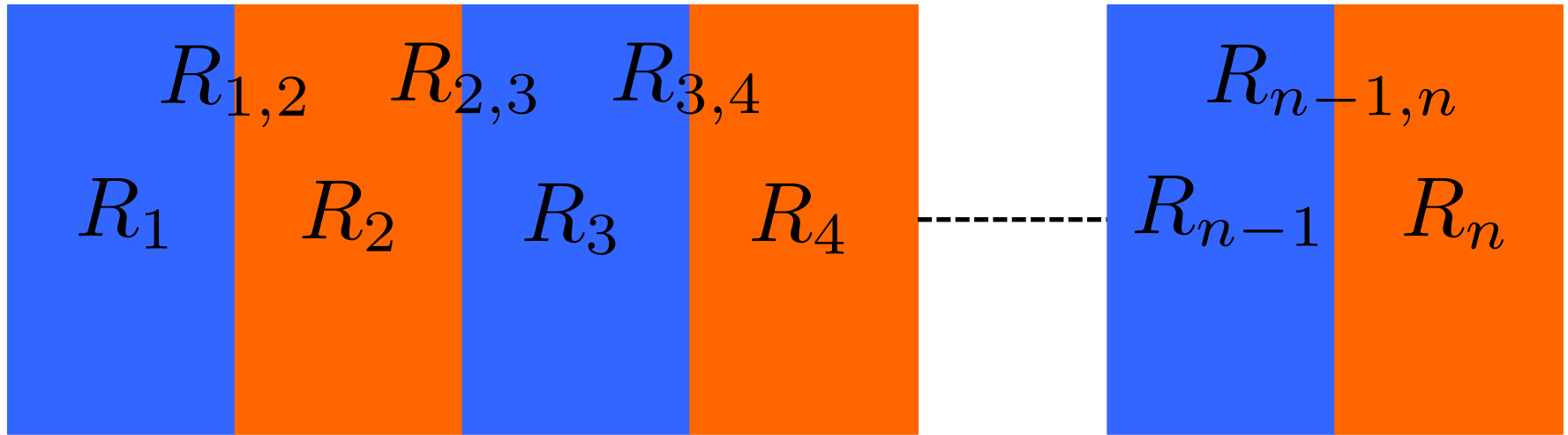
# Heat transport mechanisms in superlattices

**CaTiO<sub>3</sub>/SrTiO<sub>3</sub>**



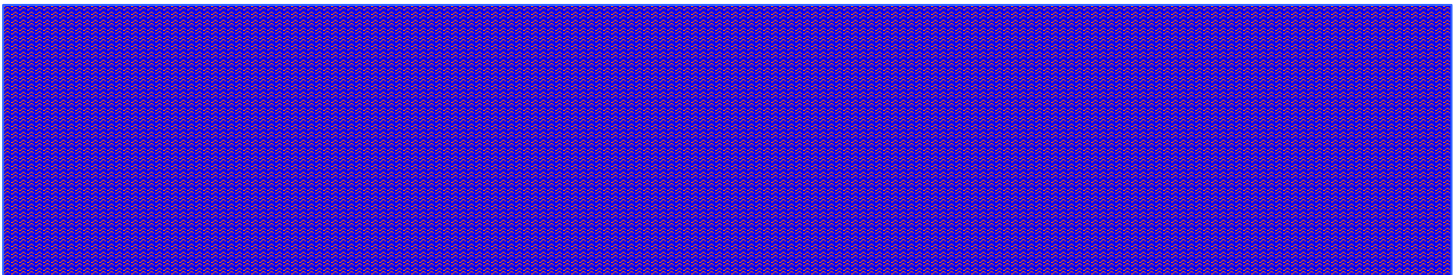
$d_{\text{SL}} = 2 \text{ nm}$  **GaAs/AlAs**  $d_{\text{SL}} = 24 \text{ nm}$

# Incoherent/particle picture of phonon transport in SLs



$$1/\kappa \propto 1/h = \sum_n R_j$$

What if layers are “linked”? – coherent transport



# Coherent transport in superlattices

PHYSICAL REVIEW B

VOLUME 25, NUMBER 6

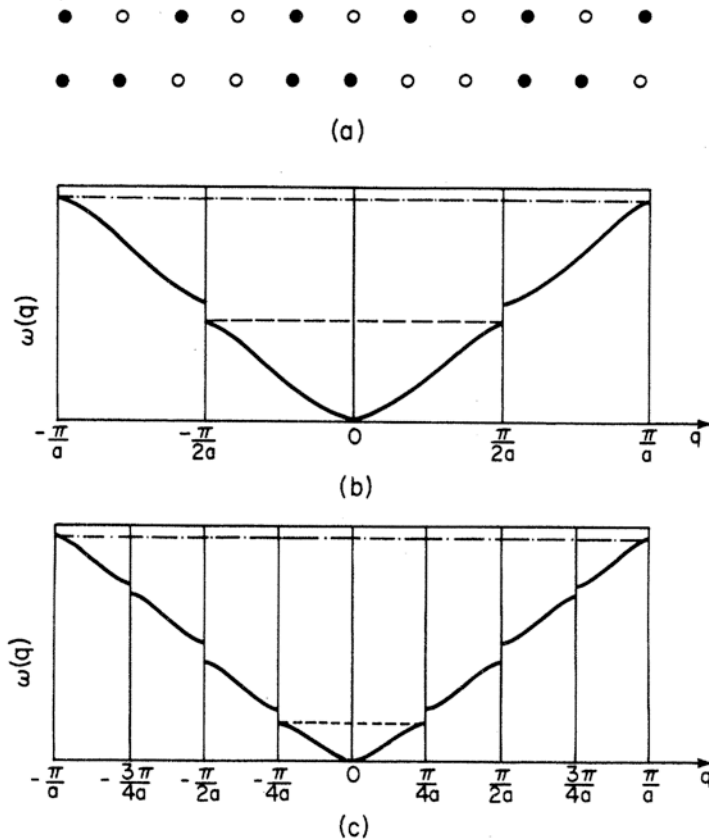
15 MARCH 1982

## Thermal conductivity of superlattices

Shang Yuan Ren\* and John D. Dow

Department of Physics and Coordinated Science Laboratory, University of Illinois at Urbana-Champaign,  
Urbana, Illinois 61801

(Received 21 September 1981)



Interfacial periodicity can  
lead to “mini-band” formation

VOLUME 84, NUMBER 5

PHYSICAL REVIEW LETTERS

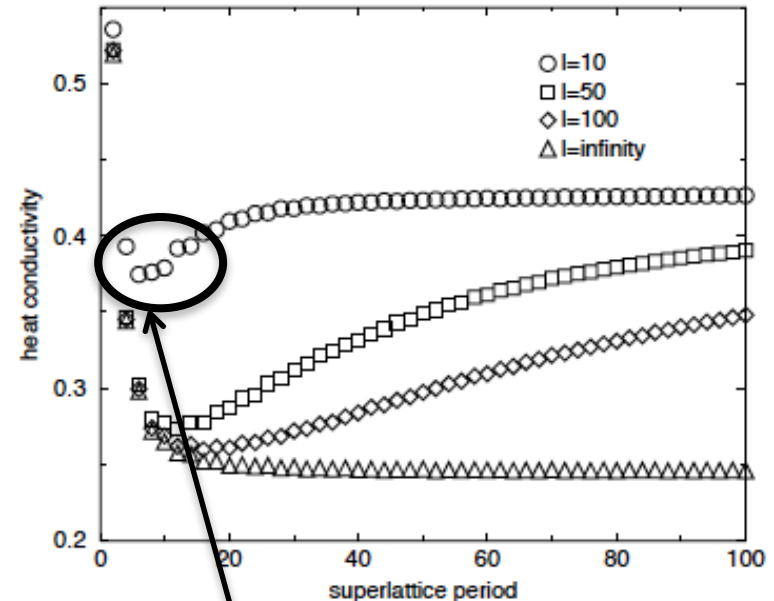
31 JANUARY 2000

## Minimum Thermal Conductivity of Superlattices

M. V. Simkin and G. D. Mahan

Department of Physics and Astronomy, University of Tennessee, Knoxville, Tennessee 37996-1200  
and Solid State Division, Oak Ridge National Laboratory, P.O. Box 2008, Oak Ridge, Tennessee 37831

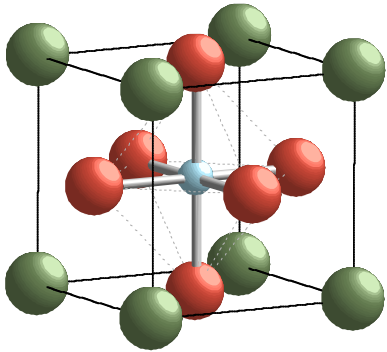
(Received 23 July 1999)



Mini-band formation leads to a  
minimum in the superlattice  
thermal conductivity

# Coherent transport in superlattices

$\text{SrTiO}_3$

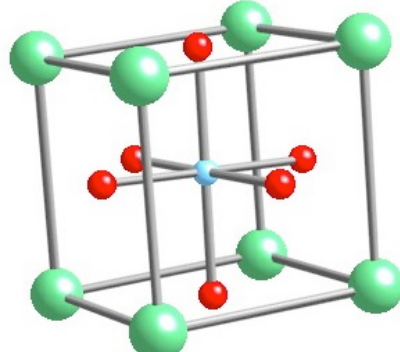


( $a = 3.905 \text{ \AA}$ )

$\rho = 5.1 \text{ g/cc}$

$v_m = 5.41 \text{ km/s}$

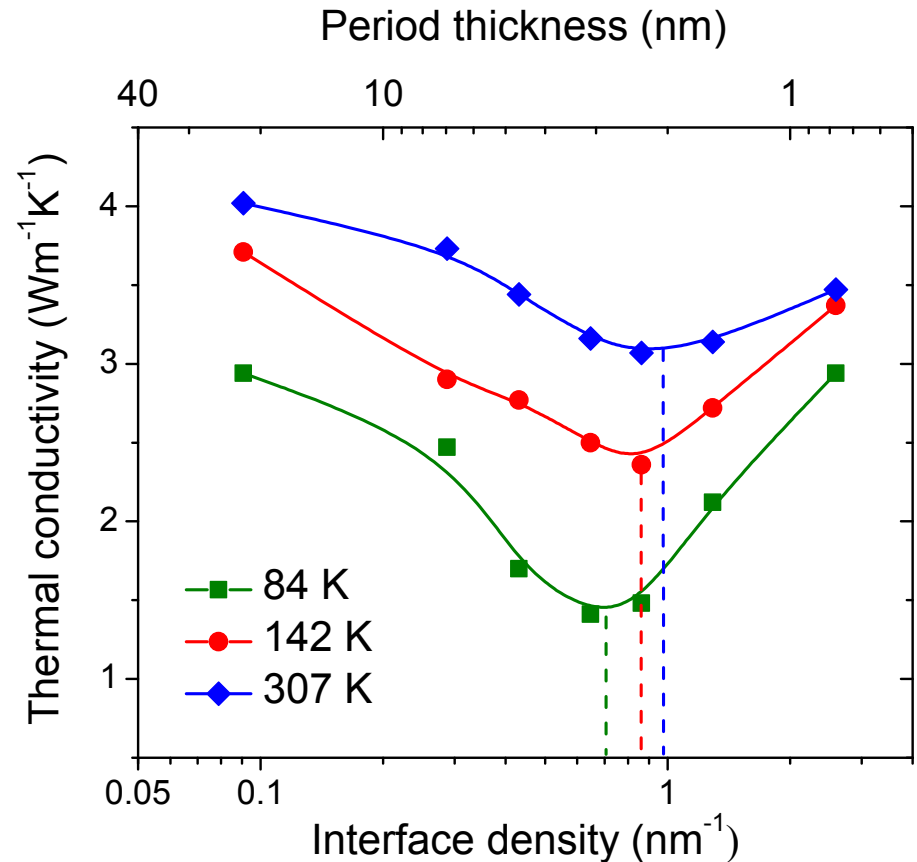
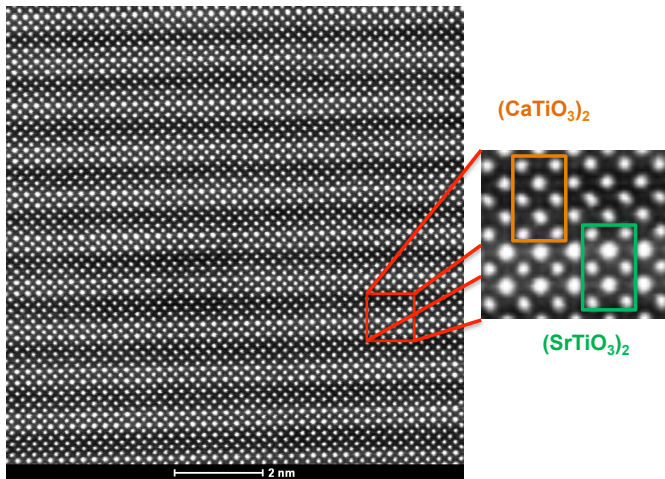
$\text{CaTiO}_3$



( $a_{pc} = 3.81 \text{ \AA}$ )

$\rho = 3.75 \text{ g/cc}$

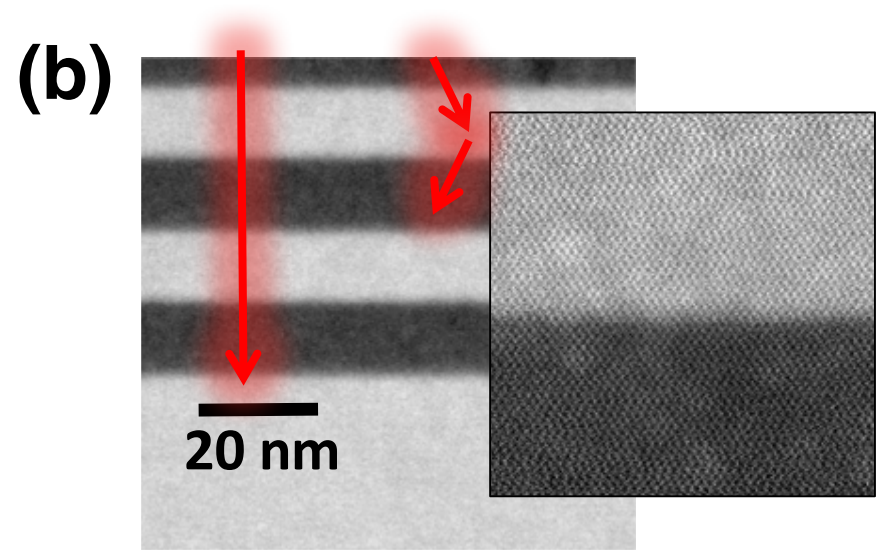
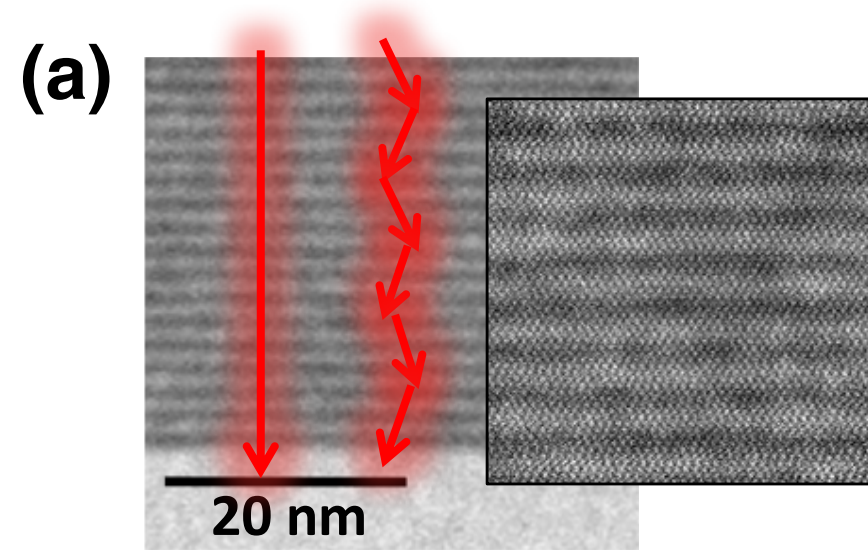
$v_m = 5.71 \text{ km/s}$



**SL design to manipulate  
coherent phonon transport**  
*Nature Materials* **13**, 168 (2013)



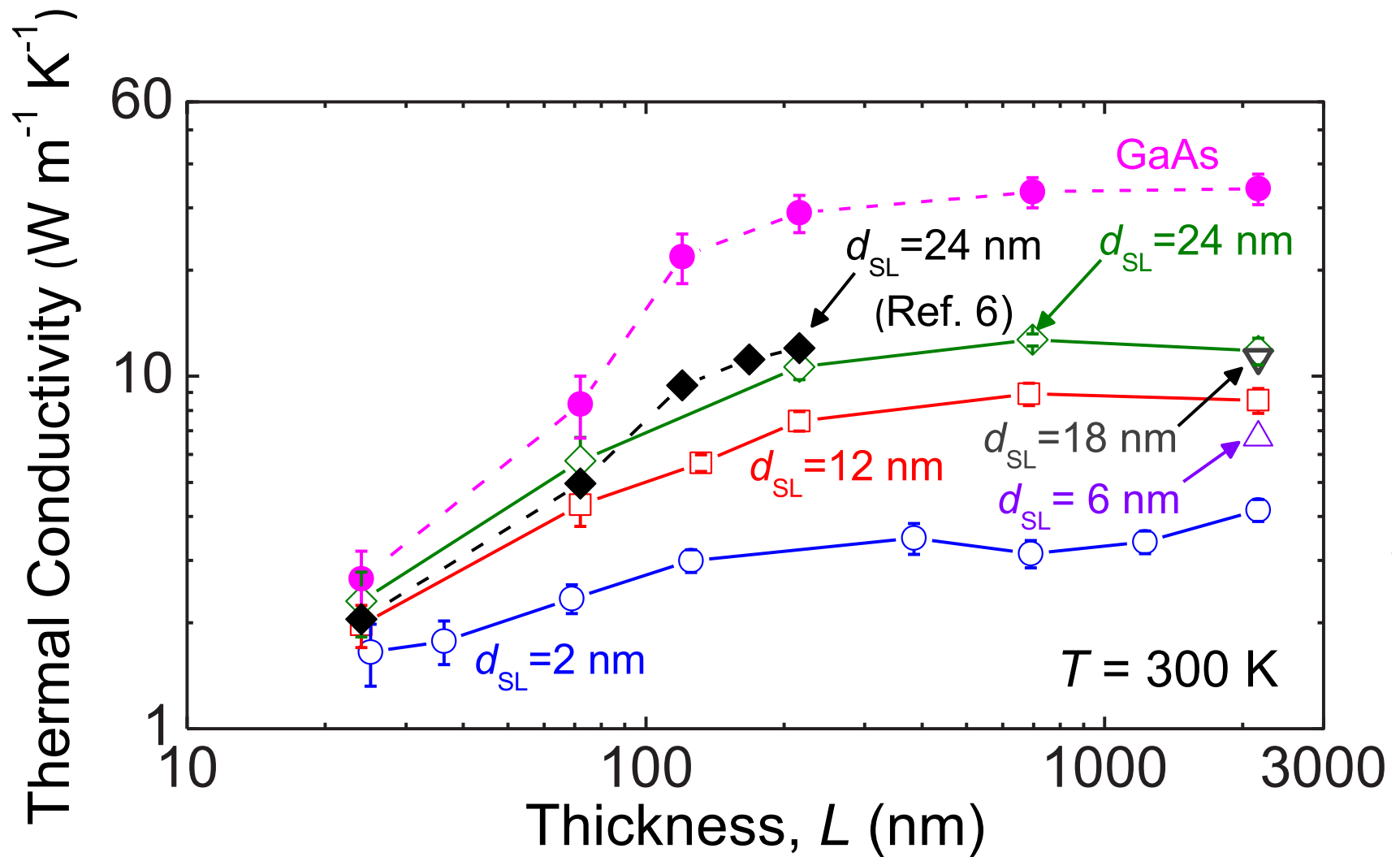
# Spectral phonon transport in SLs



**Do all phonons really  
scatter at interfaces???**

**What if wavelength of  
phonon is  $\gg$  than  
periodicity?**

# Spectral phonon transport in SLs

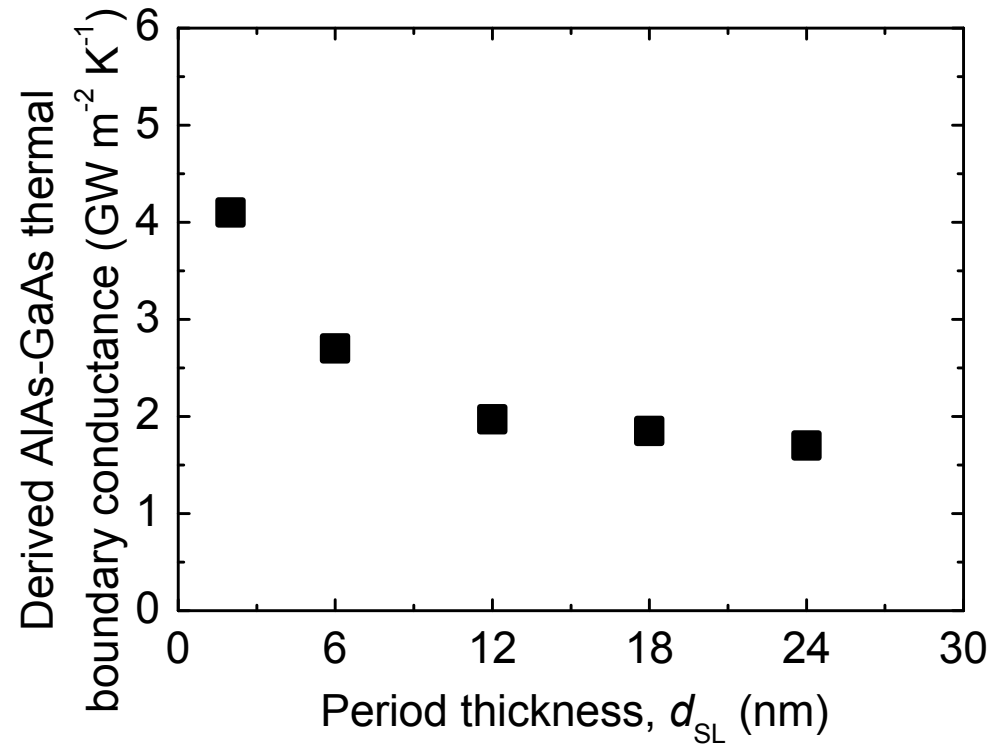
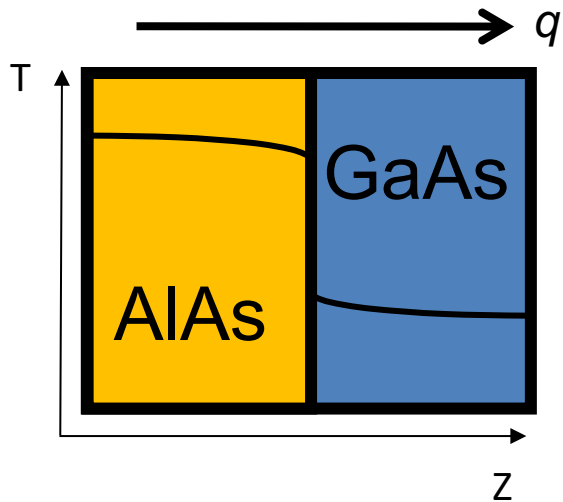


Collaboration: G. Balakrishnan (UNM)

*Phys. Rev. B* **97**, 085306

# TBC at AlAs/GaAs interface

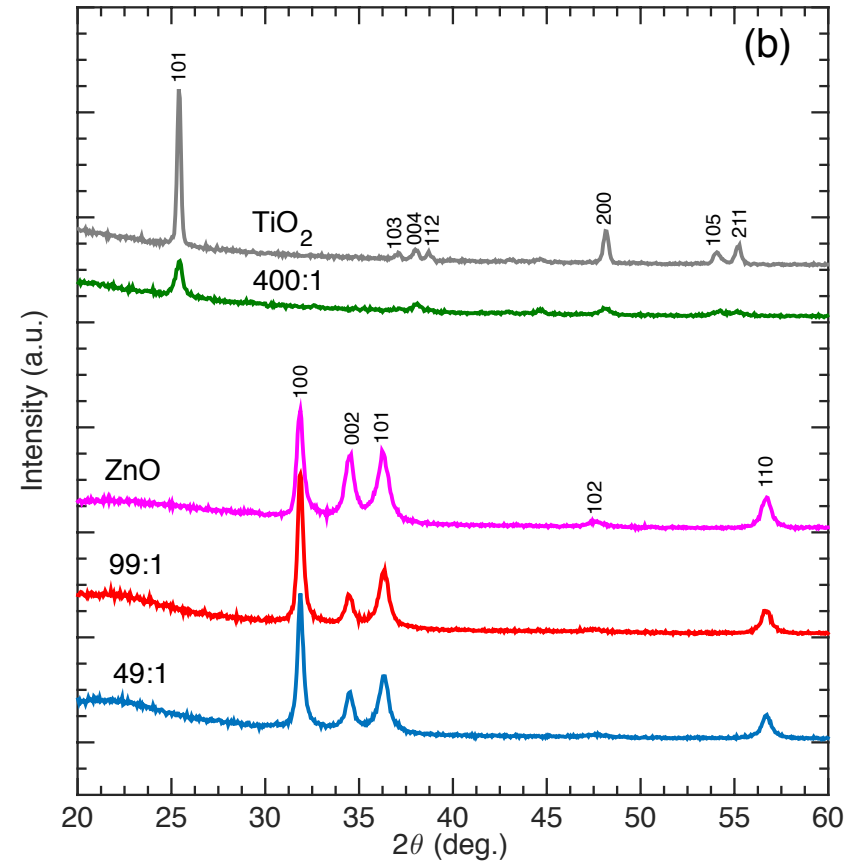
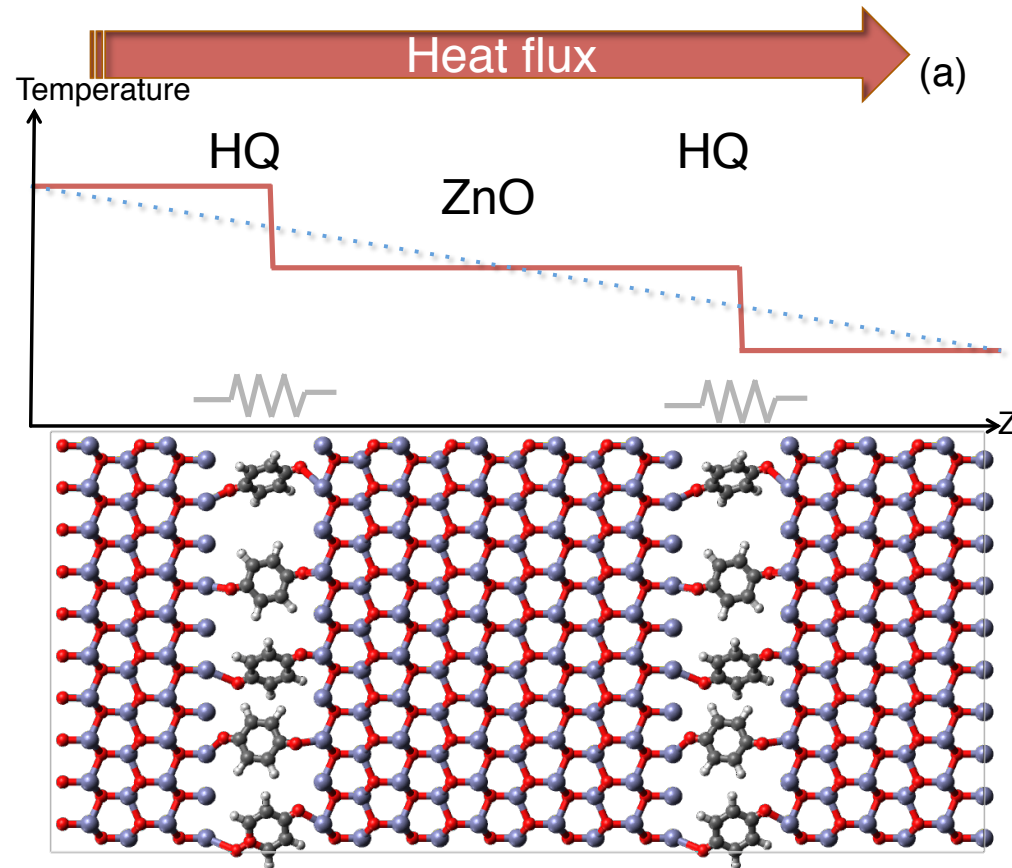
$$q = h_K \Delta T = \frac{1}{R_K} \Delta T$$



Collaboration: G. Balakrishnan (UNM)

*Phys. Rev. B* **97**, 085306

# Turn to molecular heterostructures

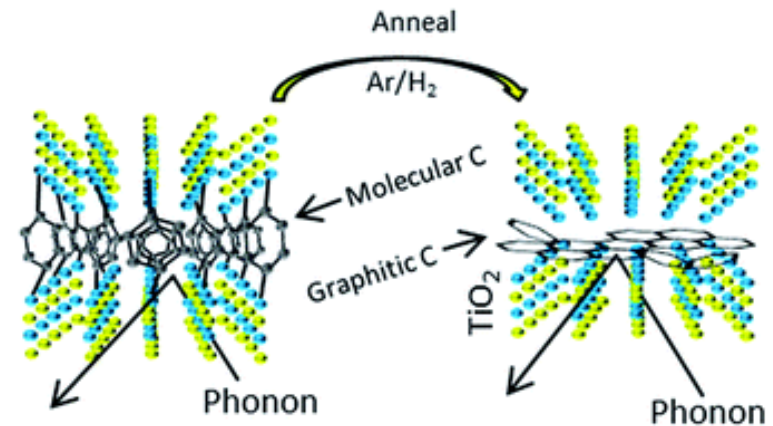
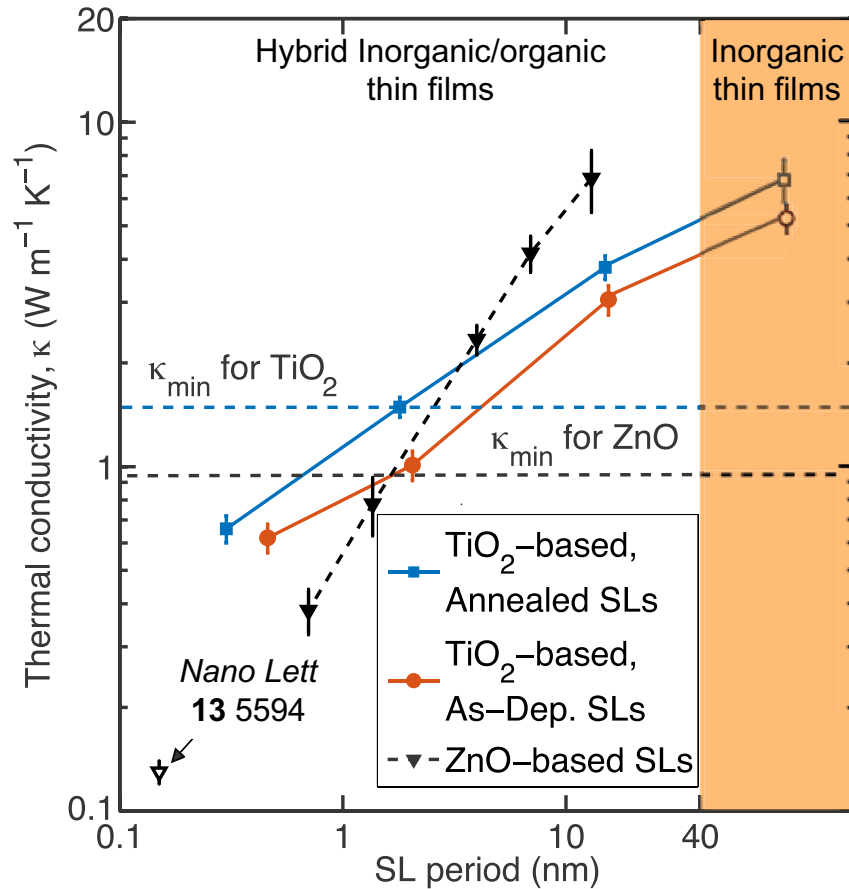


*PRB* **93**,115310; *PRB* **93**, 024201

**Collaboration:** M. Karppinen (Aalto) – ALD/MLD growth

# Turn to molecular heterostructures

Phonon scattering at organic/inorganic interface can lead crystalline composites achieving  $\kappa$  less than amorphous phase



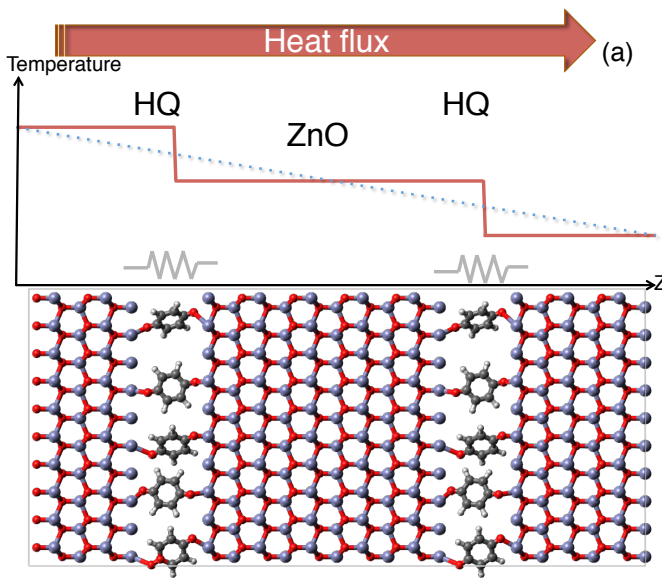
Reductions due to the thermal boundary conductance across organic/inorganic interfaces

*PRB* 93,115310; *PRB* 93, 024201

**Collaboration:** M. Karppinen (Aalto) – ALD/MLD growth

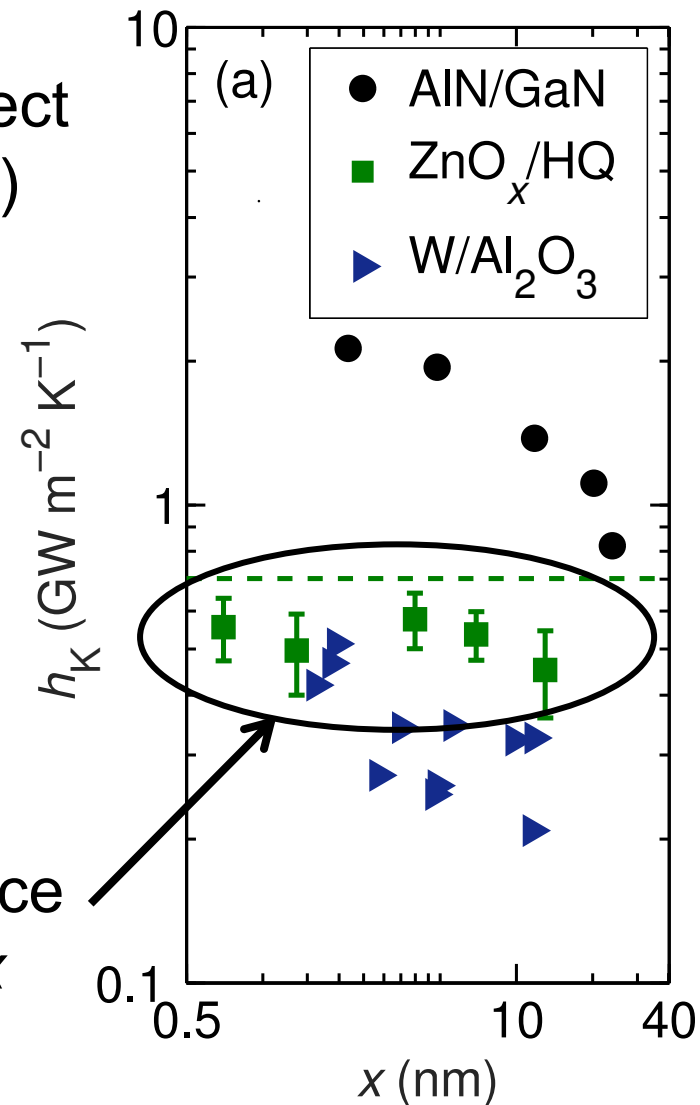
# Diffusive phonon scattering at organic/inorganic interface

Molecular interface causes all phonon modes to scatter at boundary (aka: no direct transmission of oxide modes across HQ)



$$q = h_K \Delta T = \frac{1}{R_K} \Delta T$$

Molecular interface constant with  $x$



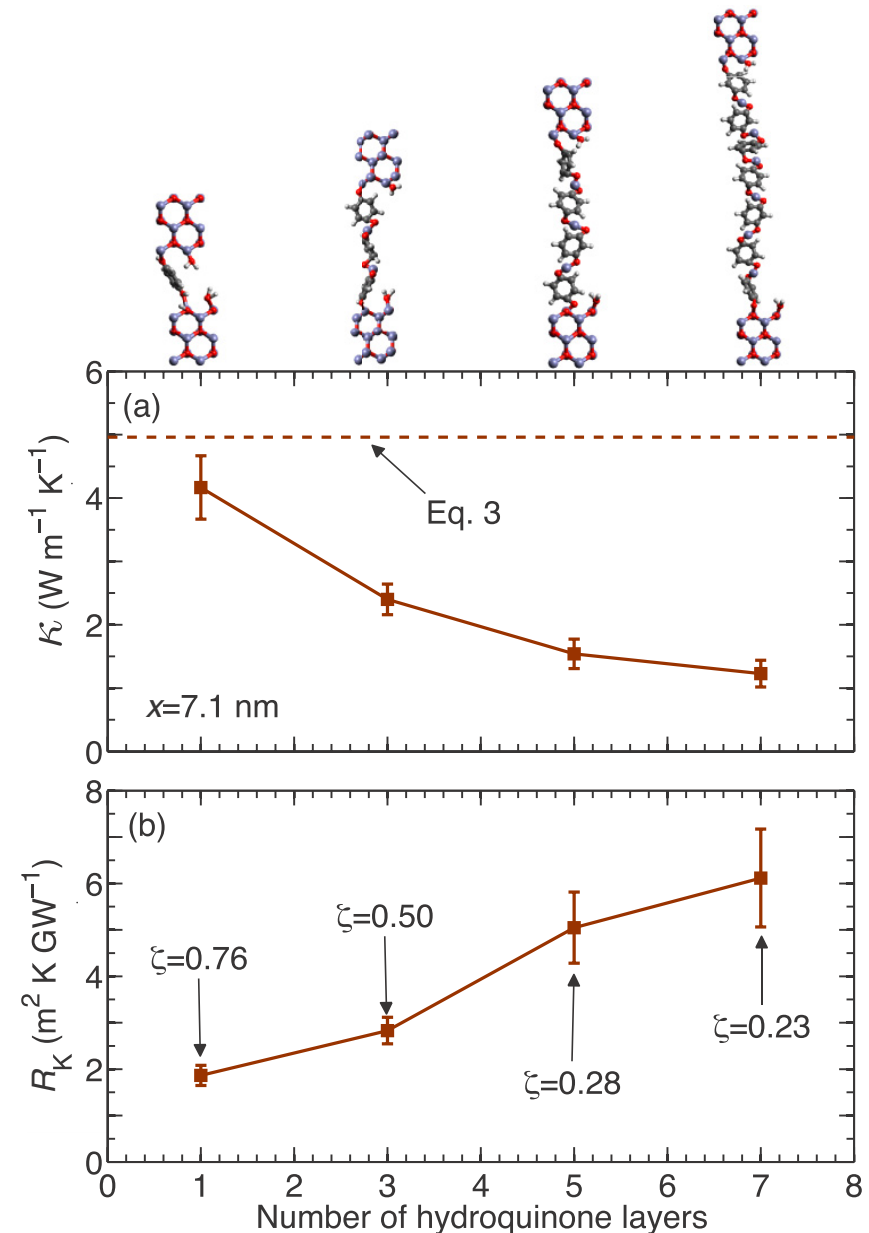
*PRB* **93**,115310; *PRB* **93**, 024201

**Collaboration:** M. Karppinen (Aalto) – ALD/MLD growth

# Diffusive phonon scattering at organic/inorganic interface

Can use hybrid materials to study diffusive scattering in few molecule thick films

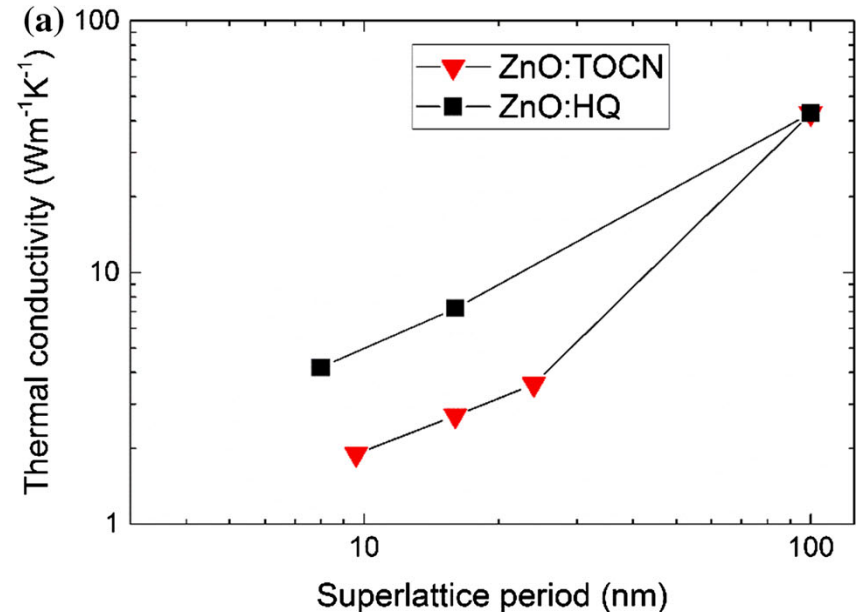
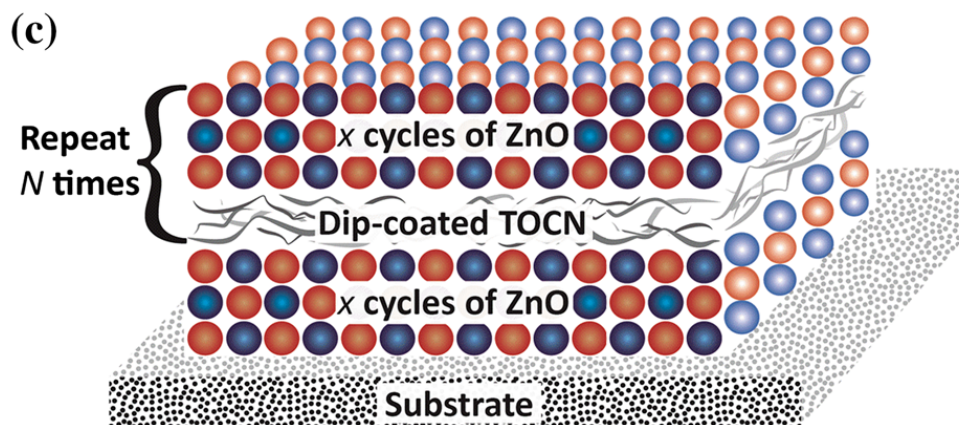
*PRB* **93**,115310; **Collaboration:** M. Karppinen (Aalto) – ALD/MLD growth





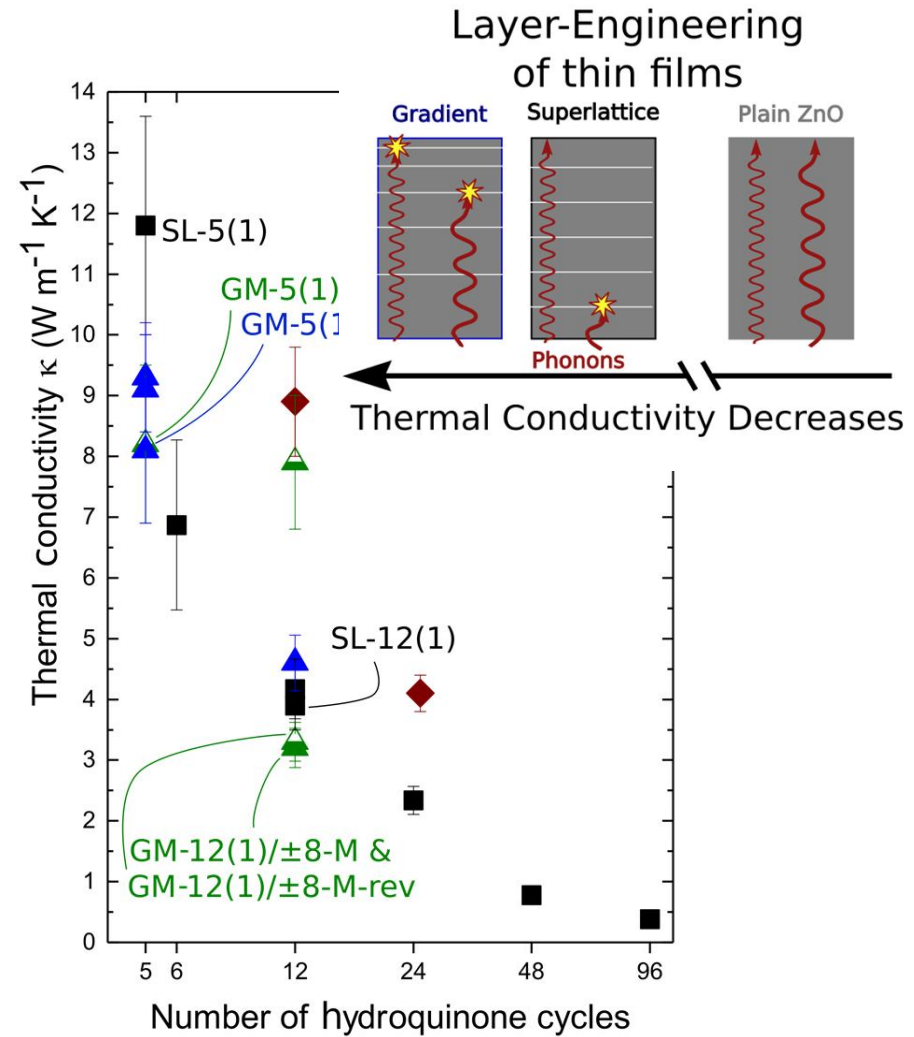
# Heat transport mechanisms in novel hybrid materials

## ZnO/nanocellulose



*J. Mat. Sci.* **52**, 6093

## A periodic hybrid “SLs”



*Appl. Mat. & Int.* (to appear)

# Conclusions/outlook – Heterogeneous material interfaces and engineering the chemical bond can lead to novel regimes of vibrational heat transport in materials

